

MAGIS-100: Large Baseline Atom Interferometry

Jeremiah Mitchell

University of Cambridge
Kavli Foundation

28 April 2021



UNIVERSITY OF
CAMBRIDGE



THE
KAVLI
FOUNDATION

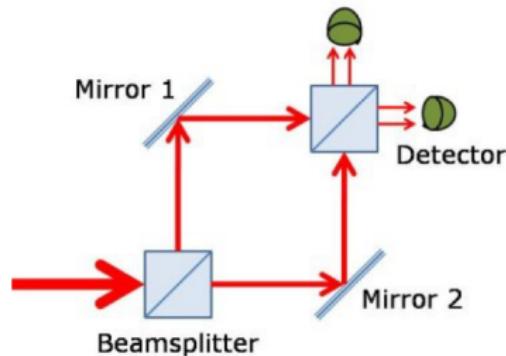
① Overview Atom Interferometry

② MAGIS-100 and AION

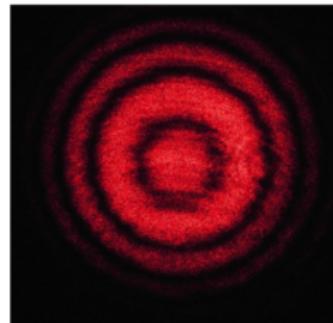
③ Systematics and Simulation

 Systematics and Error Analysis
 Modeling and Simulations

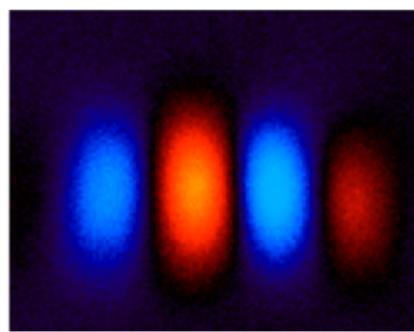
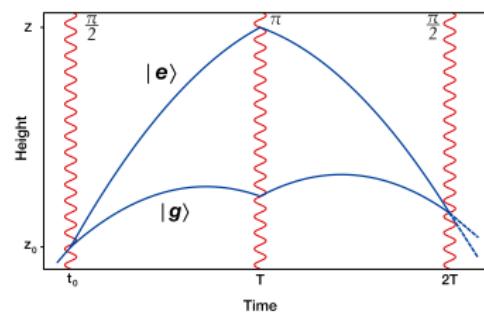
④ Summary



<http://www.cobolt.se/interferometry.html>

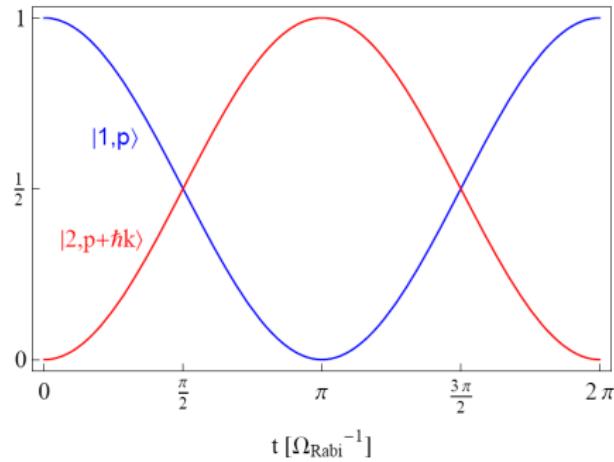
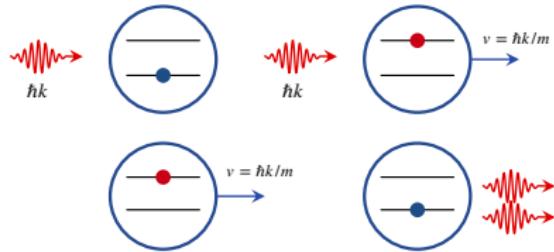


Laser interference



Matter-wave interference

- Atom optics implemented through stimulated absorption/emisison
- Raman scattering or Bragg scattering modes
- Two photon and single photon transitions



$$\Delta\phi = \Delta\phi_{\text{laser}} + \Delta\phi_{\text{propagation}} + \Delta\phi_{\text{separation}}$$

$$\Delta\phi_{\text{laser}} = \sum_j^{\text{upper}} \pm\phi_L(t_j, \mathbf{x}_u(t_j)) - \sum_j^{\text{lower}} \pm\phi_L(t_j, \mathbf{x}_l(t_j))$$

$$\Delta\phi_{\text{propagation}} = \sum_{\text{upper}} \left(\int_{t_i}^{t_f} (L_c - E_i) dt \right) - \sum_{\text{lower}} \left(\int_{t_i}^{t_f} (L_c - E_i) dt \right)$$

$$\Delta\phi_{\text{separation}} = \langle \mathbf{p} \rangle \cdot \Delta \mathbf{x}$$

$$\phi_L(t, \mathbf{x}(t)) = \mathbf{k}_{\text{eff}} \cdot \mathbf{x}(t) - \omega t + \delta$$

Matter-wave Atomic Gradiometer Interferometric Sensor (MAGIS-100)

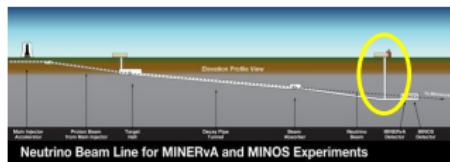
Current a 10 m EP tower and
10 m prototype
Science Reach

- Measured Earth's rotation with extreme precision
- Detected local gravity gradient of moving test-masses and testing Equivalence Principle ($< 10^{-12}$)
- Macroscopic separation of quantum states of atom ensembles (54 cm)

MAGIS-100 uses 100 m baseline with three atom sources

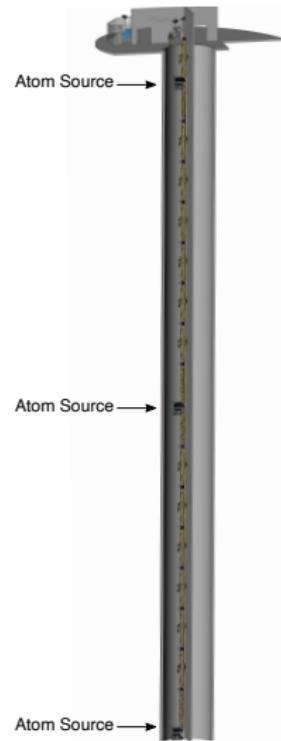
Science Goals

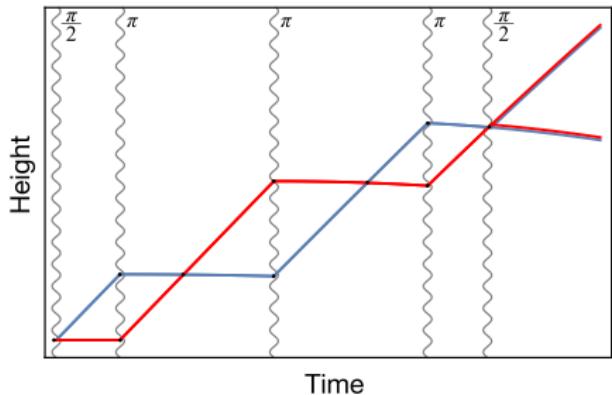
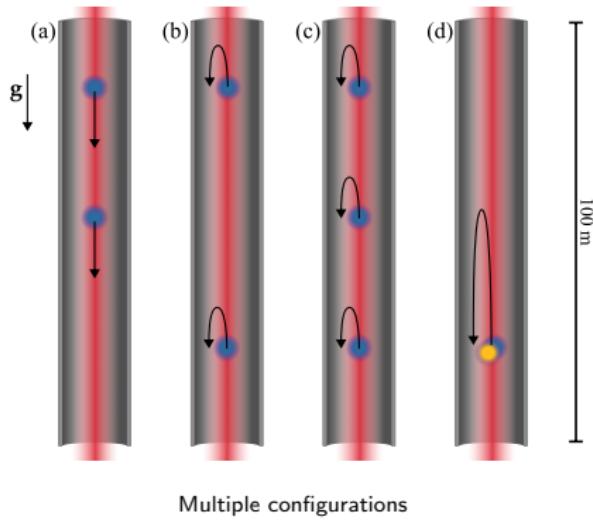
- Dark matter detection
- Gravitational wave detector prototype
- Extreme macroscopic quantum mechanics



Abe, Mahiro, et al. QST (2021).

arXiv:2104.02835v1

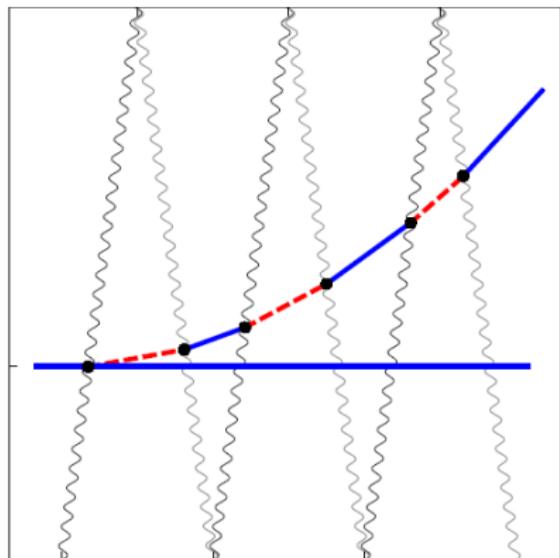




Large momentum transfer (LMT) and multi-loop atom
interferometry

- Multiple laser pulses to accelerate one of the atom interferometer arms
- Coherently enhance differential phase measurement scales with n the number of LMT pulses

$$\Delta\phi \sim 2n\omega_A(L)$$



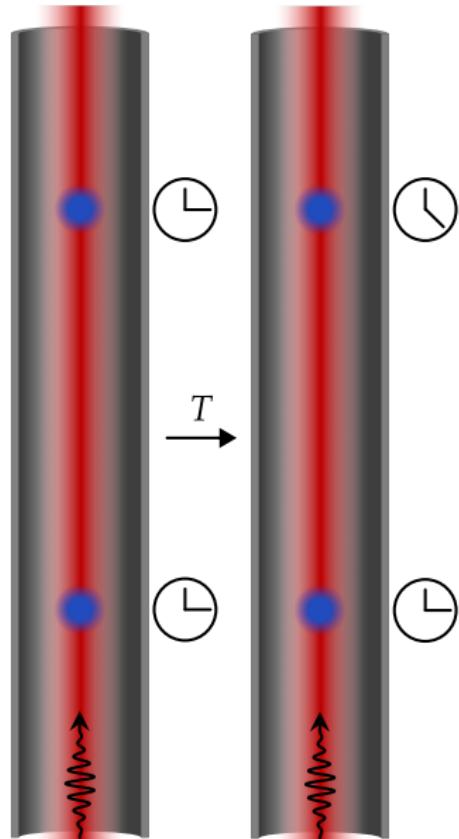
Graham, et al. PRL (2013)

- Clocks are started simultaneously ticking with frequency ω_A
- Fluctuation of laser alters clock measurement of time
 $T \rightarrow T + \Delta T$

Initial States $\frac{1}{\sqrt{2}} |g\rangle + \frac{1}{\sqrt{2}} |e\rangle$

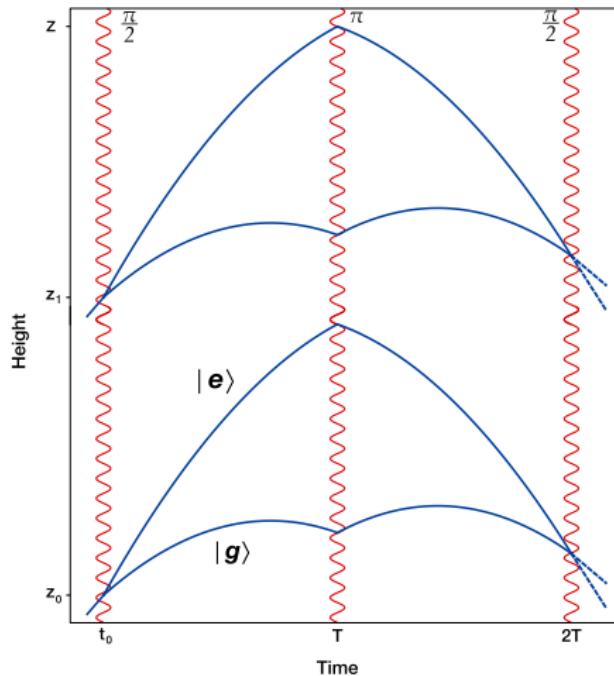
Lower Source $\frac{1}{\sqrt{2}} |g\rangle + \frac{1}{\sqrt{2}} |e\rangle e^{-i\omega_a T}$

Upper Source $\frac{1}{\sqrt{2}} |g\rangle + \frac{1}{\sqrt{2}} |e\rangle e^{-i\omega_a(T+\Delta T)}$



- Simultaneous interferometer interference measurement
- Common laser allows common mode noise suppression

$$\frac{1}{\sqrt{2}} |g_0\rangle + \frac{1}{\sqrt{2}} |e_0\rangle e^{i\phi_L}$$
$$\frac{1}{\sqrt{2}} |g_1\rangle + \frac{1}{\sqrt{2}} |e_1\rangle e^{i\phi_L}$$

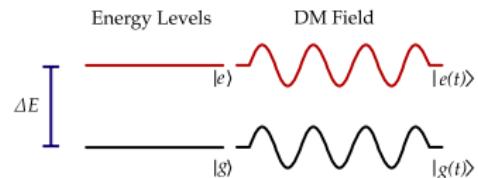


Dark Matter

- Most sensitive to general class of ultra-light scalar field dark matter
 - Time dependant fluctuation of fundamental constants (fine structure, electron mass, gravitational g)
 - Accelerations caused by gradients of field
- Includes QCD Axion, axion-like, and relaxion

Gravitational Waves

- Atoms interact with common laser beam, gravitational waves will alter the laser travel time between 2 atom clouds
- Gradiometer measurement of atoms in atomic clock modes will show phase difference

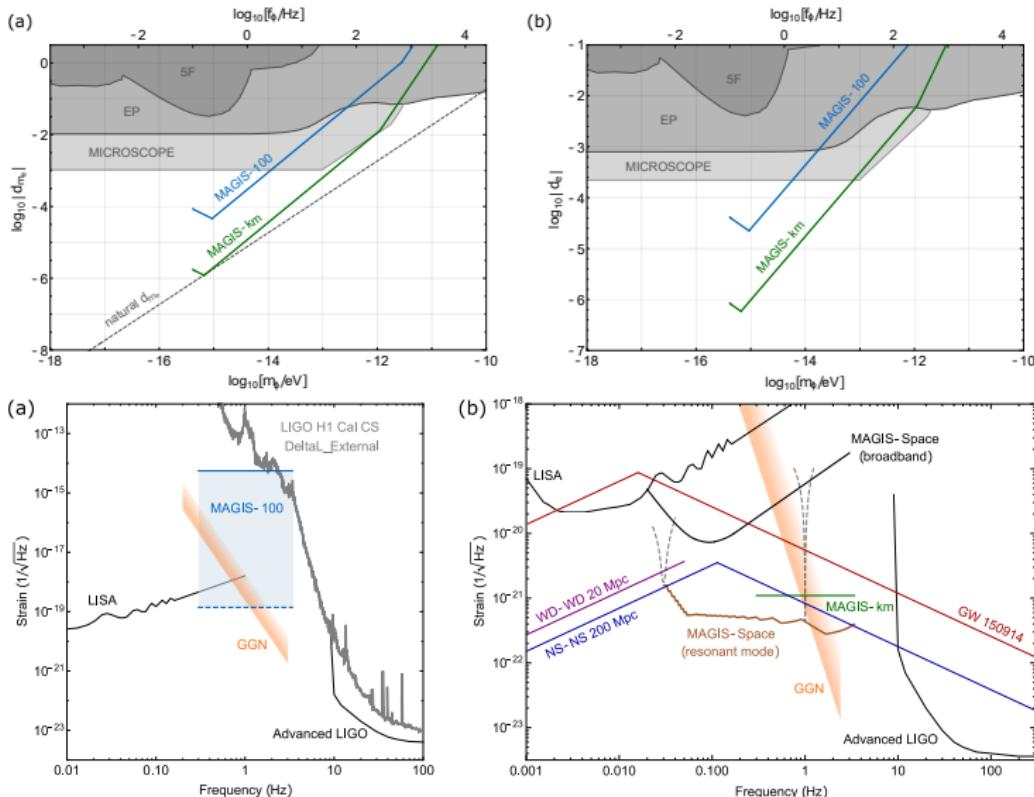


$$\Delta\phi \sim \omega_A(2L/c)$$

$$\text{Dark Matter} \quad \delta\omega_A$$

$$\text{Gravitational Waves} \quad \delta L = hL$$

MAGIS-100 Sensitivity



Abe, Mahiro, et al. QST (2021). arXiv:2104.02835v1

- Atom Interferometry Observatory Network (AION) in the UK
- Plans to work complementary to MAGIS and network the detectors together for enhanced science goals
- Will use similar technologies as MAGIS
- Provides uncommon noise background suppression
- Adds to the verification of science measurements

Main subsystems to consider for systematics and noise sources

- Laser system
- Atom Sources
- Detector environment
- Detection and readout

① Systematics and Error Analysis

- Coriolis
- Laser Wavefront Noise
- Seismic Environment Characterization
- Gravity Gradient Noise Analysis

② Modeling and Simulation

- Gravity Gradient Noise Model
- MAGIS-100 simulation for R&D

- Crucial for maximizing detector sensitivity
- For test bed systems like MAGIS-100 plays important R&D role for long baseline atom interferometers
- Goal is to understand leading order systematics and devise methods for mitigating them
- Interesting leading order effects: Coriolis effect, laser wavefront perturbations, gravity gradients, gravity gradient noise (GGN)
- Full list can be found in [Abe, Mahiro, et al. QST (2021). arXiv:2104.02835v1]

Deflection caused by atoms in inertial free-fall reference frame embedded in accelerating frame of Earth

- Causes atoms to fall out of interferometry laser beam, leads to decrease in contrast
- Atoms pick up phase shift during propagation

$$\mathbf{a}_c = -2\boldsymbol{\Omega} \times \mathbf{v}. \quad (1)$$

$$\boldsymbol{\Omega} = \begin{pmatrix} 0 \\ \omega \cos(\phi) \\ \omega \sin(\phi) \end{pmatrix}, \quad \mathbf{v} = \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix}, \quad (2)$$

$$\Delta \mathbf{r} = \omega \begin{pmatrix} v_y \sin(\phi)t^2 - \left[v_{0,z}t^2 - g\frac{t^3}{3} \right] \cos(\phi) \\ -v_x \sin(\phi)t^2 \\ v_x \cos(\phi)t^2 \end{pmatrix}. \quad (3)$$

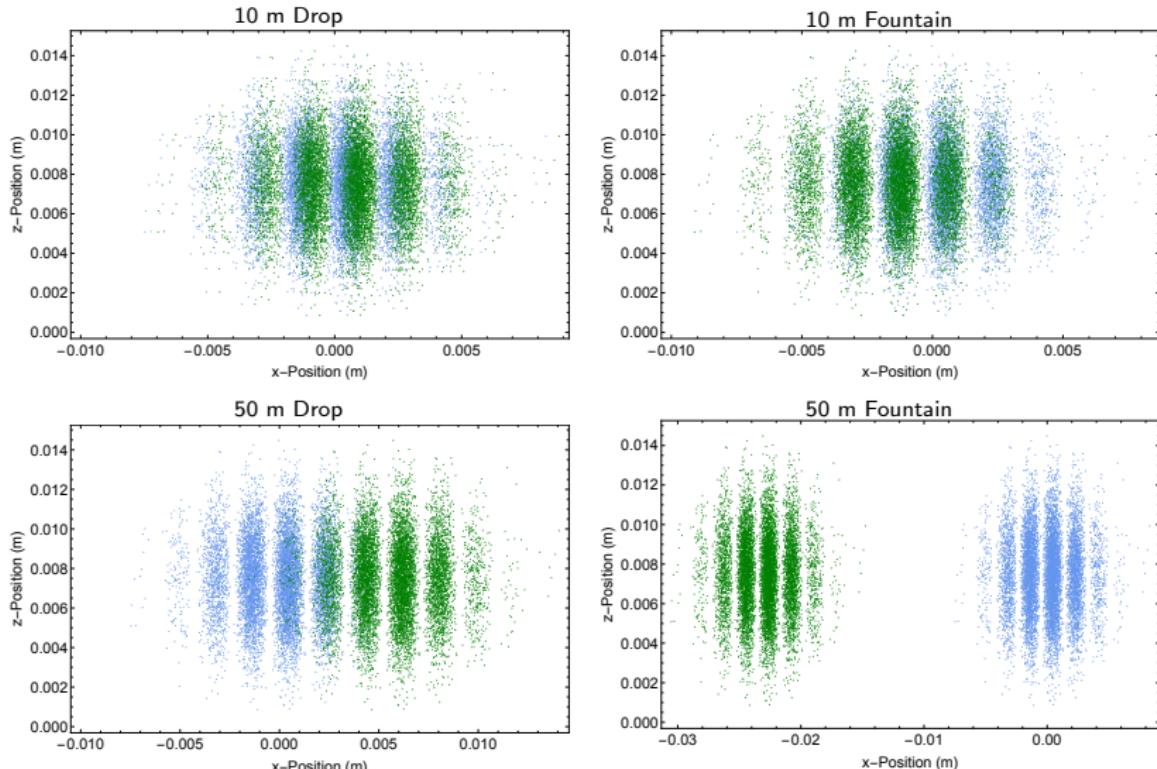
Table 1: Table of flight times for various configurations of atom launch schemes.

Configuration	T	$2T$
100 m Drop	2.26 s	4.52 s
100 m Fountain	4.52 s	9.03 s
50 m Drop	1.595 s	3.19 s
50 m Fountain	3.19 s	6.38 s
10 m Drop	0.72 s	1.43 s
10 m Fountain	1.43 s	2.86 s

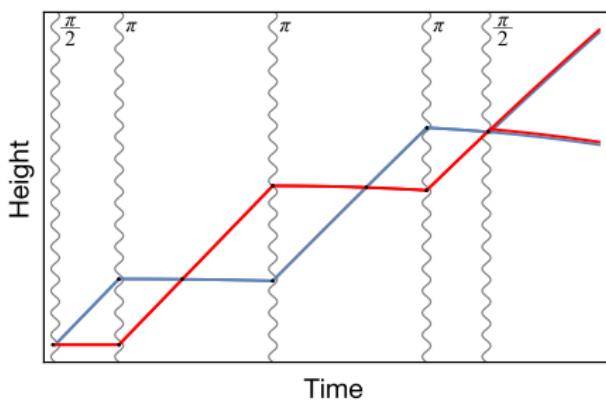
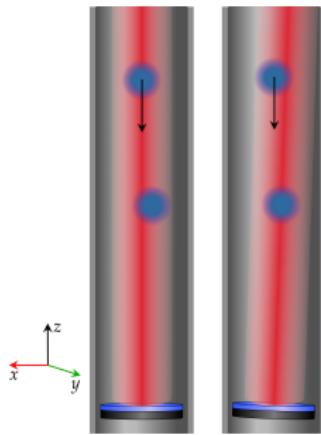
Table 2: The calculation was performed with $\omega = 7.27 \times 10^{-5}$ rad/s, $v_x = v_y = 1$ mm/s, and initial launch velocities for the fountain configurations of $v_0 = 44.29$ m/s, 31.32 m/s and 13.03 m/s for the 100 m, 50 m, and 10 m fountains respectively.

Configuration	Δx	Δz
100 m Drop	1.64 cm	1.1×10^{-6} m
100 m Fountain	-6.52 cm	4.4×10^{-6} m
50 m Drop	5.75 mm	5.5×10^{-7} m
50 m Fountain	-2.31 cm	2.2×10^{-6} m
10 m Drop	0.518 mm	1.1×10^{-7} m
10 m Fountain	-1.63 mm	4.4×10^{-7} m

Blue is unperturbed position of atom cloud during imaging, Green is position affected by unmitigated Coriolis



- Primary suppression from tip-tilt mirror rotating counter to Earth's rotation
- Multiloop sequences that can be designed for compensation of Coriolis phase shift terms



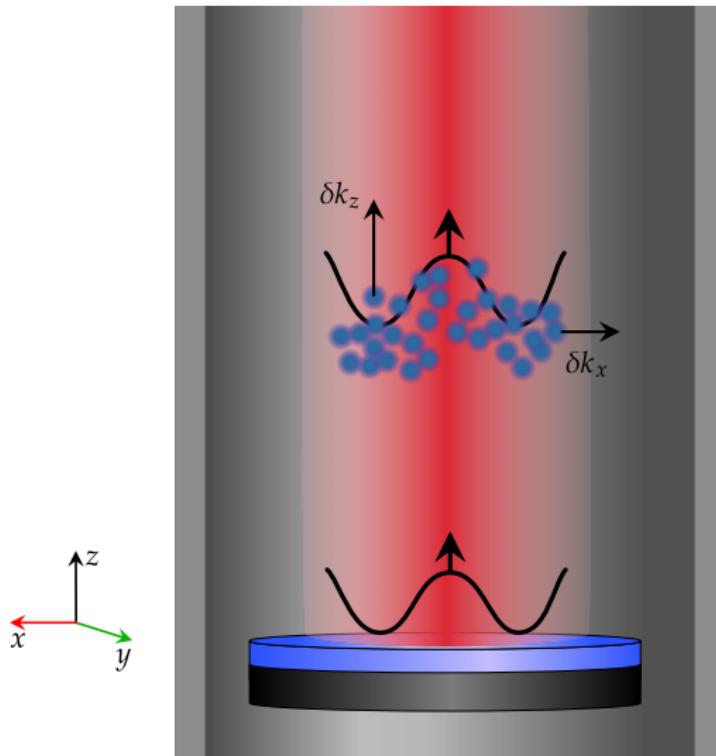
- Amplitude perturbations of the general form $E(x, z) = u(x, z)e^{ikz}$
- We then consider the initial amplitude $u(x, 0) = 1 + \delta \cos(k_x x)$ and the propagated amplitude $u(x, z) = 1 + \delta \cos(k_x x)e^{i\frac{k_x^2}{2k}z}$
- Paraxial approximation $k_z = \sqrt{k^2 - k_x^2} \approx k - \frac{k_x^2}{2k}$

$$\begin{aligned} \text{Amplitude: } & \delta \cos(k_x x) \cos\left(\frac{k_x^2}{2k} z\right), \\ \text{Phase: } & \phi_w = \delta \cos(k_x x) \sin\left(\frac{k_x^2}{2k} z\right), \\ \delta k_x = & \frac{\partial \phi_w}{\partial x} = -k_x \delta \sin(k_x x) \sin\left(\frac{k_x^2}{2k} z\right), \\ \delta k_z = & \frac{\partial \phi_w}{\partial z} = \frac{k_x^2}{2k} \delta \cos(k_x x) \cos\left(\frac{k_x^2}{2k} z\right). \end{aligned} \tag{4}$$

Considered effective momentum kicks because terms scale as

$$\frac{\hbar \delta k_i^2}{2m} \quad i = x, z \tag{5}$$

Laser Wavefront Perturbations

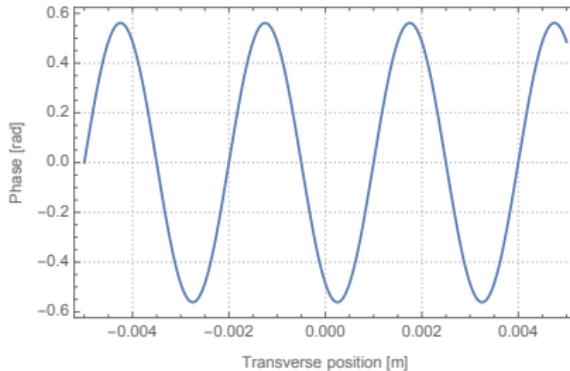
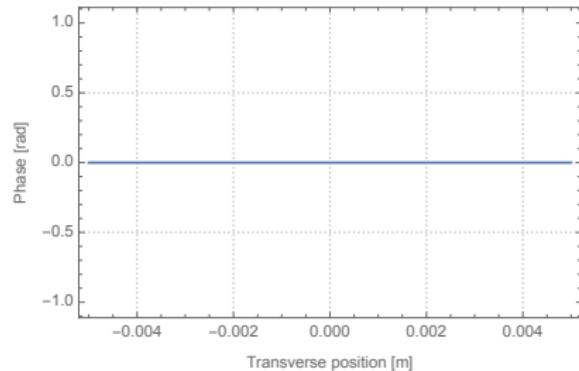


Laser Wavefront Phase Terms cont.

Table 3: Selection of example values for wavefront phase shifts at the first and second order in δ .

Term	Phase shift	Notes
1	$n\delta \cos(k_t x_i) \sin\left(\frac{k_t^2}{2k} H\right)$	First Order in δ
1	$(0.325 \text{ rad}) \left(\frac{n}{100}\right) \left(\frac{\delta}{0.005}\right) \left(\frac{\sin(k_t^2 H/2k)}{0.65}\right)$	
2	$-\frac{k_t^4 n^2 \delta^2 T \hbar \cos(k_t x_i) \cos\left(\frac{k_t^2}{2k} H\right) \cos(k_t T v_t + k_t x_i) \cos\left(\frac{g k_t^2 T^2}{4k} - \frac{k_t^2 T v_z}{2k} - \frac{k_t^2}{2k} H\right)}{8k^2 m}$	Longitudinal Kicks
2	$(-1.25 \times 10^{-11} \text{ rad}) \left(\frac{n}{100}\right) \left(\frac{\delta}{0.005}\right) \left(\frac{\cos(k_t^2 H/2k)}{0.7}\right)$	
3	$\frac{k_t^2 n^2 \delta^2 T \hbar \sin(k_t x_i) \sin\left(\frac{k_t^2}{2k} H\right) \sin(k_t T v_t + k_t x_i) \sin\left(\frac{g k_t^2 T^2}{4k} - \frac{k_t^2 T v_z}{2k} - \frac{k_t^2}{2k} H\right)}{2m}$	Transverse Kicks
3	$(-4.33 \times 10^{-9} \text{ rad}) \left(\frac{n}{100}\right) \left(\frac{\delta}{0.005}\right) \left(\frac{\sin(k_t^2 H/2k)}{-0.2}\right)$	

- Apply laser beam tilt of $d_{\text{atom}} \theta$ to final pulse from top or bottom tip-tilt mirror
- Allows measuring first derivative of wavefront via spatially resolved imaging
- Apply transverse initial velocity to atom cloud for information about wavefront in direction of kick

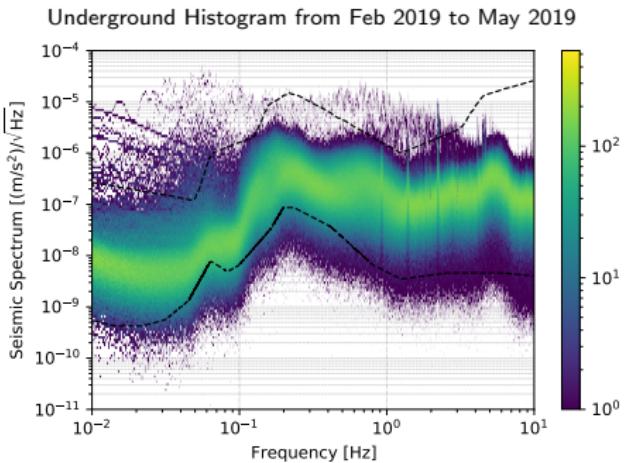
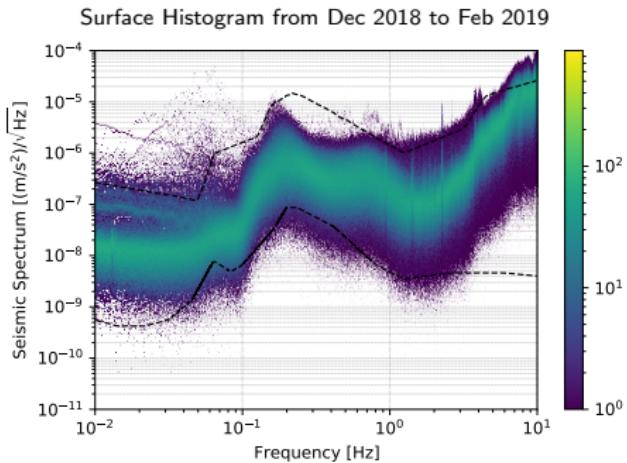


- To understand noise limits at Fermilab a state-of-the-art seismometer was purchased and installed on site
- Data was taken over the course of two years at one location underground and four surface locations
- Power spectra from the ambient seismic environment were made to understand vibrational noise sources as well as gravity gradient noise

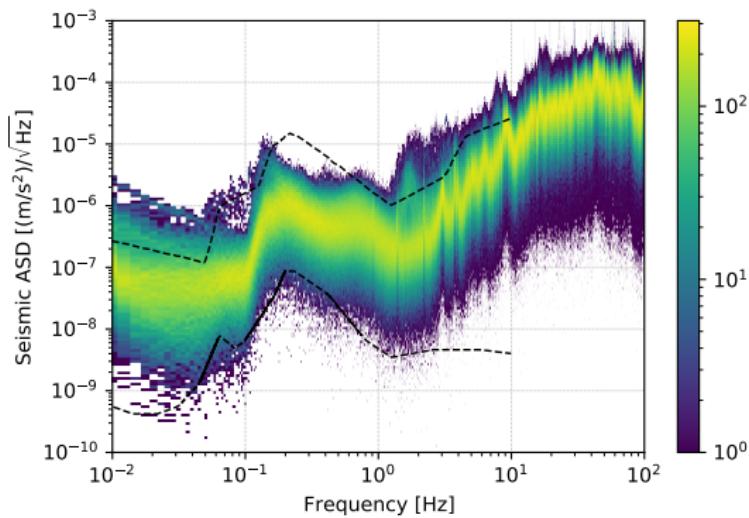


Ambient Seismic Noise

$$\begin{aligned} PSD\{x(t)\} &= \frac{2(\Delta t)^2}{T} |x(\omega)|^2, \\ x(\omega) &= DFT\{x(t)\} = \sum_{t=0}^{N-1} x(t) e^{2\pi i t/N}. \end{aligned} \quad (6)$$



Measurements at proposed location of initial laser transport pedestal



- Vibration couples to lasers and initial atom ensemble kinematics

$$\delta\phi_{\text{vibration}} \sim \left(10^{-9} \text{ rad}/\sqrt{\text{Hz}}\right) \left(\frac{n}{100}\right) \left(\frac{\Delta\nu}{10 \mu\text{m}/\text{s}}\right) \left(\frac{T}{1\text{s}}\right) \left(\frac{\delta a}{10^{-4} \text{ m}/\text{s}^2/\sqrt{\text{Hz}}}\right) \quad (7)$$

$$\delta\phi_{\text{RGGV}} \sim \left(2 \times 10^{-6} \text{ rad}/\sqrt{\text{Hz}}\right) \left(\frac{n}{100}\right) \left(\frac{\Delta\nu_x}{1 \mu\text{m}/\text{s}/\sqrt{\text{Hz}}}\right) \left(\frac{T}{1\text{s}}\right) \quad (8)$$

Can be mitigated by vibrational isolation on laser tables, large stable pedestal mounts for the laser transport system

- Gravity gradient noise is second order effect where seismic surface waves (Rayleigh waves) couple to gravitational potential field
- Perturbations in potential field lead to acceleration perturbations
- Atom trajectories are perturbed under free fall and pickup phase shift

$$\delta h_{ggn} = \frac{\pi G \gamma \nu \rho_0}{2LT\omega_{ggn}^3} \langle \delta \xi_z \rangle \left(e^{-\sqrt{2}L\omega_{ggn}/c_R} - 1 \right) \sin(\omega_{ggn} T) \cos(\phi_{ggn}) \quad (9)$$

$$\delta h_{ggn} \sim \left(10^{-19} / \sqrt{\text{Hz}} \right) \left(\frac{100 \text{ m}}{L} \right) \left(\frac{1.5 \text{ s}}{T} \right) \left(\frac{1 \text{ Hz}}{\omega_{ggn}} \right)^3 \left(\frac{\delta \xi_z}{1 \mu\text{m}} \right) \left(\frac{300 \text{ m/s}}{c_R} \right) \quad (10)$$

- Atom interferometer phase response calculated by semi-classical analytical solver using the GGN phase term $\delta\phi_{ggn}$

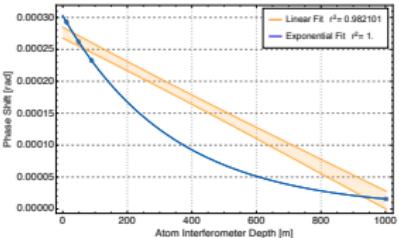
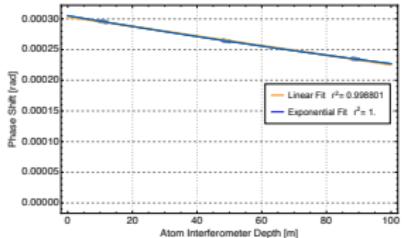
$$\delta\phi_{ggn} = \frac{\pi G \gamma_\nu \rho_0}{\omega_{ggn}} \langle \delta\xi_z \rangle n k_{\text{eff}} T e^{-\sqrt{2}k_I z} \sin(\omega_{ggn} T) \cos(\phi_{ggn}). \quad (11)$$

$$L_c = \frac{1}{2} m \dot{\mathbf{r}}^2 - m\phi(\mathbf{r}) - m\delta\phi_{ggn} \quad (12)$$

Stochastically simulated 1000 measurements and compute the RMS measurement to be fit

Exponential Model

$$\phi(z) = A \exp(-\alpha z)$$

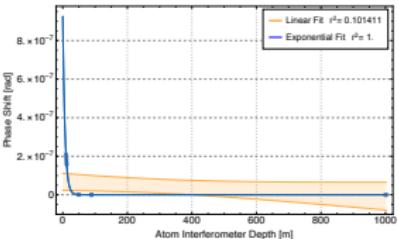
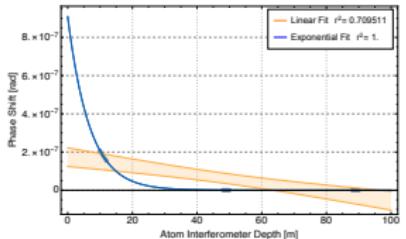
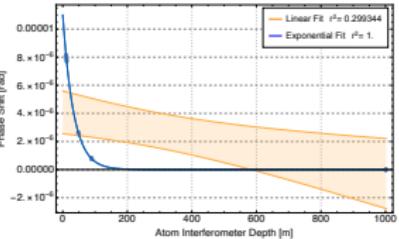
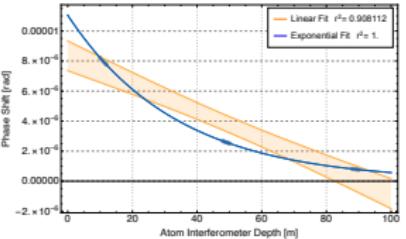


Linear Model

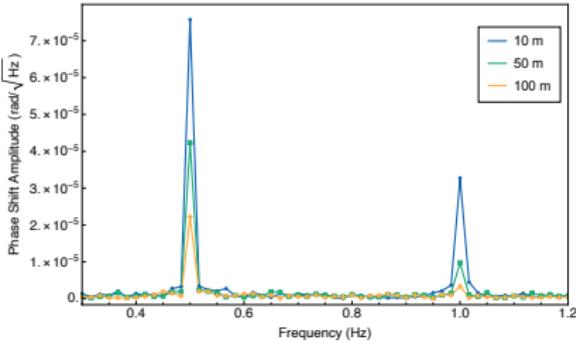
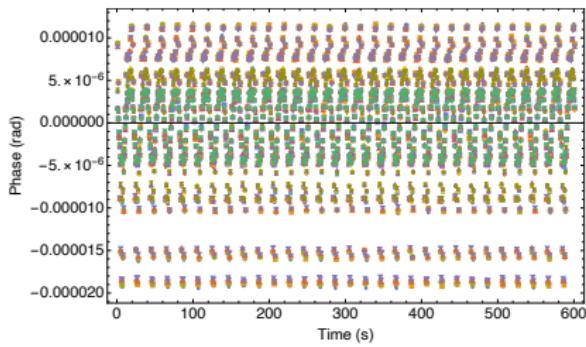
$$\phi(z) = Cz$$

Assumed Rayleigh wave frequencies:

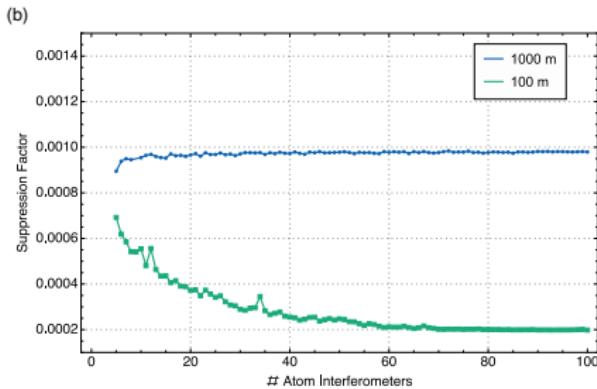
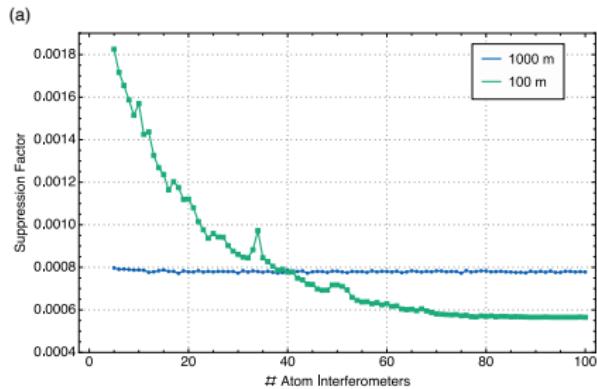
0.1 Hz
1 Hz
5 Hz



- GGN analysis using Fourier components simulated from time sampled phase measurements
- Allows for correlation between atom sources for extraction of exponential character
- No loss of shot-to-shot information by averaging
- Possibly more sensitive than absolute phase measurements



- String-of-pearls characterization
- Residual linear strain noise interferes with GW and DM measurements
- Suppression factor defined by ratio of linear fitting and simulated phase difference along the baseline
- Correlated seismometer array and *in situ* calibration data



- Modeled gravity gradient noise (GGN)
 - Test mass acceleration perturbation caused by density perturbations of ground and atmosphere
 - Important systematic that limits MAGIS-100 maximum sensitivity
- Simulation package developed for MAGIS-100
 - Generates realistic wavefunctions and simulated CCD images
 - Used for R&D and evaluating parameter choice effects on phase extraction precision

Short wavelength approximation:

$$h_{GGN} = \frac{G\rho_0}{2\pi f^2 L} \langle \xi_z \rangle$$

Long wavelength approximation:

$$h_{GGN} = \frac{G\rho_0}{fc_R} \langle \xi_z \rangle$$

Using simple model convert between measured seismic displacements into inferred GGN strain

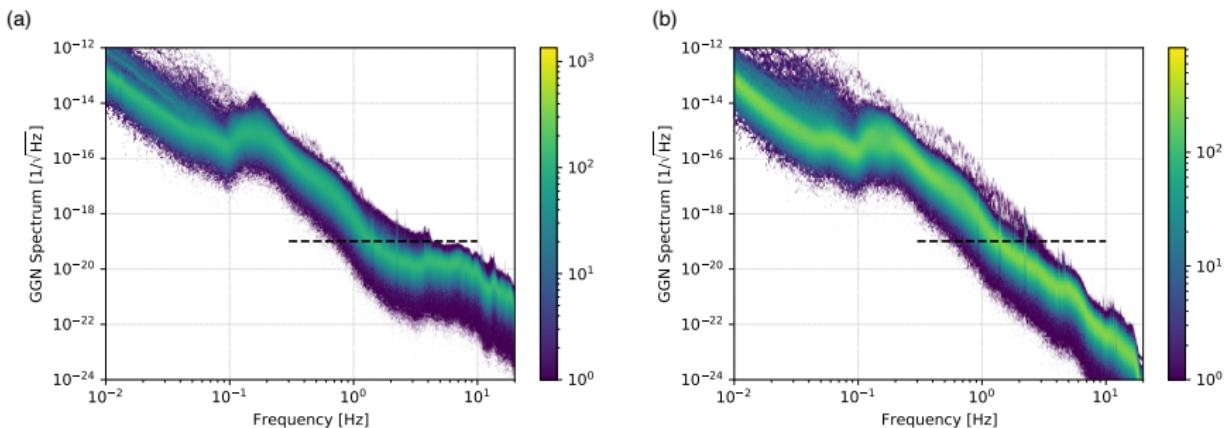


Figure 1: Inferred GGN strain spectrum (a) from surface measurements, (b) from underground measurements. Black dashed line is strain sensitivity after advancements.

- Began as case study into optimizing the optical system parameters
- Evolved into a simulation package for generating images after an atom interferometric measurement

2-port Wavefunction

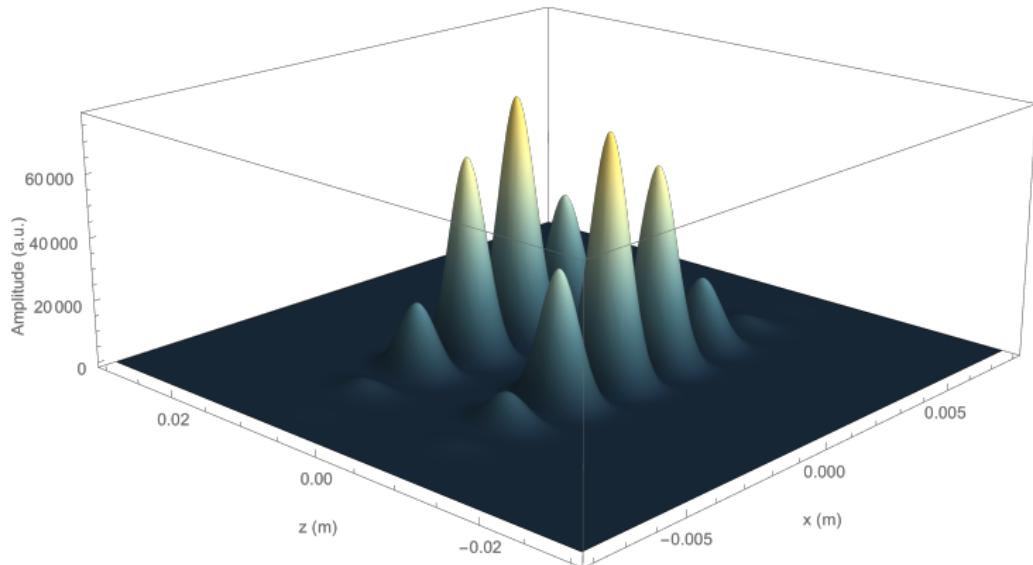


Figure 2: Plot of analytical 2 port wavefunction. Parameters used for strontium atoms with interferometry time $T = 3\text{ s}$, extra drift time $T_{\text{extra}} = 0.15\text{ s}$, cloud radius $w_0 = 0.002\text{ m}$, port separation velocity $\Delta v_z = \hbar k_{z,\text{spread}}/m = 0.10\text{ m/s}$, and a fringe shear $k_{\text{fringe}} = 1/0.0003\text{ m}^{-1}$ with a relative phase $\phi = \pi/2$ between the ports.

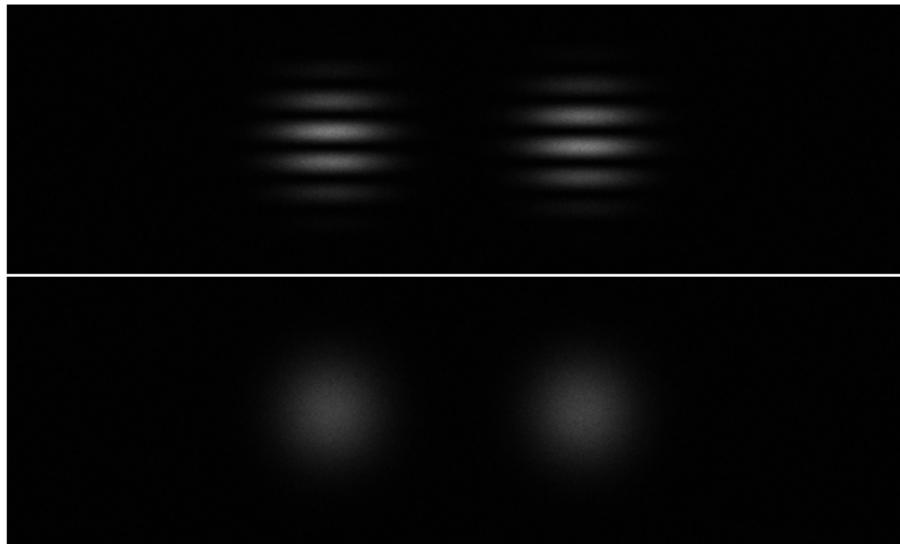


Figure 3: Simulated atom cloud interference pattern pixelated images. Images, going from the top to the bottom, are of the xz projection where z runs horizontally, yz projection with z running horizontally, and xy projection from top down view.

- GGN and environmental noise analysis and site characterization for AION
- Characterization and measurement of atmospheric Newtonian noise
- Modeling of deep underground GGN for kilometer length baselines

- Leading systematics have been studied and analyzed for MAGIS-100
- Mitigation strategies will ensure crucial precision of MAGIS-100 when operation begins
- Helped develop MAGIS simulation and analysis tools that will become part of the MAGIS-100 data analysis pipeline

MAGIS-100 Collaboration

Fermilab Team: Rob Plunkett and others

Stanford: Jason Hogan and his students

Northwestern: Tim Kovachy and his students

Liverpool: Jon Coleman and his students



Northern Illinois
University



JOHNS HOPKINS
UNIVERSITY



UNIVERSITY OF
CAMBRIDGE



UNIVERSITY OF
OXFORD

Support from the Kavli Foundation

Backup Slides

- Gravitational waves will cause acceleration difference between freely falling atoms
- Measure differential acceleration where atoms are proof masses
- Light flight time is mass separation
- Reference is atomic phase

$$ds^2 = dt^2 - (1 + h \sin(\omega(t - z))) dx^2 - (1 - h \sin(\omega(t - z))) dy^2 - dz^2$$

$$L(1 + h \sin(\omega t))$$

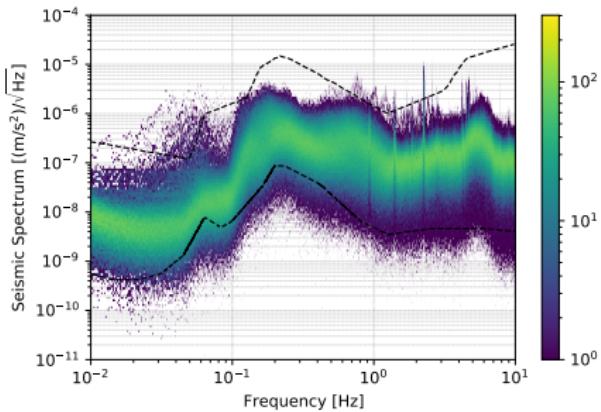
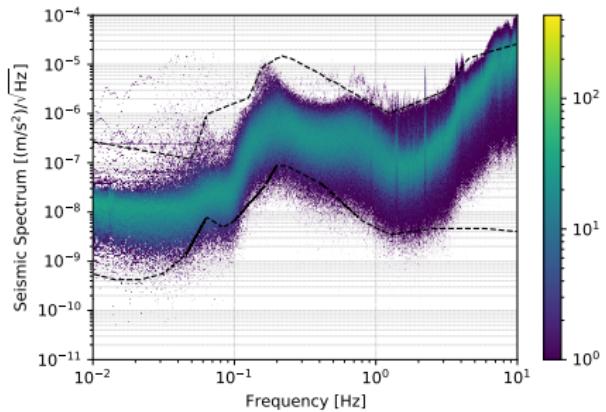
$$\delta\phi_{gw} = 2k_{eff} h L \sin^2\left(\frac{\omega_{gw} T}{2}\right) \sin(\phi_0),$$

- Our analysis of laser wavefront provides the required specs to achieve maximum sensitivity
- More suppression can be gained by averaging effects across the atom cloud during large momentum transfer
- Mitigation strategies include:
 - Tilted laser pulse sequences to image the wavefront across the atom cloud
 - Advanced array of image sensors to image the wavefront from reconstruction

$$\Delta x \sim 0.1 \mu\text{m}$$

$$\delta \sim 3 \times 10^{-4}$$

Seismic Plots



Laser Transport Vibrations

