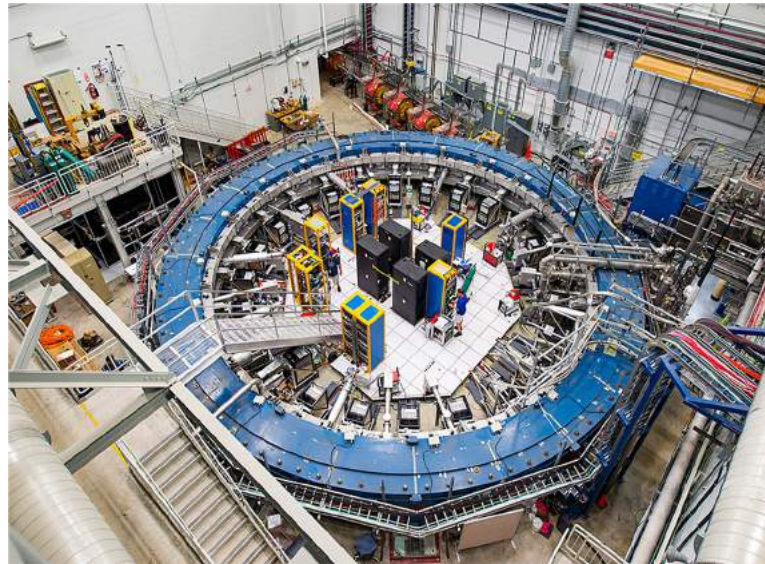


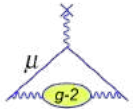
First Results From The Fermilab Muon g-2 Experiment



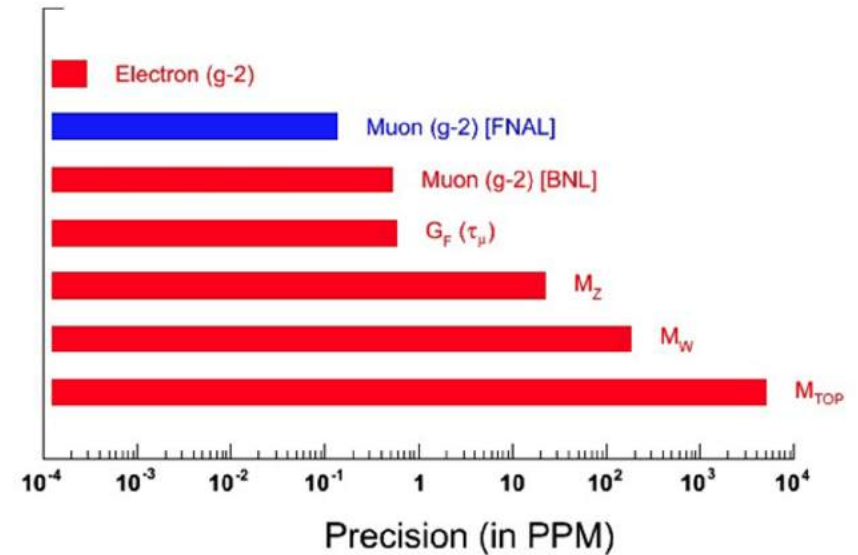
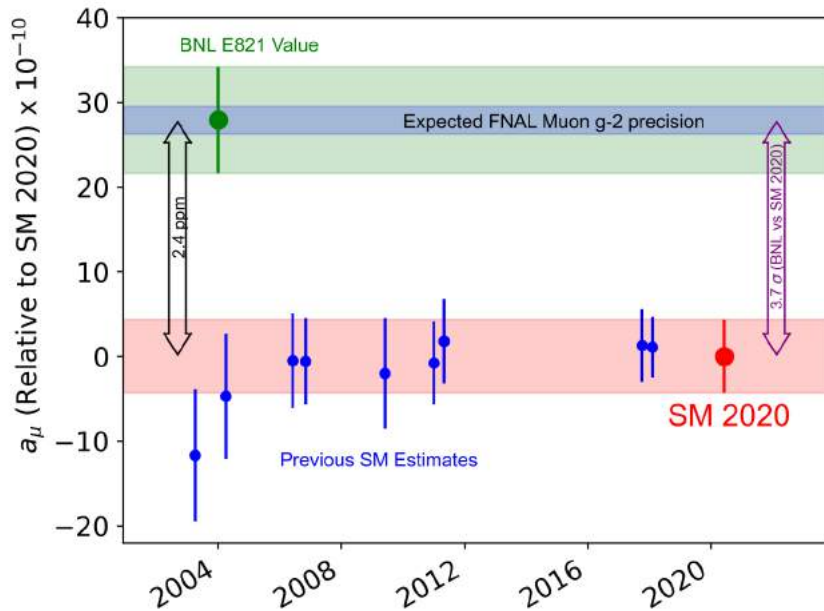
Alex Keshavarzi
Kings College London Seminar
9th June 2021



Precision



The BNL E821 measurement had a 0.54 ppm (540 ppb) uncertainty

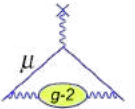


BNL-SM discrepancy: 2.4 ppm

FNAL aim is 100 ppb stat. \oplus 100 ppb syst.

Today's talk is on a dataset of similar size to BNL ~ 10 billion μ^+

Magnetic moments



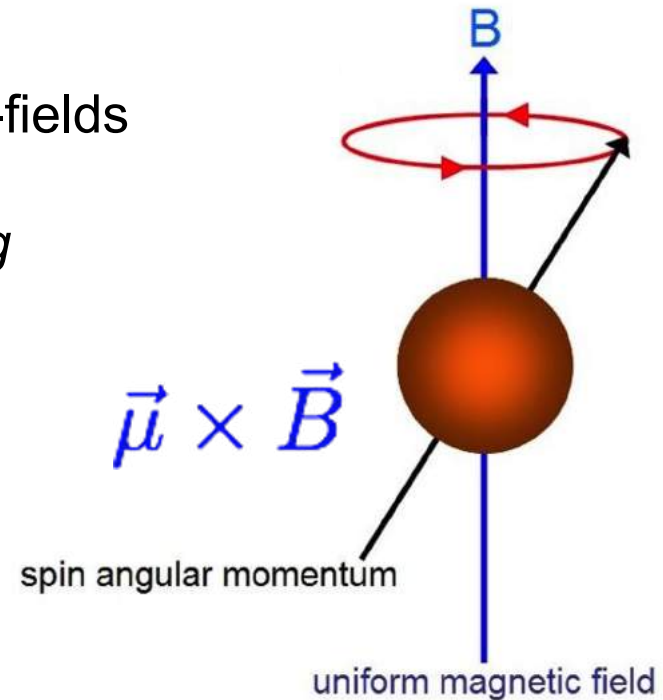
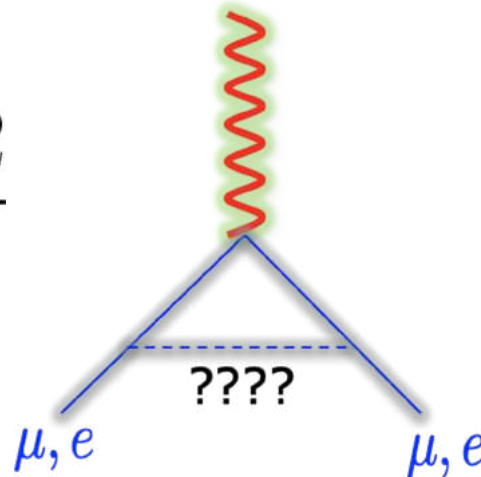
The muon has an intrinsic magnetic moment that is coupled to its spin via the gyromagnetic ratio g :

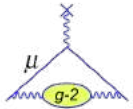
$$\vec{\mu} = g \frac{e}{2m_\mu} \vec{S}$$

Magnetic moment (spin) interacts with external B-fields

Makes spin precess at frequency determined by g

$$a_\mu = \frac{g - 2}{2}$$





Muon $g-2$ Theory

arXiv.org > hep-ph > arXiv:2006.04822

Search...
Help | Advanced

High Energy Physics – Phenomenology

[Submitted on 8 Jun 2020]

The anomalous magnetic moment of the muon in the Standard Model

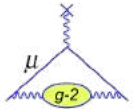
T. Aoyama, N. Asmussen, M. Benayoun, J. Bijnens, T. Blum, M. Bruno, I. Caprini, C. M. Carloni Calame, M. Cè, G. Colangelo, F. Curciarello, H. Czyż, I. Danilkin, M. Davier, C. T. H. Davies, M. Della Morte, S. I. Eidelman, A. X. El-Khadra, A. Gérardin, D. Giusti, M. Golterman, Steven Gottlieb, V. Gülpers, F. Hagelstein, M. Hayakawa, G. Herdoiza, D. W. Hertzog, A. Hoecker, M. Hoferichter, B.-L. Hoid, R. J. Hudspith, F. Ignatov, T. Izubuchi, F. Jegerlehner, L. Jin, A. Keshavarzi, T. Kinoshita, B. Kubis, A. Kupich, A. Kupś, L. Laub, C. Lehner, L. Lellouch, I. Logashenko, B. Malaescu, K. Maltman, M. K. Marinković, P. Masjuan, A. S. Meyer, H. B. Meyer, T. Mibe, K. Miura, S. E. Müller, M. Nio, D. Nomura, A. Nyffeler, V. Pascalutsa, M. Passera, E. Perez del Rio, S. Peris, A. Portelli, M. Procura, C. F. Redmer, B. L. Roberts, P. Sánchez-Puertas, S. Serednyakov, B. Schwartz, S. Simula, D. Stöckinger, H. Stöckinger-Kim, P. Stoffer, T. Teubner, R. Van de Water, M. Vanderhaeghen, G. Venanzoni, G. von Hippel, H. Wittig, Z. Zhang, M. N. Achasov, A. Bashir, N. Cardoso, B. Chakraborty, E.-H. Chao, J. Charles, A. Crivellin, O. Deineka, A. Denig, C. DeTar, C. A. Dominguez, A. E. Dorokhov, V. P. Druzhinin, G. Eichmann, M. Fael, C. S. Fischer, E. Gámiz, Z. Gelzer, J. R. Green, S. Guellati-Khelifa, D. Hatton, N. Hermansson-Truedsson et al. (32 additional authors not shown)

The Muon $g-2$ Theory Initiative



Muon g-2 in the SM

$$\Delta a_\mu = 279(76) \times 10^{-11} \rightarrow 2.39(0.65) \text{ ppm}$$



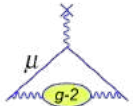
- a_μ arises due to quantum corrections / higher order interactions / loop contributions
- All SM particles contribute \rightarrow Calculate and sum all sectors of the SM:

$$a_\mu^{\text{SM}} = a_\mu^{\text{QED}} + a_\mu^{\text{EW}} + a_\mu^{\text{HVP}} + a_\mu^{\text{HLbL}}$$

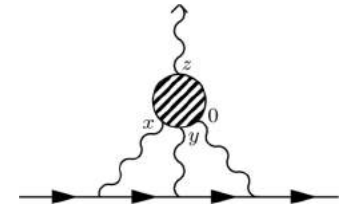
			a_μ^{SM} portion	δa_μ^{SM} portion
QED	<p>1-loop + 2-loop + ...</p>	Perturbative (Known to five-loop)	$\sim 99.99\%$	$\sim 0.001\%$
EW	<p>γ, W, ν_μ, Z, H</p>	Perturbative (Known to two-loop)	$\sim 1 \text{ ppm}$	$\sim 0.2\%$
HVP	<p>had, had</p>	Non-perturbative (Data-driven & lattice)	$\sim 59 \text{ ppm}$	$\sim 84\%$
HLbL	<p>had</p>	Non-perturbative (Data-driven & lattice)	$\sim 1 \text{ ppm}$	$\sim 16\%$

Muon g-2 in the SM: HLbL

$$\Delta a_\mu = 279(76) \times 10^{-11} \rightarrow 2.39(0.65) \text{ ppm}$$



- HLbL scattering - hadronic blob coupled to 3 off-shell/1 on-shell photon.
- Four point function - notoriously difficult to calculate.
- Previously only calculated from models with large systematics.

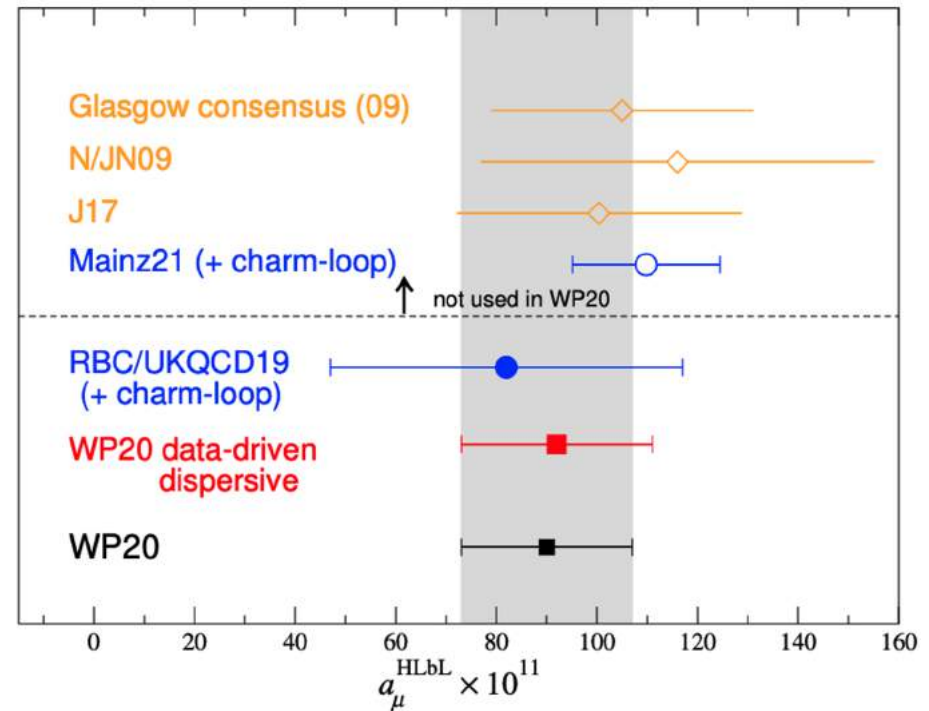


Data-driven (error ~ 0.2 ppm of a_μ^{SM})

- Model-independent dispersive evaluation, using data (e.g. π , η , η' TFFs) as input for hadronic insertions.

Lattice (error ~ 0.3 ppm of a_μ^{SM})

- Model-independent evaluation, computed on discretized Euclidean spacetime (lattice) in finite volume.



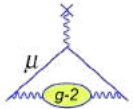
Recommended Muon g-2
TI result (before Mainz):

$$a_\mu^{\text{HLbL}} = 92(18) \times 10^{-11}$$

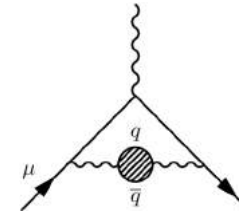
Improved, but still evolving.
Still systematics dominated
(goal < 10% uncertainty)

Muon g-2 in the SM: HVP

$$\Delta a_\mu = 279(76) \times 10^{-11} \rightarrow 2.39(0.65) \text{ ppm}$$



- Hadronic Vacuum Polarisation - hadronic blob coupled to 2 photons.
- Two point function - in principal, much easier than HLbL.
- Most precisely calculated from $e^+e^- \rightarrow$ hadrons cross section data.



Lattice (error ~ 1.6 ppm of a_μ^{SM})

- Uncertainties dominated by finite volume, discretisation and isospin breaking systematics.

Data-driven (error ~ 0.3 ppm of a_μ^{SM})

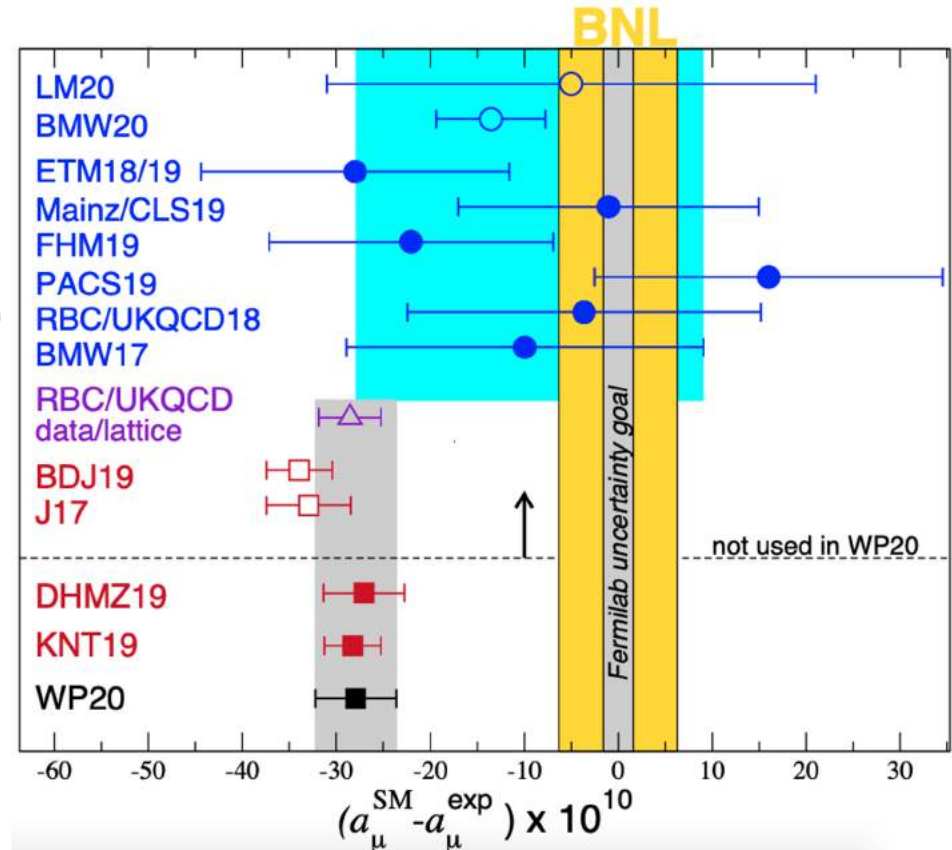
- Cross section data consistently combined and input into dispersion integral:

$$a_\mu^{\text{LOHVP}} = \frac{1}{4\pi^3} \int_{s_{th}}^{\infty} ds K(s) \sigma_{\text{had}}(s)$$

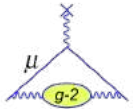
- Several groups have achieved this (most precisely in the UK).

Recommended Muon g-2 TI value from data-driven result:

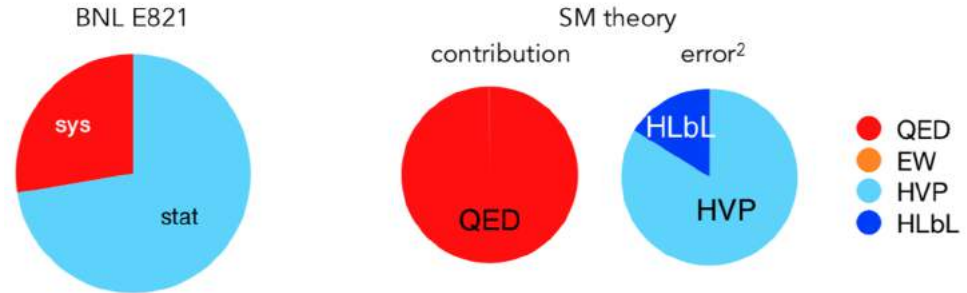
$$a_\mu^{\text{HVP}} = 6845(40) \times 10^{-11}$$



Muon g-2 in the SM and Outlook



Contribution	Value $\times 10^{11}$
Experiment (E821)	116 592 089(63)
HVP LO (e^+e^-)	6931(40)
HVP NLO (e^+e^-)	-98.3(7)
HVP NNLO (e^+e^-)	12.4(1)
HVP LO (lattice, $udsc$)	7116(184)
HLbL (phenomenology)	92(19)
HLbL NLO (phenomenology)	2(1)
HLbL (lattice, uds)	79(35)
HLbL (phenomenology + lattice)	90(17)
QED	116 584 718.931(104)
Electroweak	153.6(1.0)
HVP (e^+e^- , LO + NLO + NNLO)	6845(40)
HLbL (phenomenology + lattice + NLO)	92(18)
Total SM Value	116 591 810(43)
Difference: $\Delta a_\mu := a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	279(76)

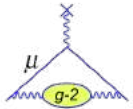


$$\Delta a_\mu = 279(76) \times 10^{-11} \rightarrow 2.39(0.65) \text{ ppm}$$

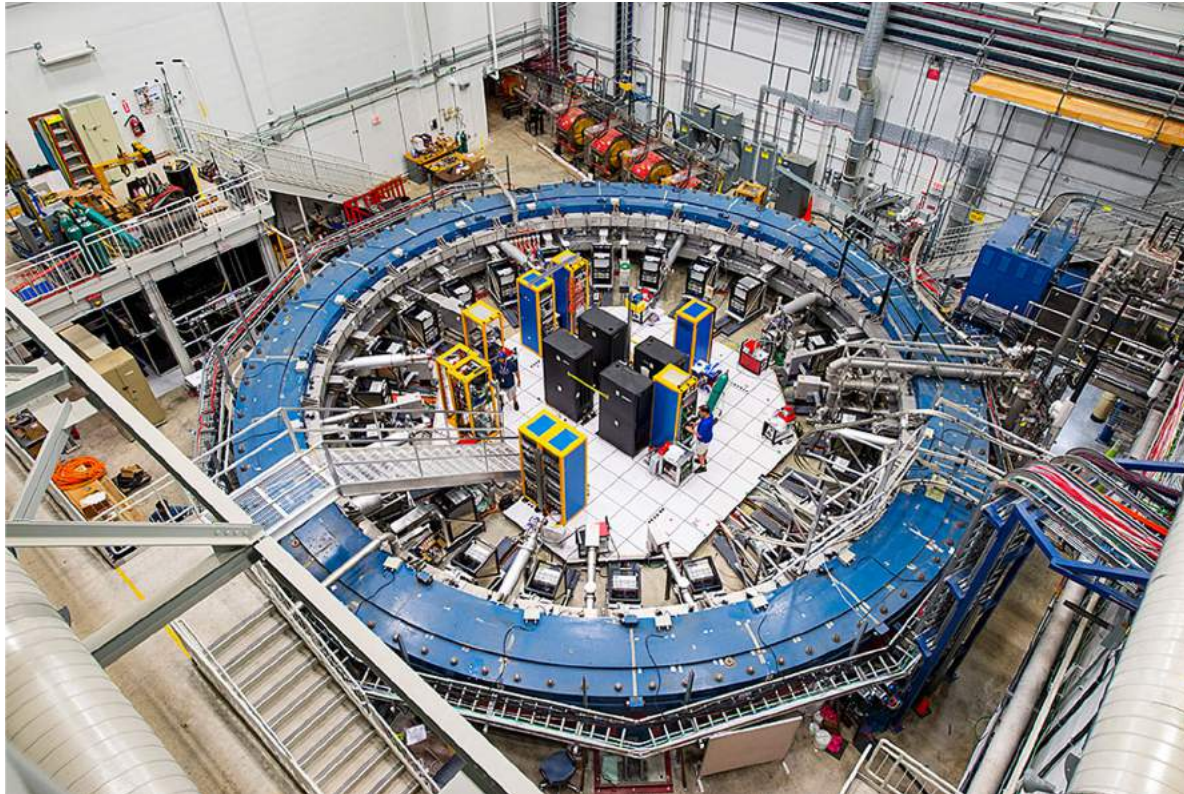
Muon g-2 theory initiative recommended result:

$$a_\mu^{\text{SM}} = 116\,591\,810(43) \times 10^{-11} \text{ (0.37 ppm)}$$

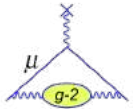
Results in 3.7σ discrepancy when compared to BNL measurement.



The Fermilab Muon $g-2$ Experiment



Muons at Fermilab



Lower instantaneous rate but larger integrated rate than BNL

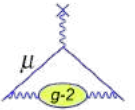


$\sim 10,000\mu^+$ (from 10^{12} p) at 3.1 GeV every 10 ms

(g-2): $\frac{1}{3}$ of proton cycles, neutrino expts: $\frac{2}{3}$

Extra 900m of instrumented beamlines

4 years to build (2 years magnet 'shimming' ...)



May 2013



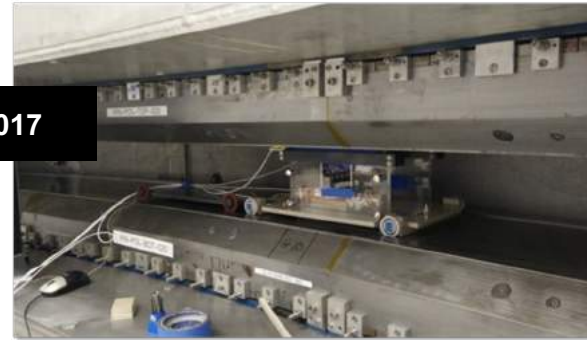
2013



2015



2016-2017

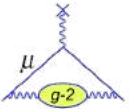


May 31 2017 (g-2)

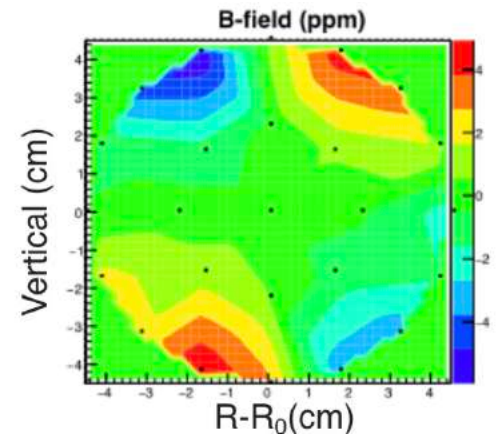
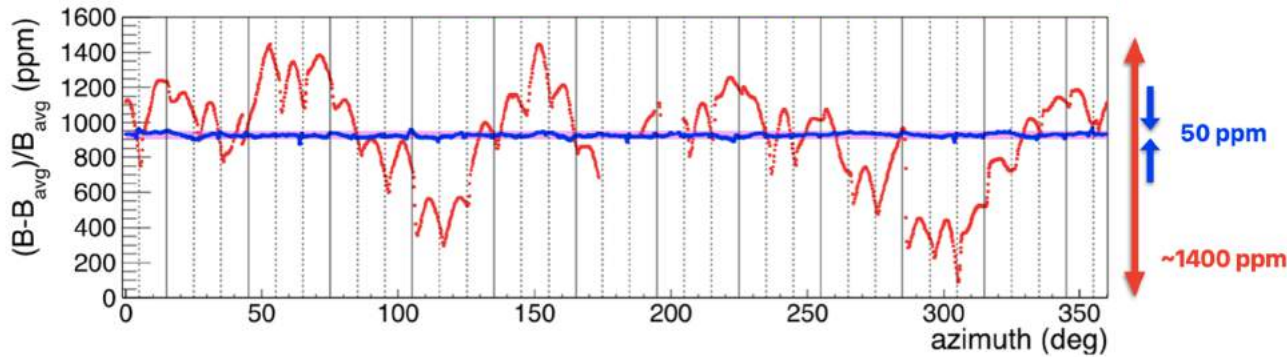


Run-1 data taking started Feb. 2018

The g-2 ring



Magnetic field uniformity 3 times better than the goal (BNL)



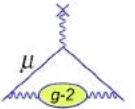
Measurement principle

- Inject polarised muon beam into magnetic storage ring
- Measure difference between spin precession and cyclotron frequencies

$$g = 2, \omega_a = 0$$

- $g \neq 2, \omega_a \propto a_\mu$

$$\omega_a = \omega_s - \omega_c = a_\mu \frac{eB}{mc}$$



Spin precession freq.

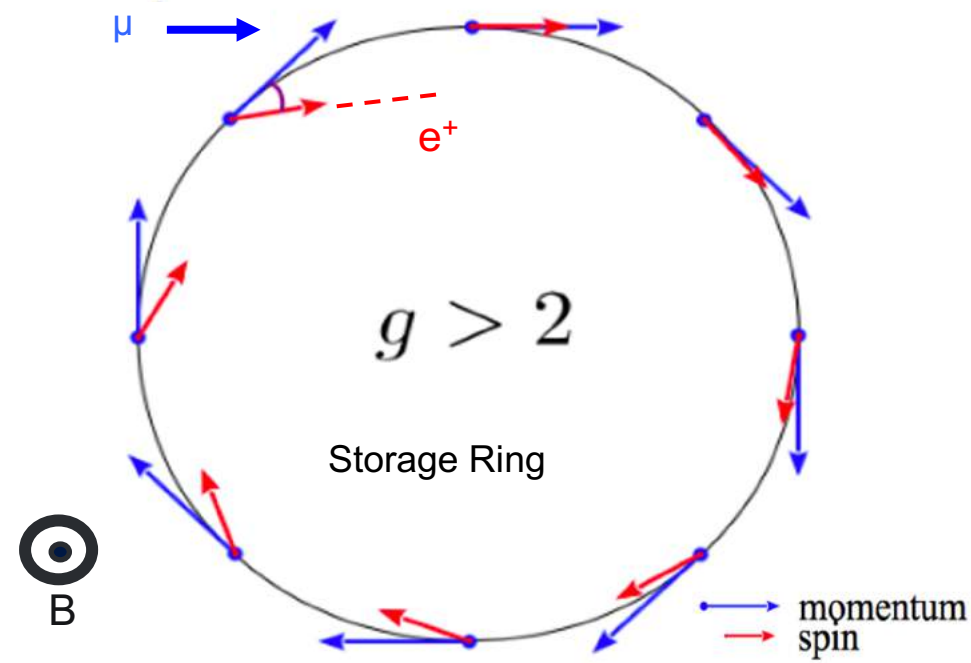
$$\omega_s = \frac{geB}{2mc} + (1 - \gamma) \frac{eB}{\gamma mc}$$

Larmor precession

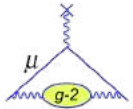
Cyclotron freq.

$$\omega_c = \frac{eB}{\gamma mc}$$

Thomas precession



Measurement details



The experiment actually measures two frequencies

$$a_\mu = \frac{\omega_a}{\tilde{\omega}_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

What we measure

3ppb 0.0003ppb
22ppb

$$\mathcal{R}'_\mu = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

Unblinding conversion factor

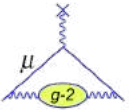
Measured $g - 2$ frequency

Corrections from the beam dynamics systematic effects

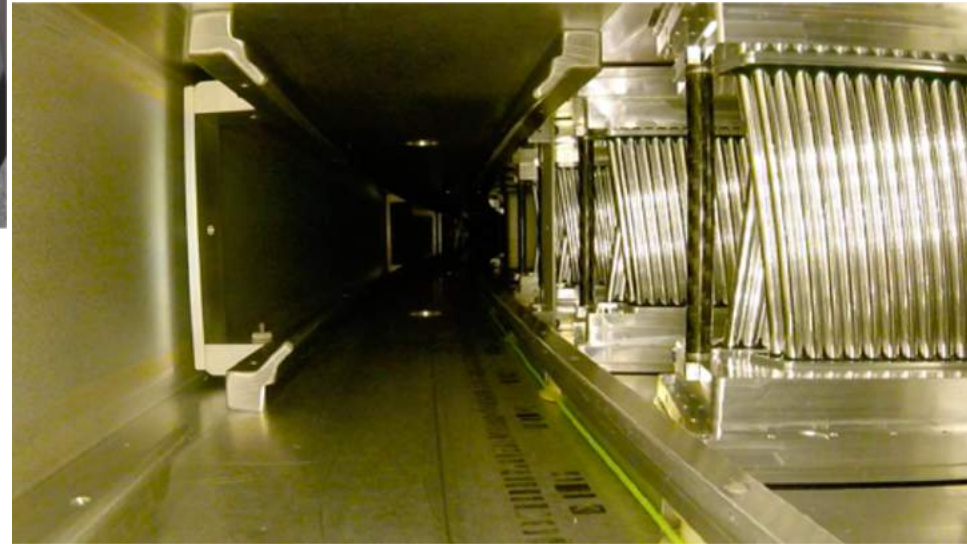
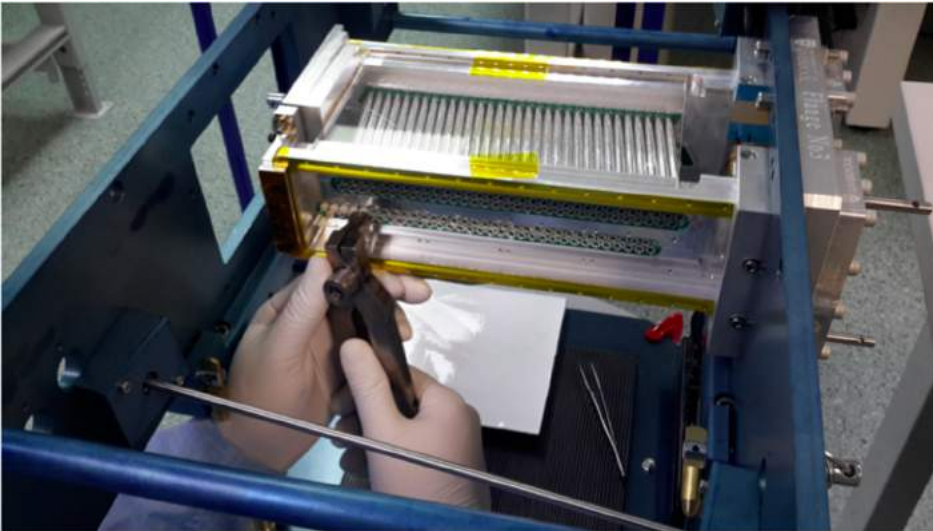
NMR probe calibration factor

Magnetic field weighted over the muon distribution and azimuthally averaged

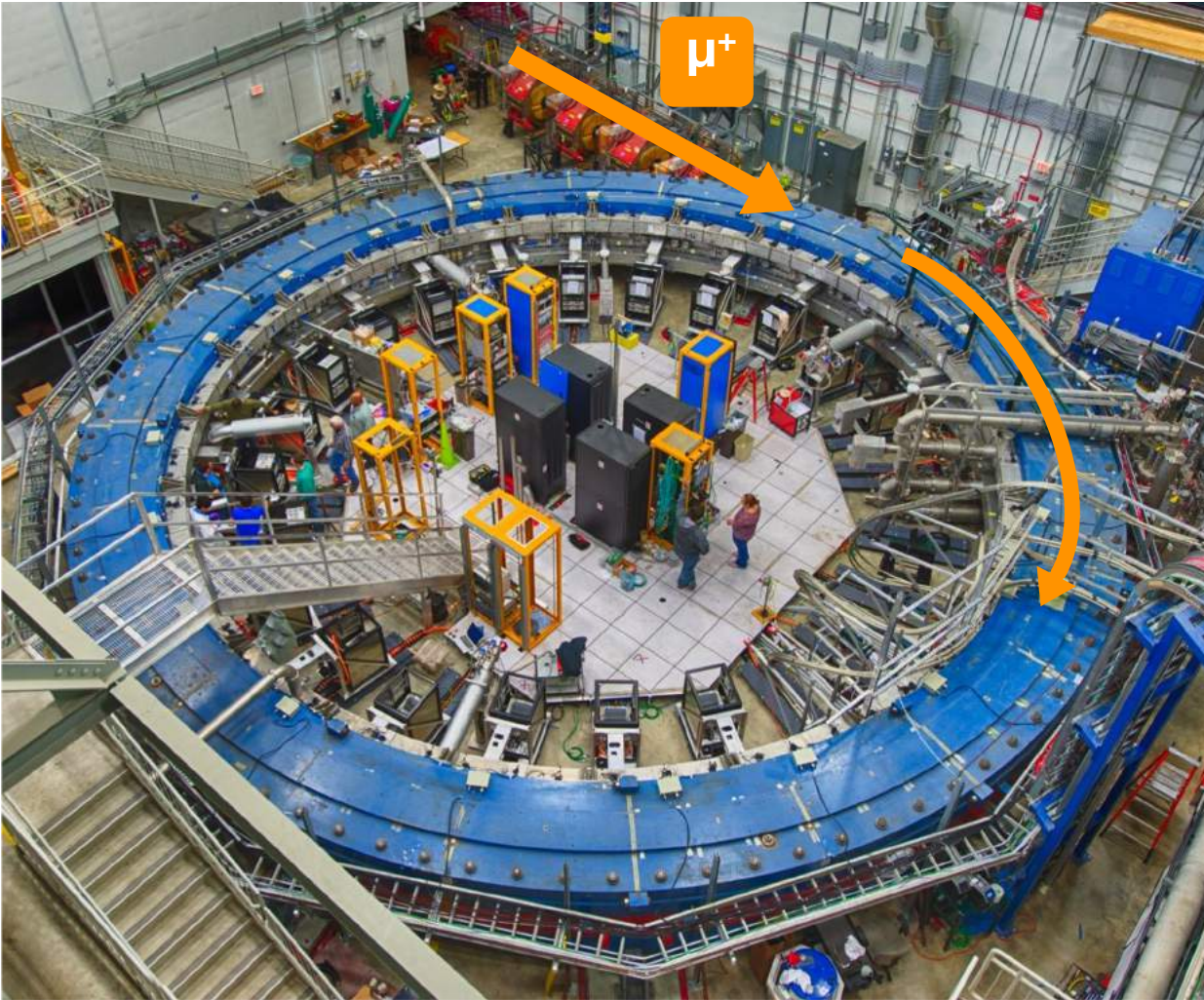
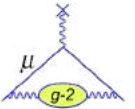
Corrections from the transient magnetic field



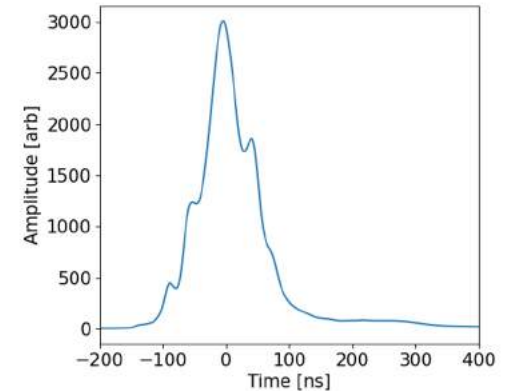
Storing and detecting the beam...



Beam injection

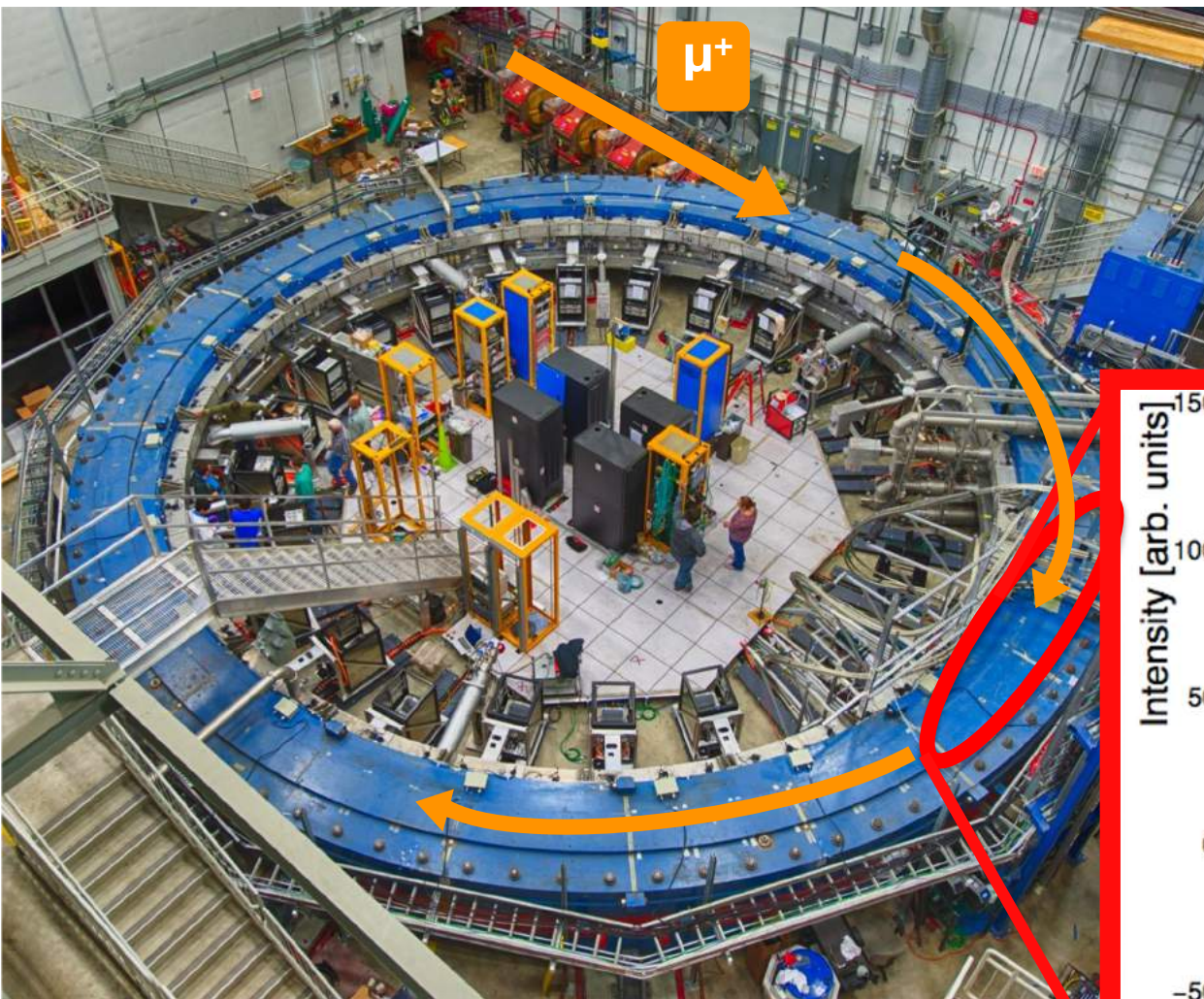
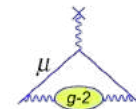


- Monitor beam profile before entrance with scintillating X and Y fibres
- Get time profile of beam using scintillating pad
- $\sim 125\text{ns}$ wide

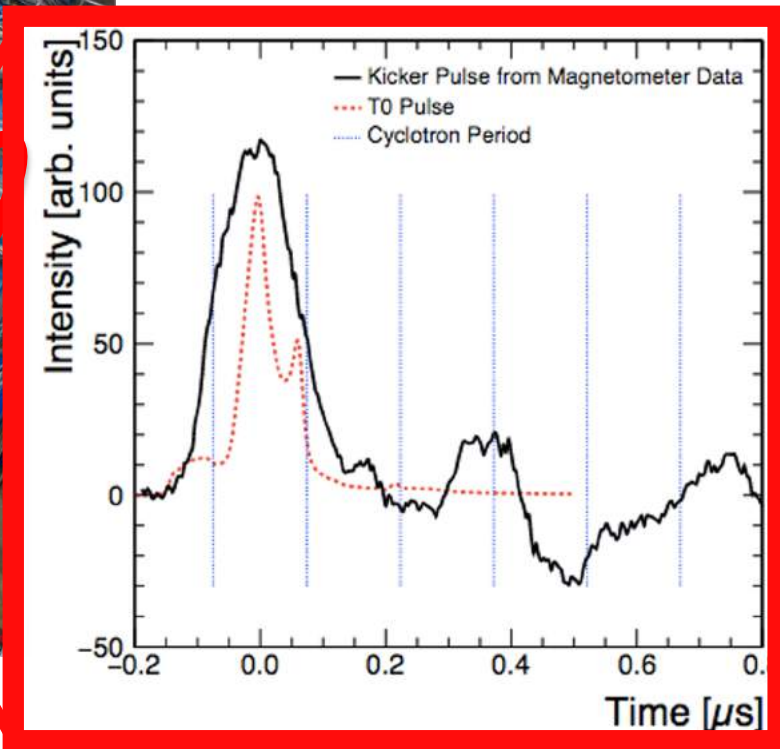


- Cancel B-field during injection using Inflector, so muons can get into the ring

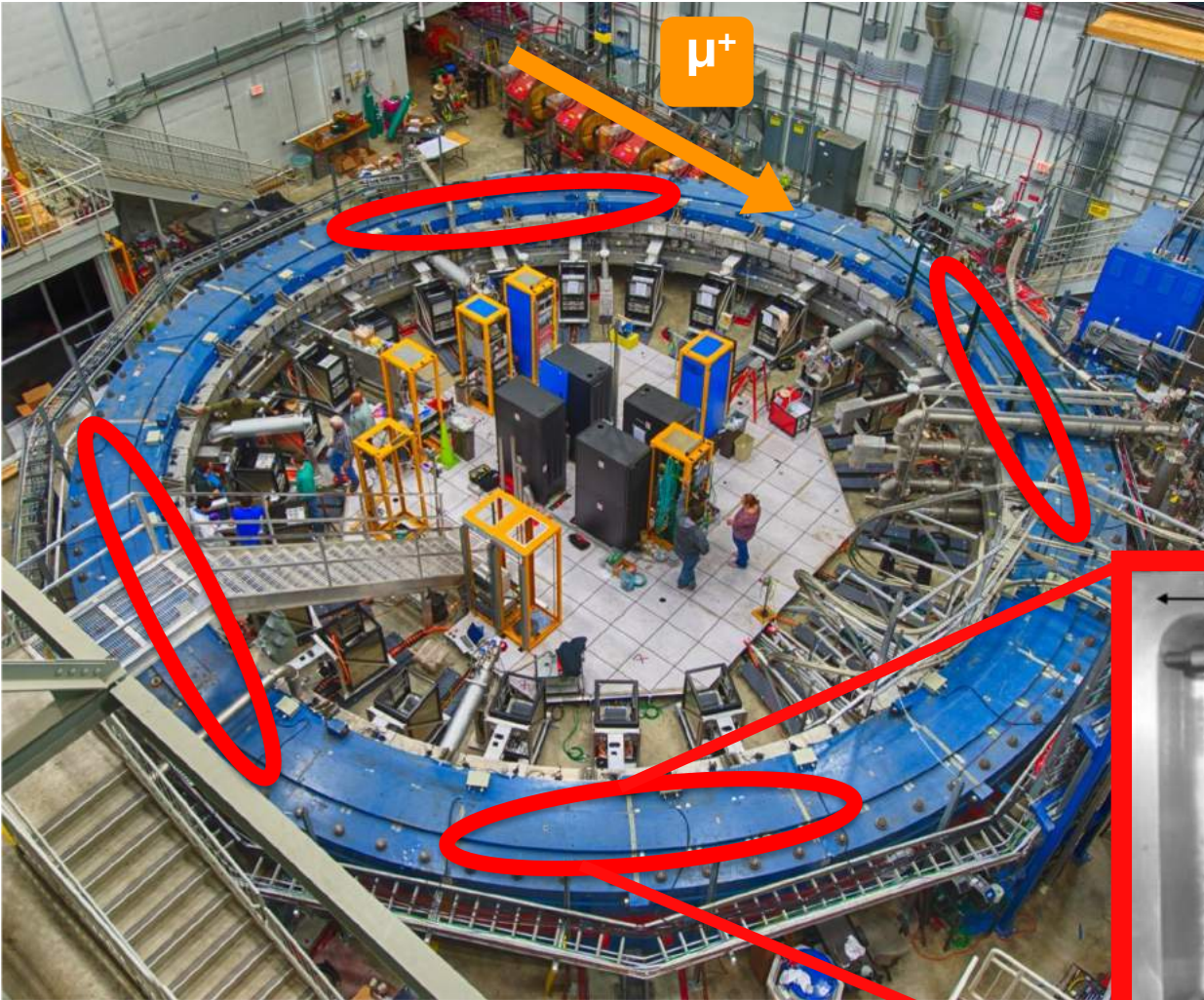
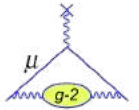
'Kick' onto correct orbit



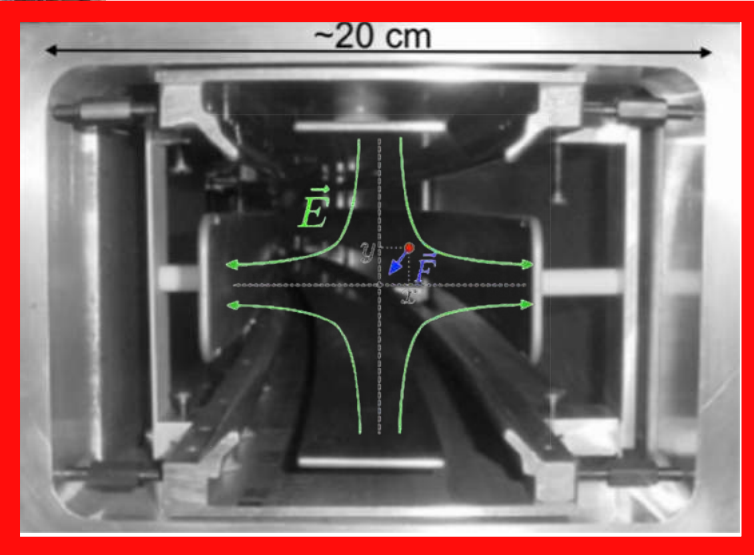
- After Inflector muons are 77mm away from ideal radius
- Apply short magnetic pulse to 'kick' muons onto the correct orbit



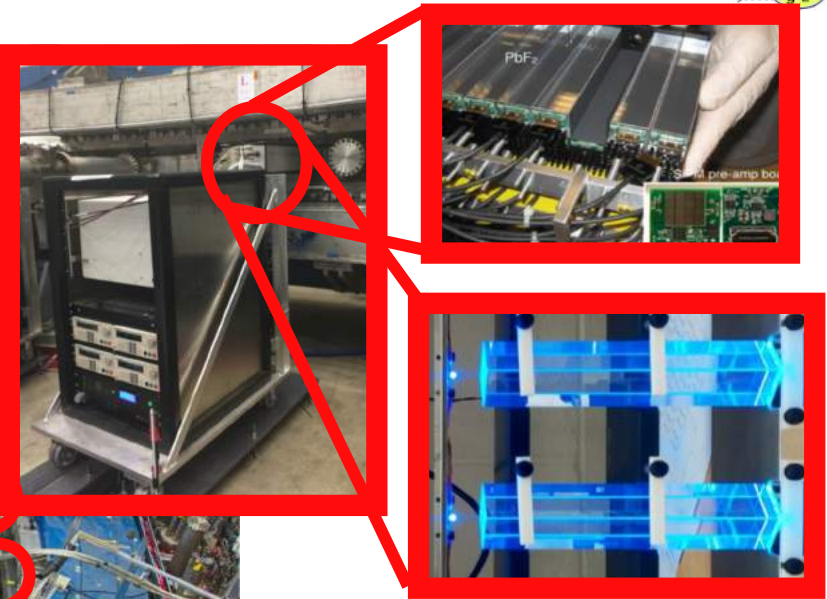
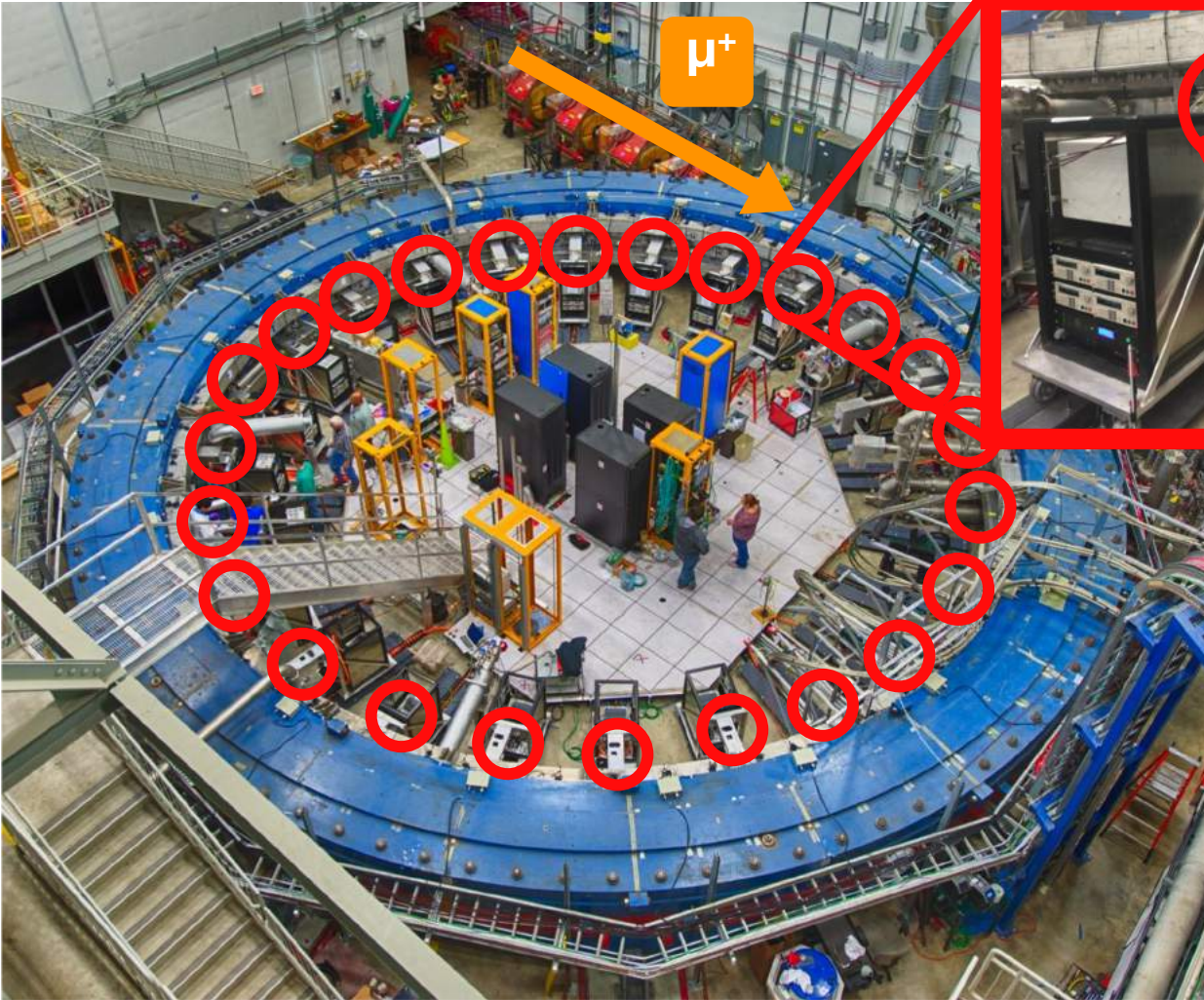
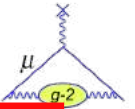
Beam focusing



- Focus the muons vertically
- Aluminium electrodes cover $\sim 43\%$ of total circumference



Calorimeters

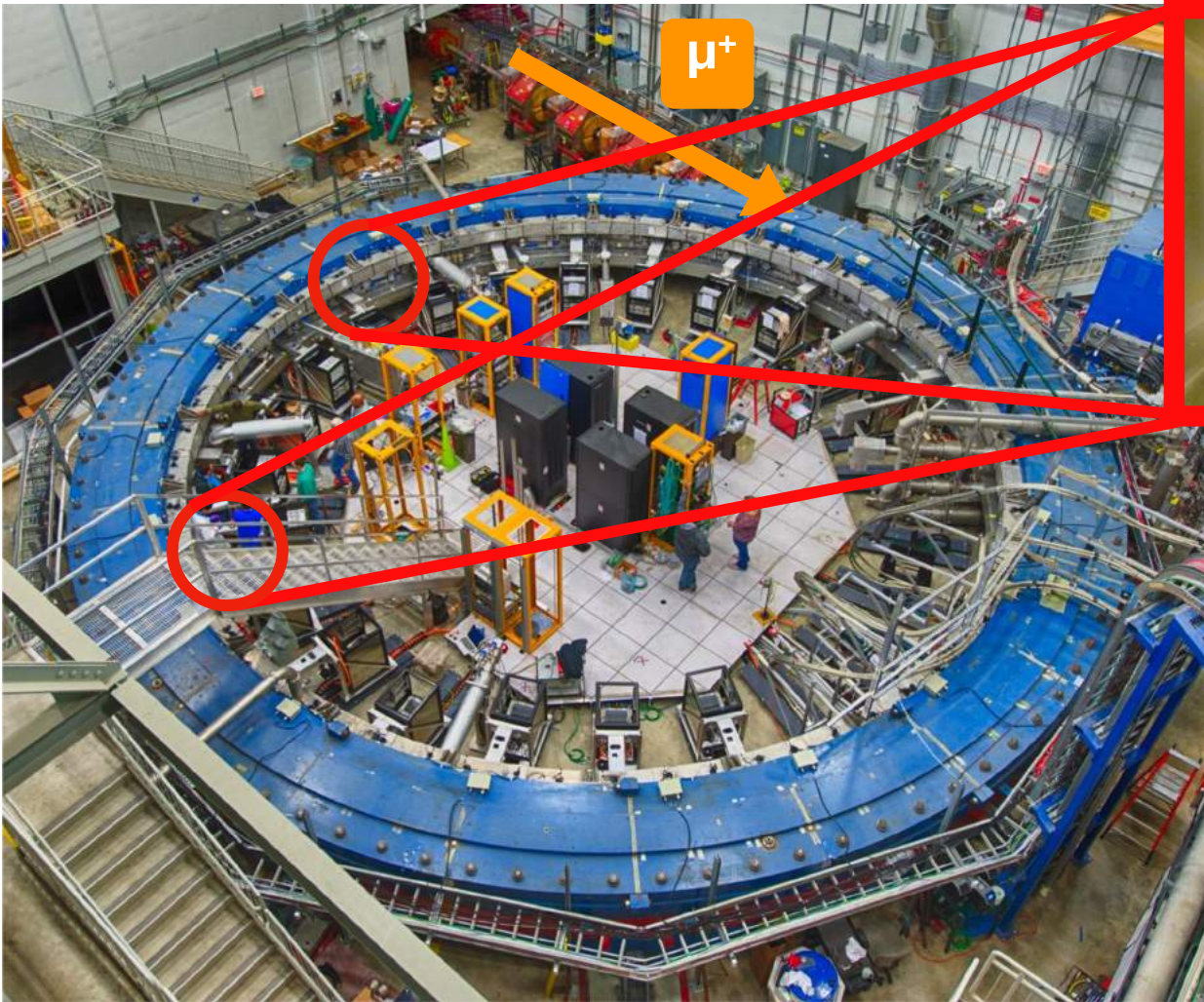
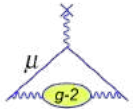


24 Calorimeters

Arrays of 6 x 9 PbF₂ crystals
2.5 x 2.5 cm² x 14 cm (15X₀)

Readout by SiPMs to 800
MHz WFDs

Tracking Detectors



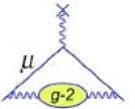
2 Tracking stations

Each contain 8 modules

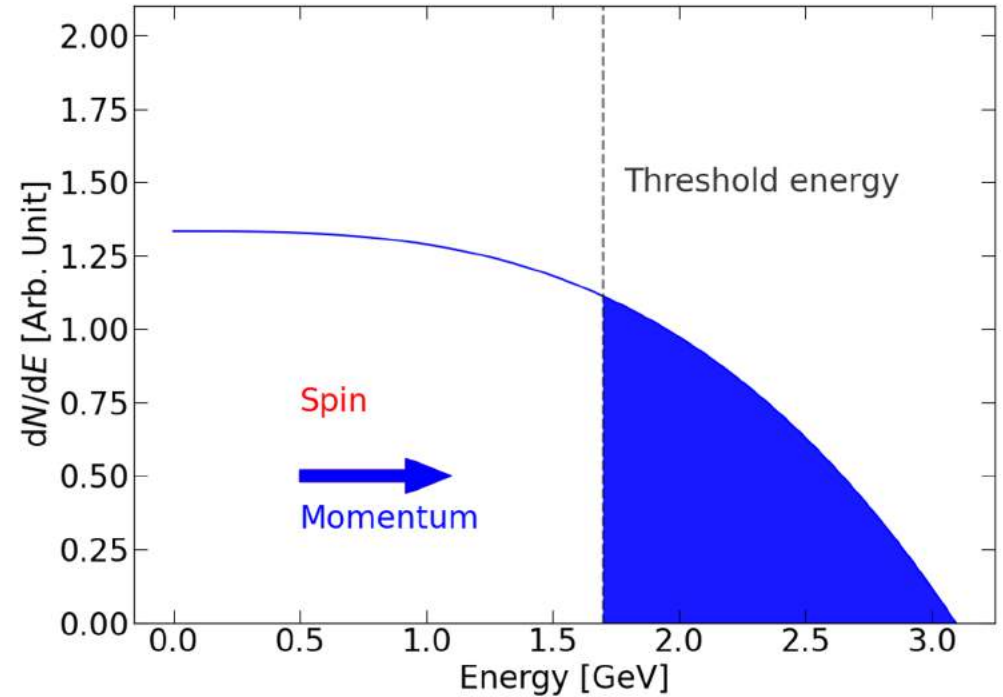
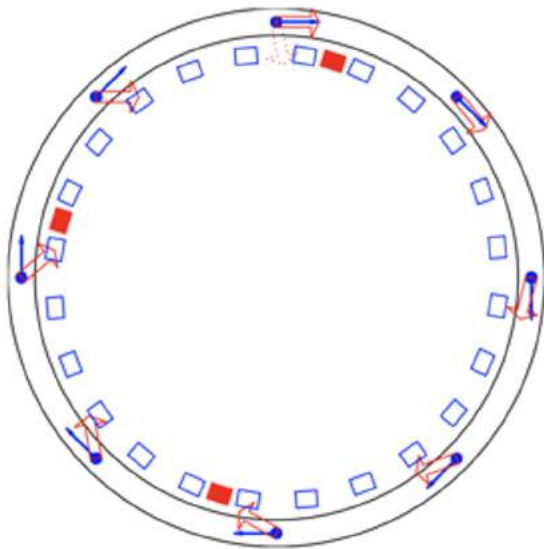
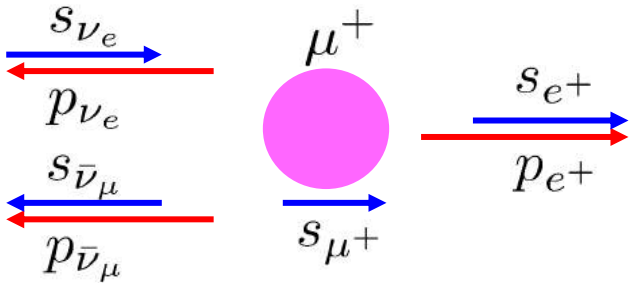
128 gas filled straws in each module

Traceback positrons to their decay point

Measuring ω_a



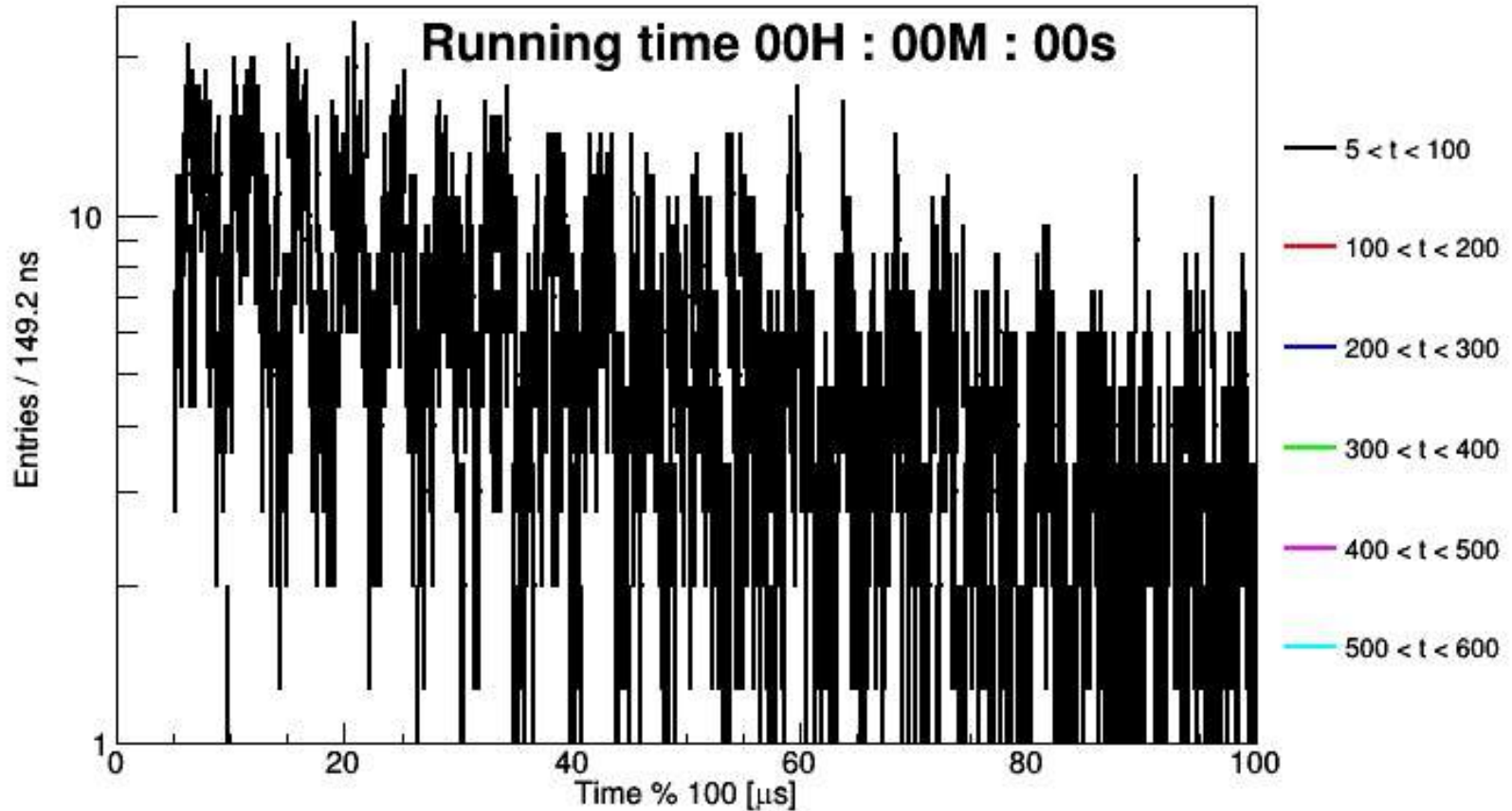
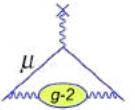
- e^+ preferentially emitted in direction of muon spin



The number of high momentum positrons above a fixed energy threshold oscillates at precession frequency

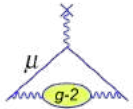
Simply count the number above an energy threshold vs time

Precession in 1 hour of data



$$N_e(t) \simeq N_0 e^{-\frac{t}{\gamma\tau}} [1 - A \cos(\omega_a t + \phi_a)]$$

Beam corrections



- Injected beam has a small vertical component
- Need to use electrostatic quadrupoles to focus the beam vertically

$$\vec{\omega}_a = \frac{e}{mc} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) \left(\vec{\beta} \cdot \vec{B} \right) \vec{\beta} \right]$$

- This introduces 2 additional terms reducing the precession frequency
- **We can minimise the first by choosing $\gamma = 29.3$ to give $p_\mu = 3.1 \text{ GeV}$**
- For a 1.45T field, this sets the radius of the ring to 7.11m
- However we now have 2 corrections to make to a_μ because:

Not all muons are at the 'magic' momentum of 3.1 GeV

E-field correction

$$C_E = \frac{\Delta\omega_a}{\omega_a}$$

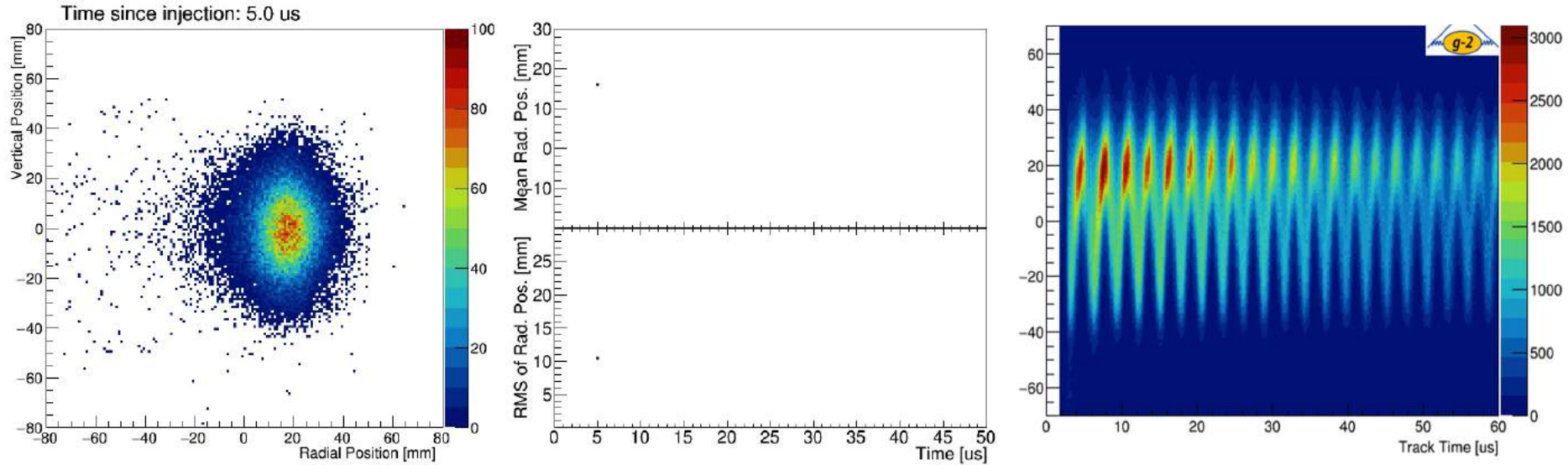
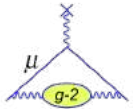
Vertical momentum component aligned with B field

Pitch correction

$$C_P = \frac{\Delta\omega_a}{\omega_a}$$

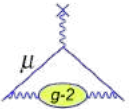
- Both corrections depend on the quadrupole field strength, and are $< 0.5 \text{ ppm}$

Beam Measurements

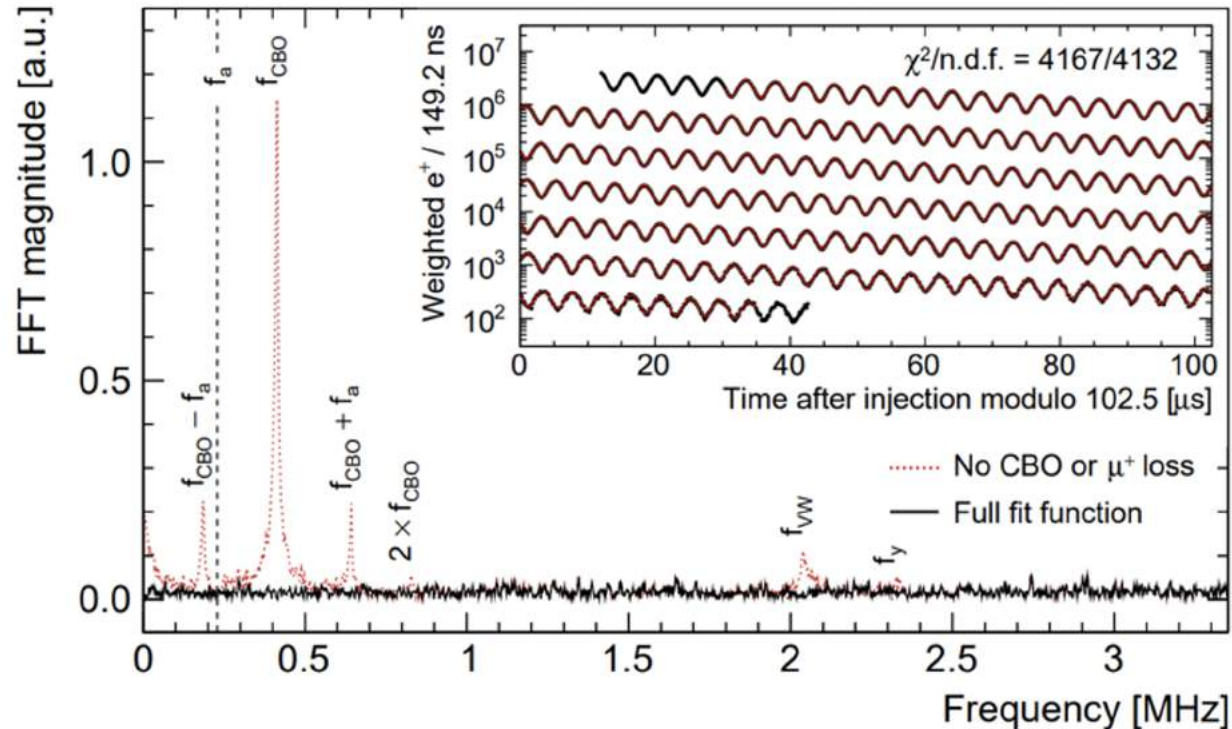


- Use the tracking detectors to measure the decay positrons to infer the decay position
- Muons oscillate radially and vertically at different frequencies, according to the quadrupole strength

Fitting for ω_a

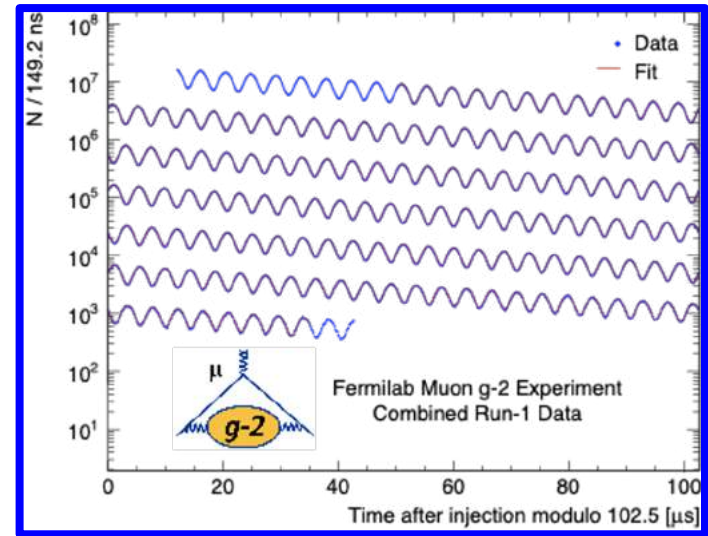
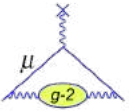


- A fourier transform of the residuals to the fit shows contributions from the movements of the beam, pileup and muon losses

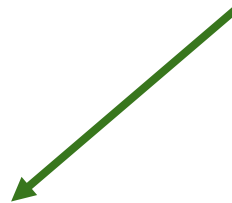


- To account for these effects additional terms are included in the final 24 parameter fit function

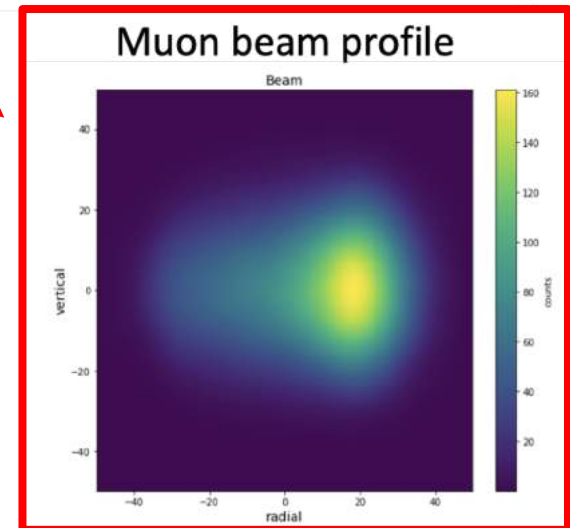
Field measurement

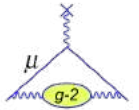


$$\mathcal{R}'_{\mu} = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

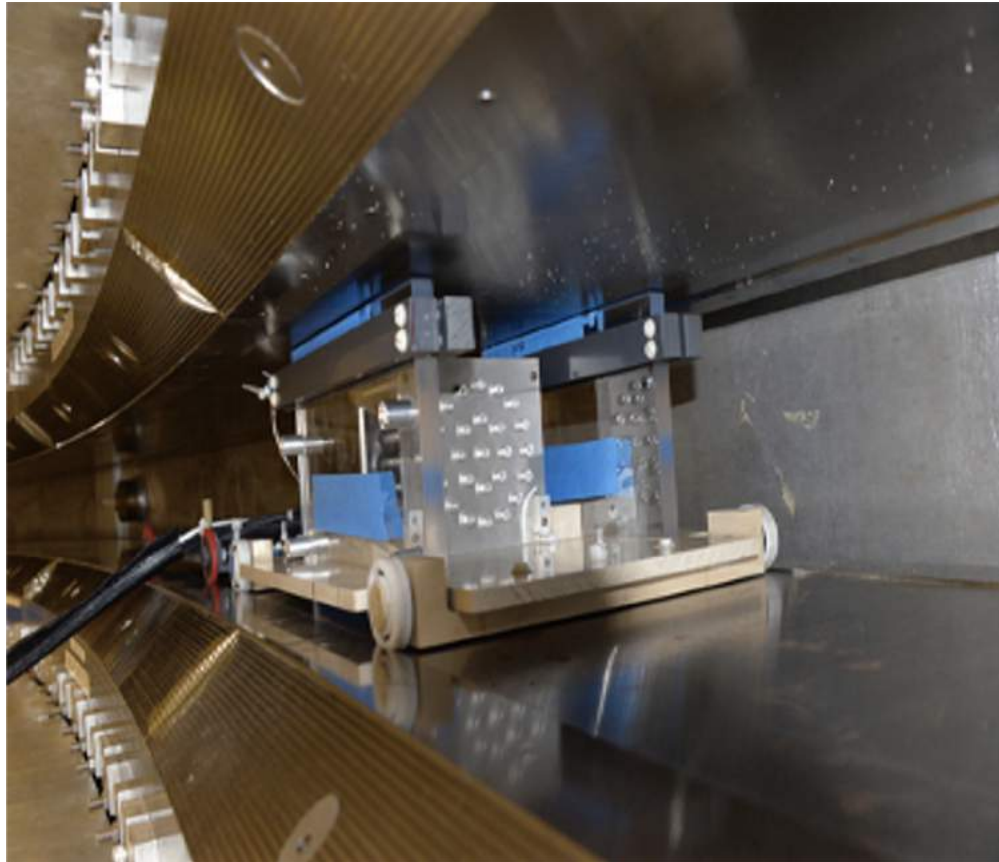


Measuring the magnetic field is the last piece



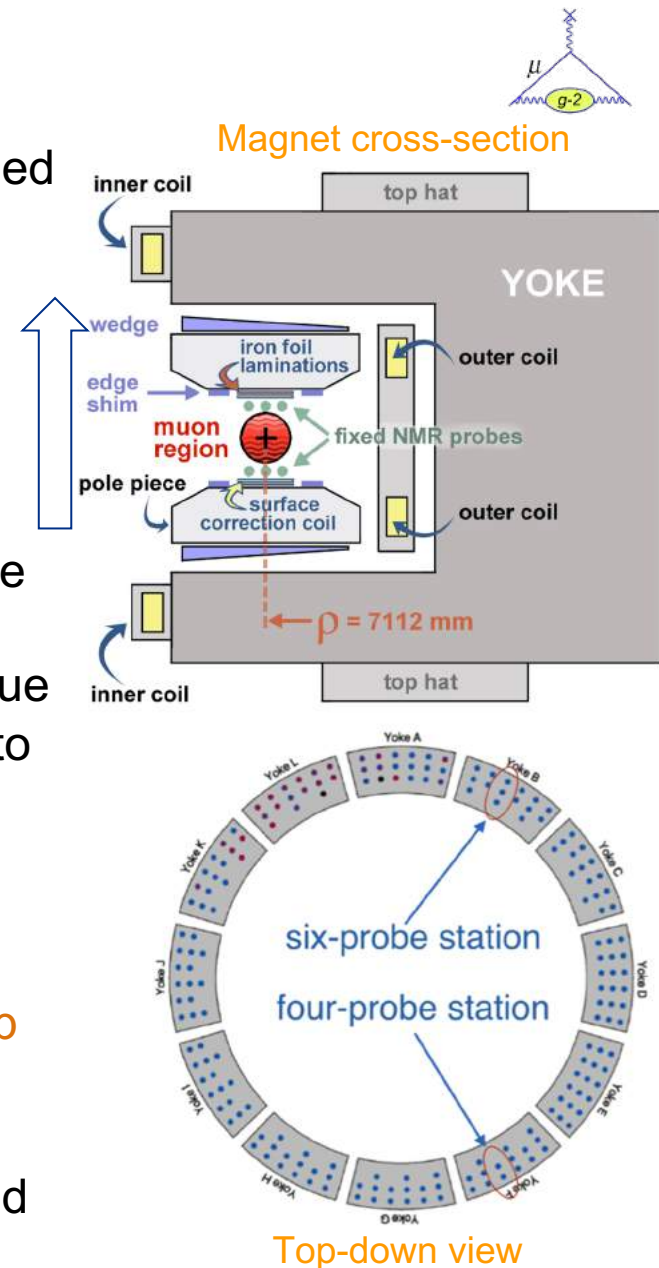


Magnetic field measurement

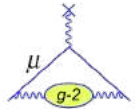


The g-2 storage ring magnet

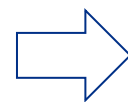
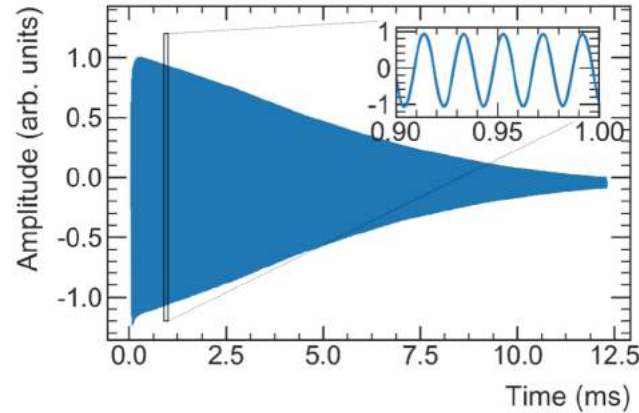
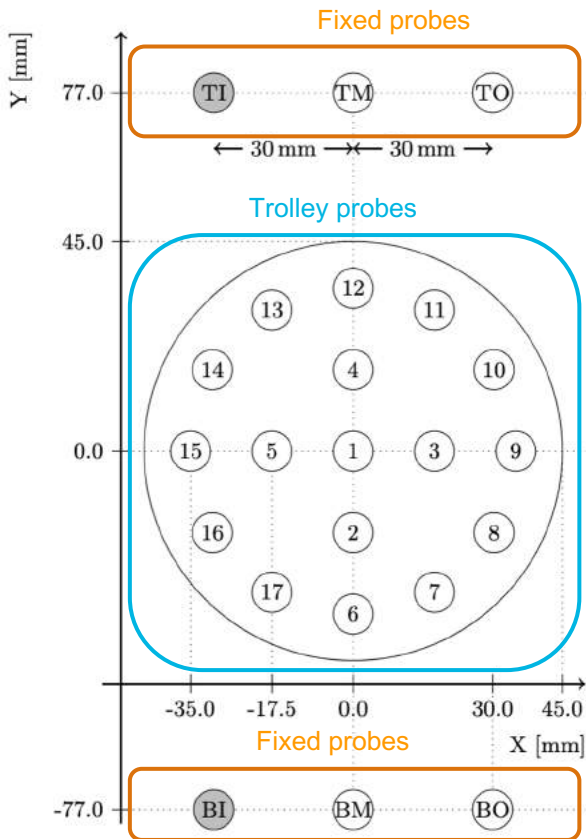
- 7.112 m radius 'C'-shape magnet with vertically-aligned field $B = 1.45$ T
- Dipole field has ppm-level uniformity (14 ppm RMS across the full azimuth)
- Tiny (ppm) changes in magnet geometry, driven by temperature changes, cause the field to drift over time
- Measured using pulsed NMR – a well-known technique that is routinely used in a wide range of applications to measure magnetic fields at the ppb level
- 378 'fixed' NMR probes, built for this experiment, around the ring measure the drift continuously, and provide feedback to the magnet power supply to keep the dipole (vertical) term constant
- Shimming devices minimise gradients (transverse and azimuthal field components).



Measuring the field: the NMR Trolley



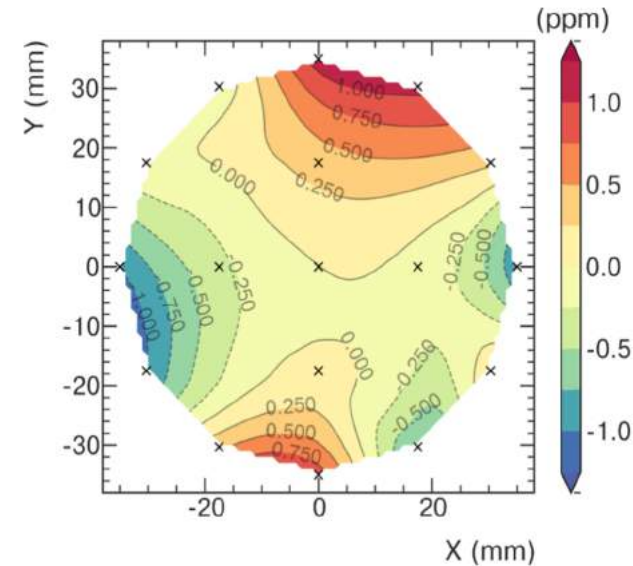
- An in-vacuum trolley with 17 NMR probes drives around the ring every ~ 3 days, mapping out the field components



Field measured by extracting frequency from a Free Induction Decay (FID) spectrum

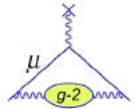


Field gradients in an azimuthal "slice"



At ~ 8000 azimuthal locations, obtain a field contour plot from the 17 probes

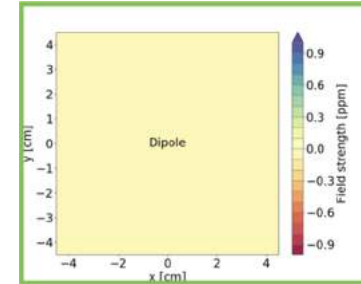
Spatial dependence of B



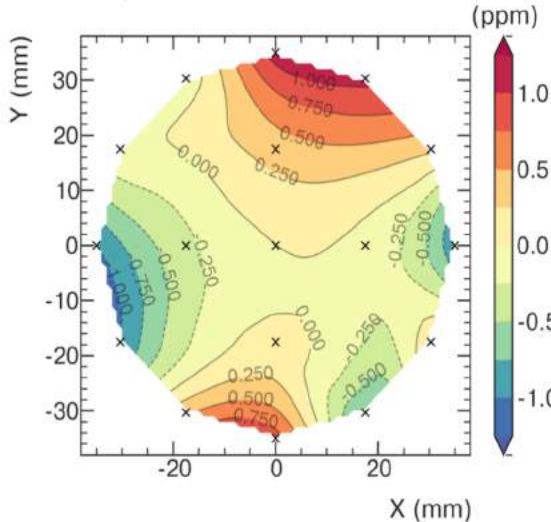
- Extract terms from a multipole (m) expansion of B in r and θ :

$$B \approx B_y = A_0 + \sum_{n=1} \left(\frac{r}{r_0} \right)^n (A_n \cos(n\theta) + B_n \sin(n\theta))$$

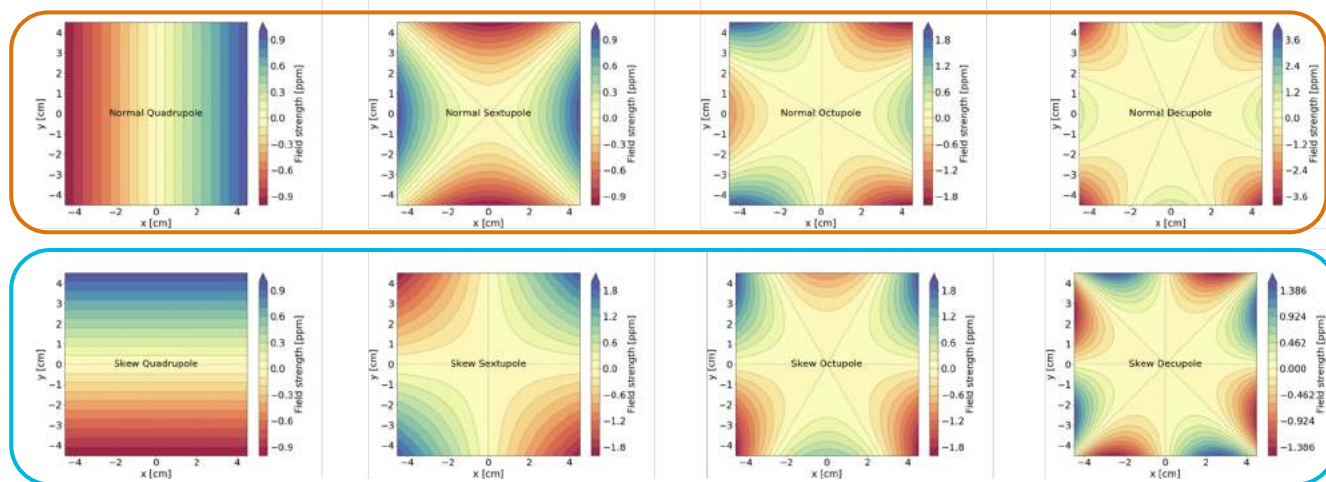
Dipole (m1)



Field gradients in an azimuthal “slice”



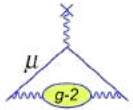
“Normal” terms: m2, m4, m6, m8, ...



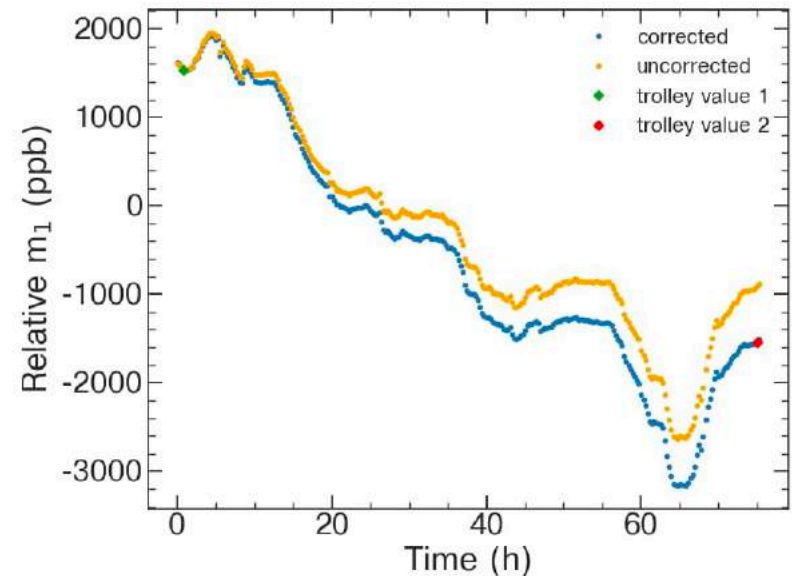
“Skew” terms: m3, m5, m7, m9, ...

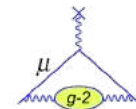
- Trolley: Fit the 2D contour plot to extract the multipole terms (m1, m2, m3, ...)
- Fixed probes: extract terms from geometric combination of probe frequencies
- Fixed probes can track m1, m2, m3, m4 only

Interpolating between trolley runs



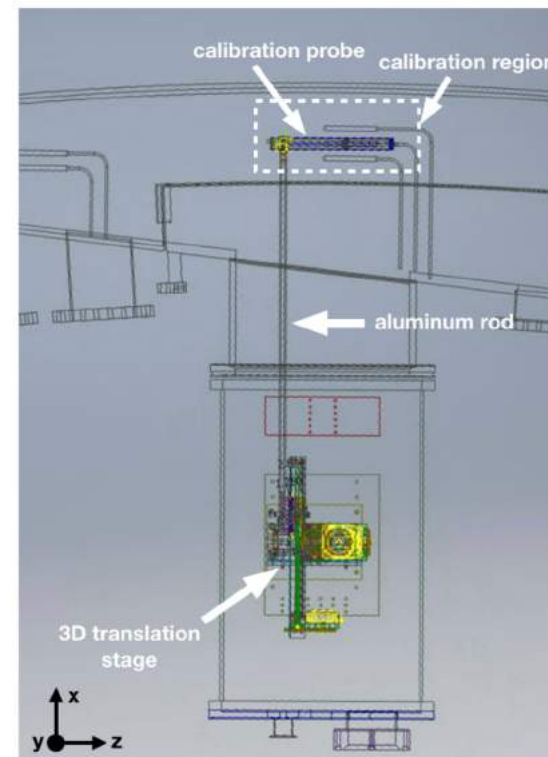
- Need to know the field experienced by the muons, but the trolley cannot take data when the muons are present. **One trolley run takes 3 hours, every ~3 days.**
- Fixed probes take data continuously during muon fills. Use this data to **interpolate** between trolley runs.
- There are 72 fixed probe 'stations' around the ring, every ~5 degrees
- The fixed probe measurements are calibrated using the trolley measurements both times the trolley passes
- Calibration drifts over time, due to changes in higher-order terms that cannot be tracked by the fixed probes
- Leads to the **tracking error uncertainty** (22 - 43 ppb in the run 1 datasets)





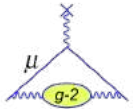
Absolute calibration

- Trolley and fixed NMR probes use petroleum jelly as the proton sample. Chosen for low volatility.
- Must calibrate with protons in a water sample (measurement standard) in order to measure a_μ
- A dedicated calibration probe with a cylindrical H_2O sample is installed inside the vacuum chamber.
- In a dedicated calibration campaign, trolley and calibration probes switch places to repeatedly measure the same field in the same place
- Calibration probe is calibrated against a different probe with a spherical water sample.
- Both calibration probes were cross-checked with a spherical 3He sample (different systematics)

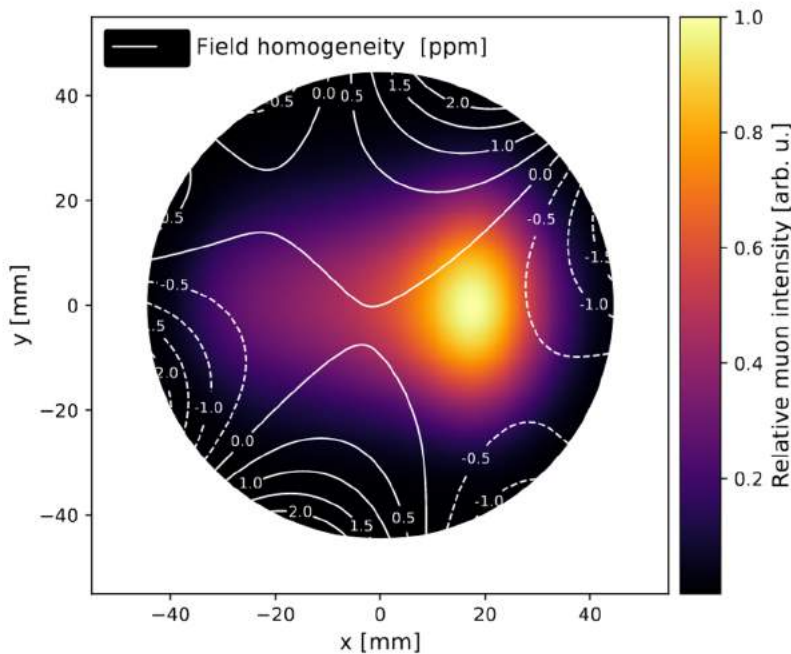


Agreement between all three calibration probes at 10 ppb level

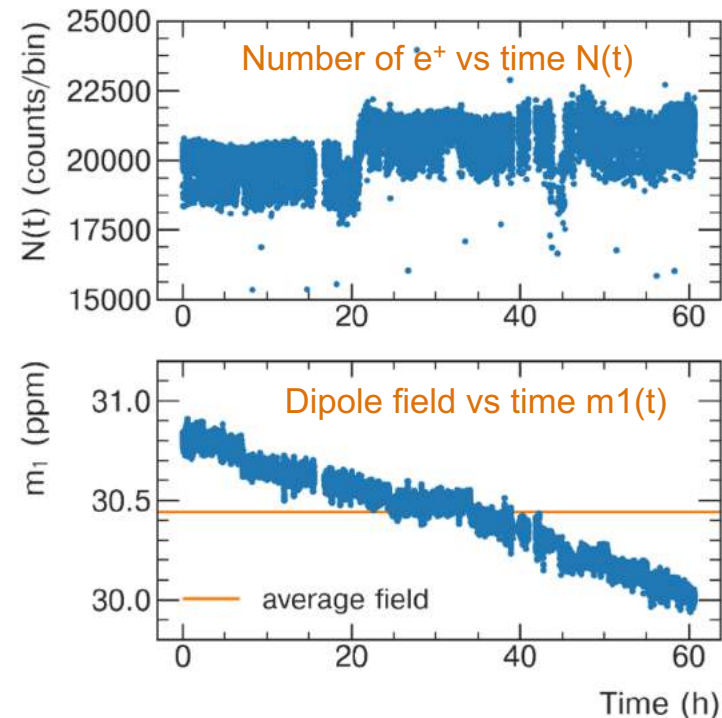
The muon-weighted field



- To obtain the field experience by the muons, the magnetic field distribution as a function of time must be weighted by:
 - The number of muons as a function of time, $N(t)$
 - The beam distribution as a function of time

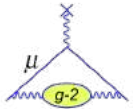


The field is weighted by the 2D beam distribution. An average beam distribution for every 3 hours is used.

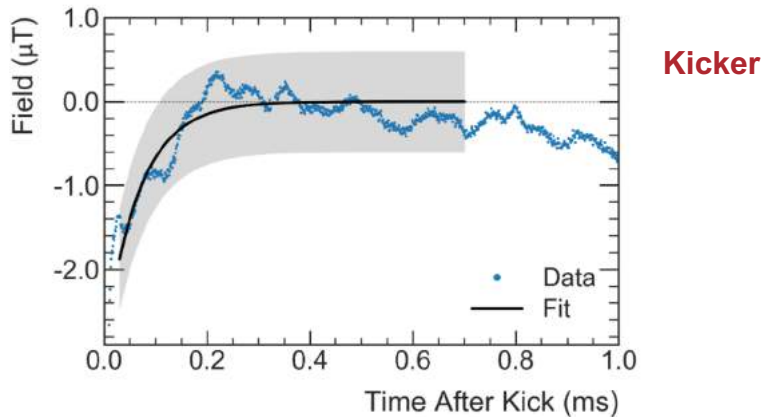


Measured field (every 1.7 s) is weighted by the number of detected e^+

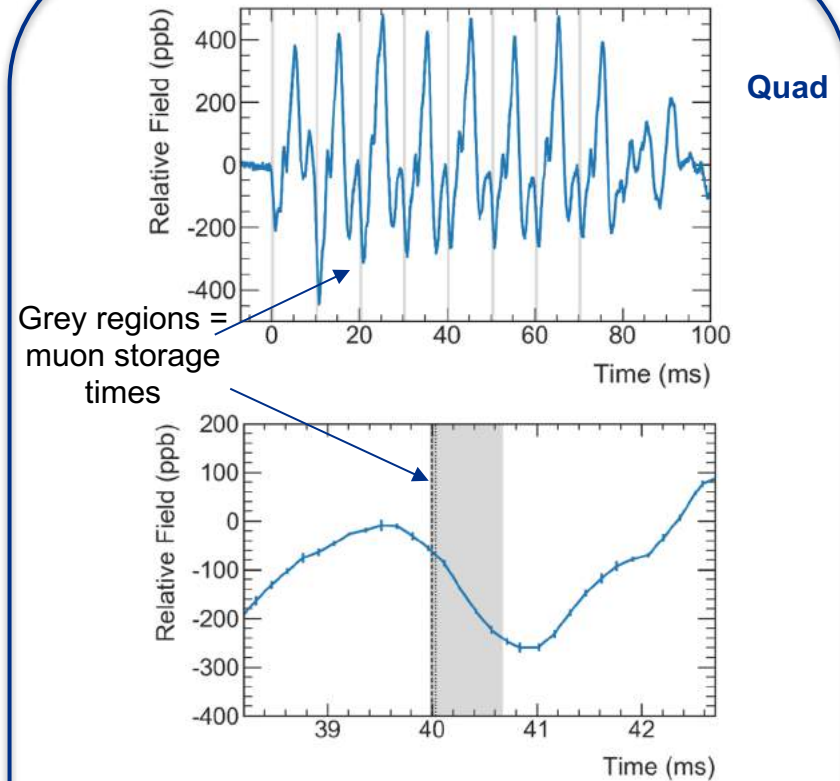
Transient fields



- Largest uncertainties come from “fast transient” fields generated by the pulsed systems (kickers and quads)
- Muons experience a field change which the fixed probes do not see (due to shielding)
- Effects were measured separately during dedicated measurement campaigns.

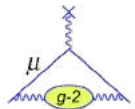
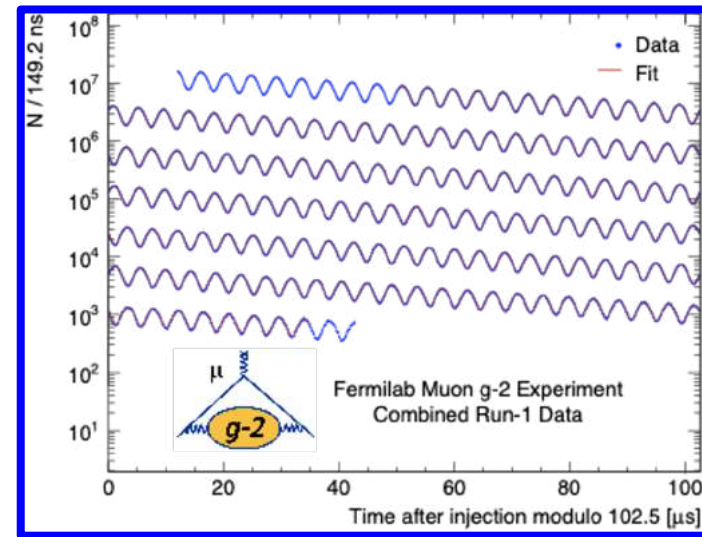


- Kicker pulse of **22 mT for 150 ns** just after muon injection.
- Field change caused by residual field after kicker pulse. Muons present from **30 μs to 700 μs** after the kick (fit region)
- Kicker correction: **-27 (37) ppb**

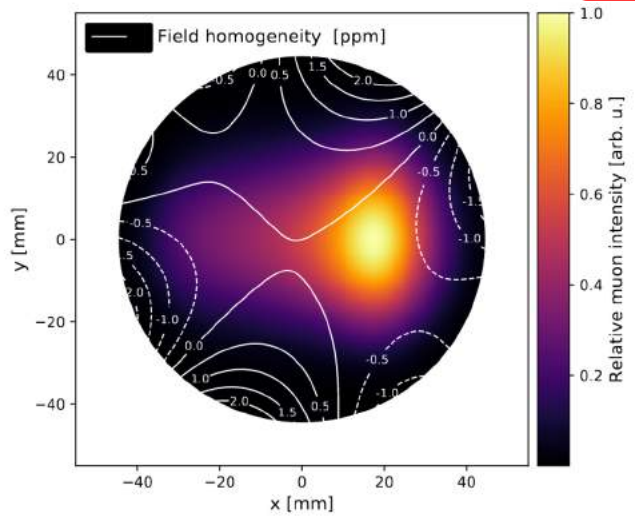


- Measured with a dedicated in-vacuum NMR probe located between quad plates during pulsing
- Quad correction: **-17 (92) ppb**

Bringing it all together

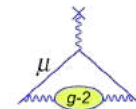


$$\mathcal{R}'_{\mu} = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

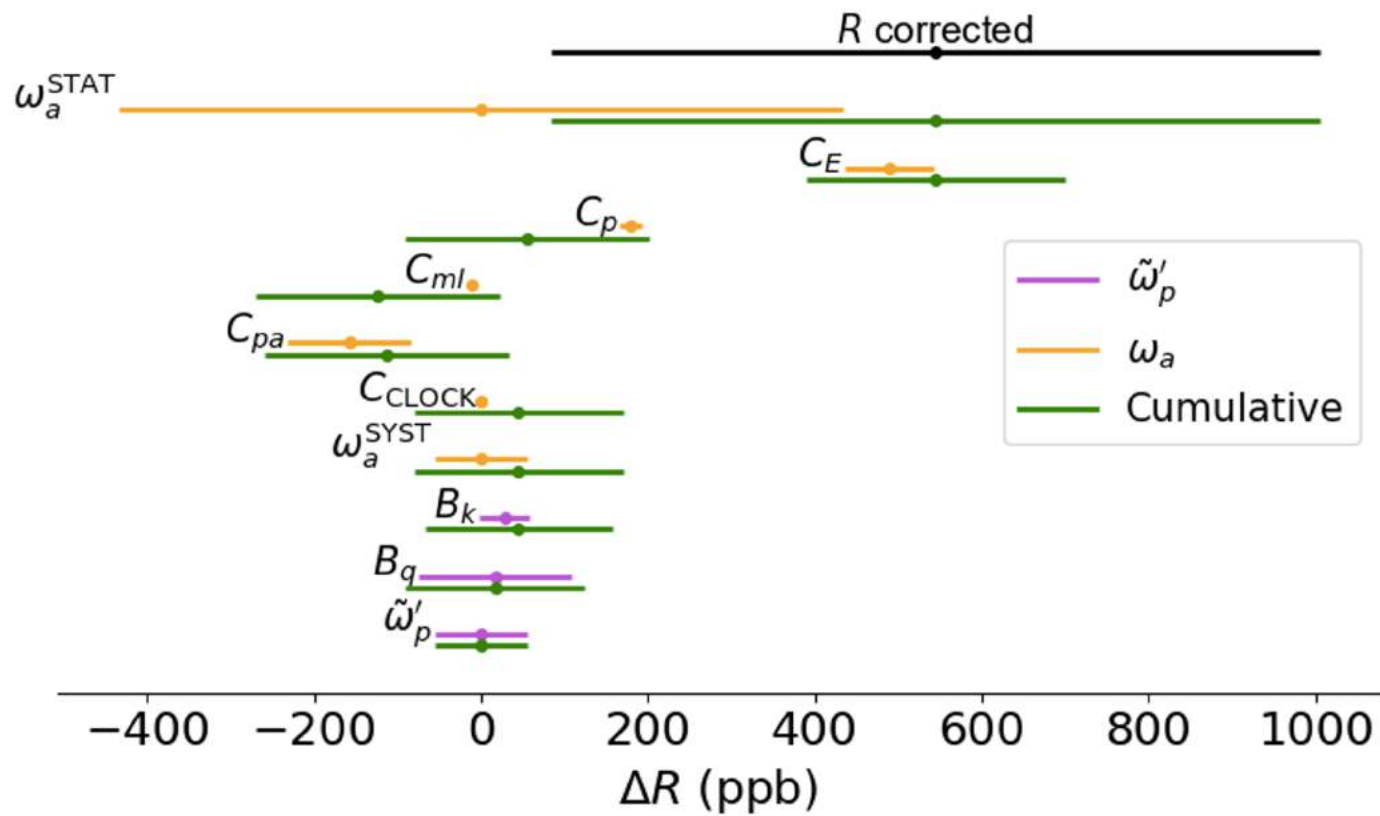


Corrections

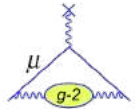
Correcting Measured R



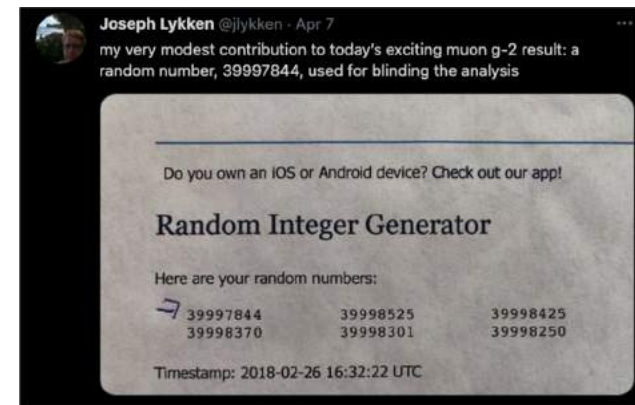
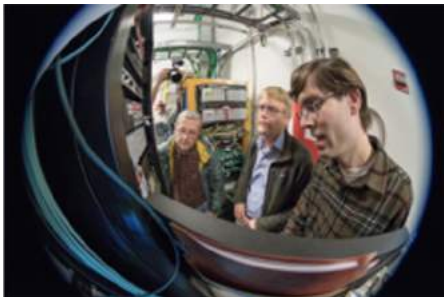
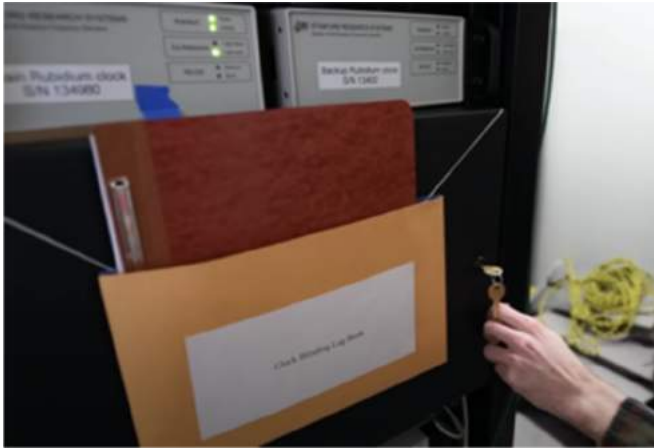
$$\mathcal{R}'_{\mu} = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$



Clock Blinding

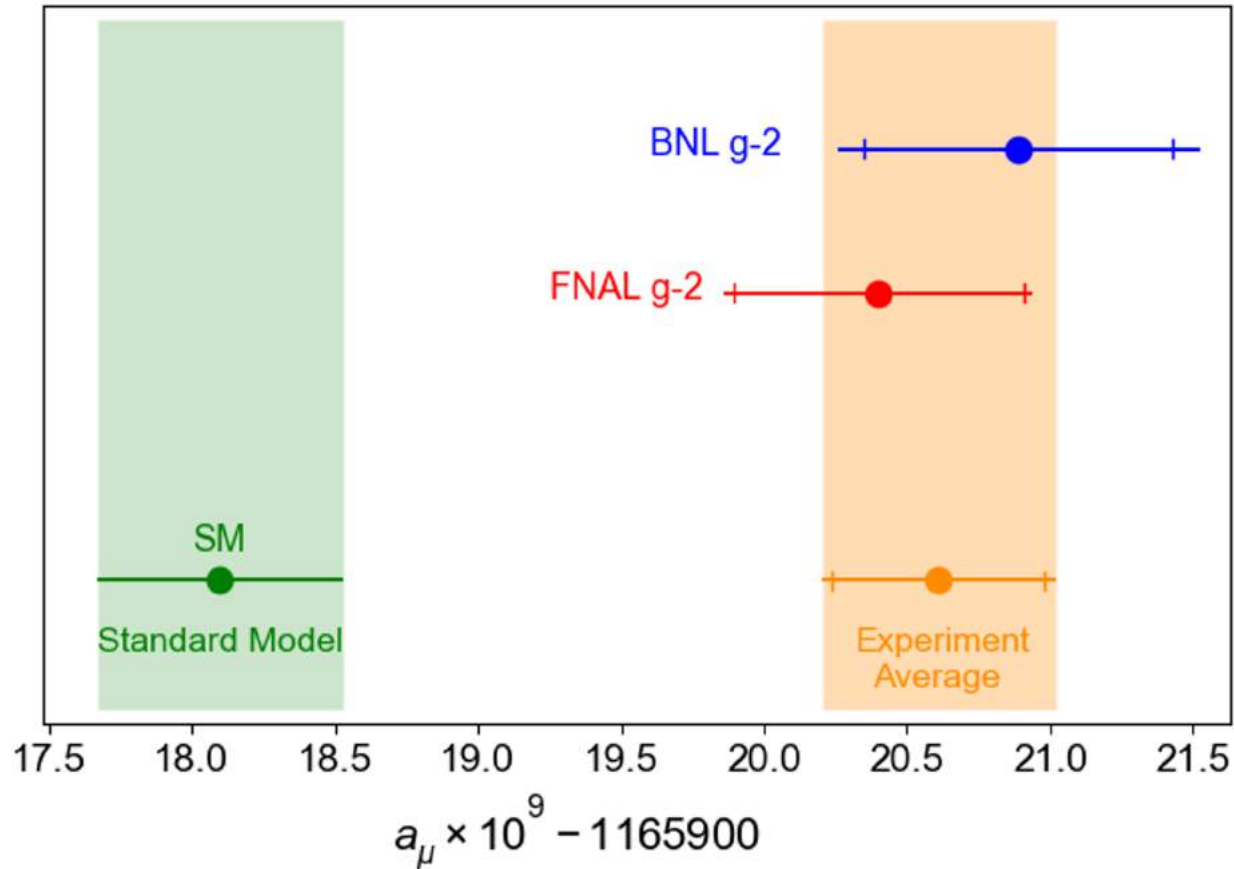
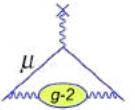


- The clock is hardware blinded to have a frequency of $(40 \pm \epsilon)$ MHz
- Only 2 people outside of the collaboration set and know the number
- Blinding offset was ± 25 ppm (approx $\times 10$ BNL-SM difference)

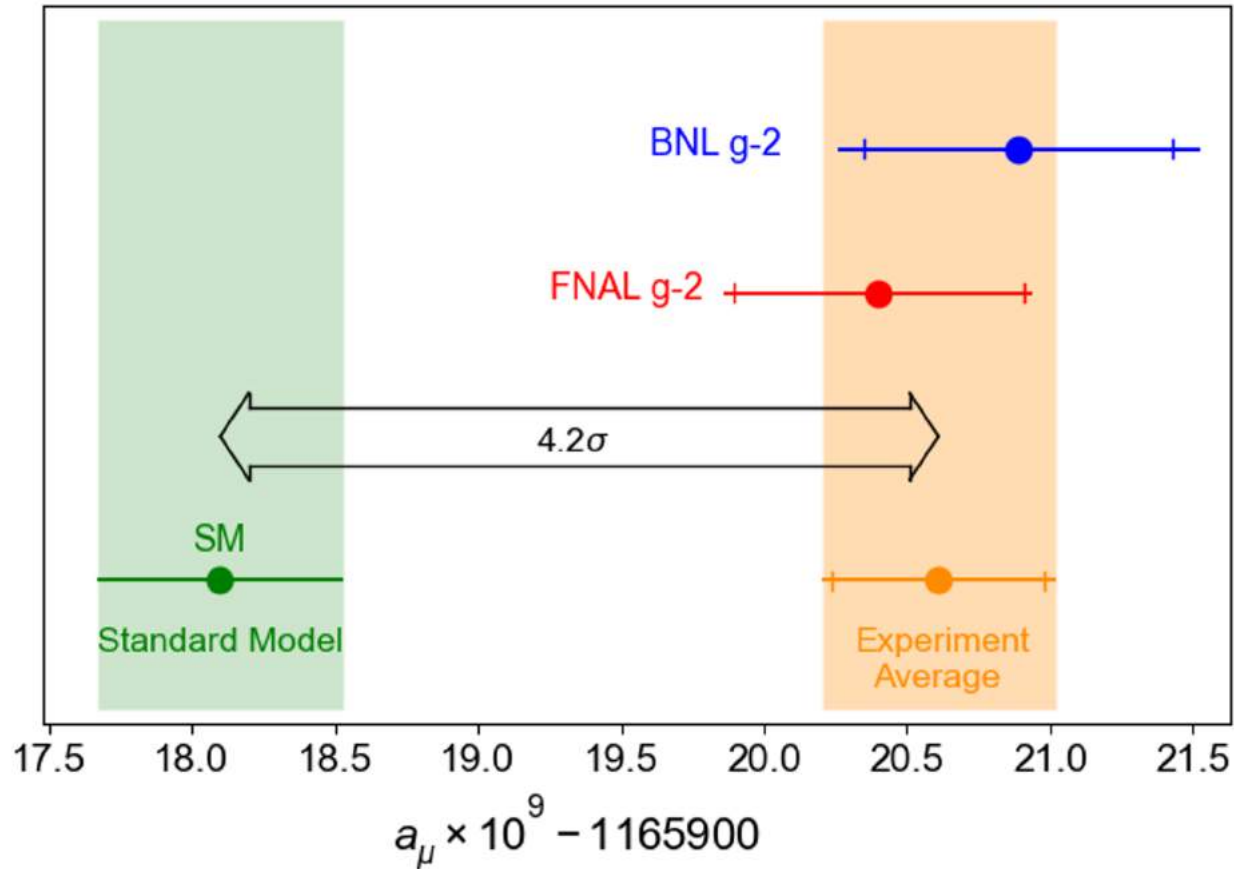
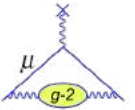


- Additionally each analysis is blinded in software

Unblinded result

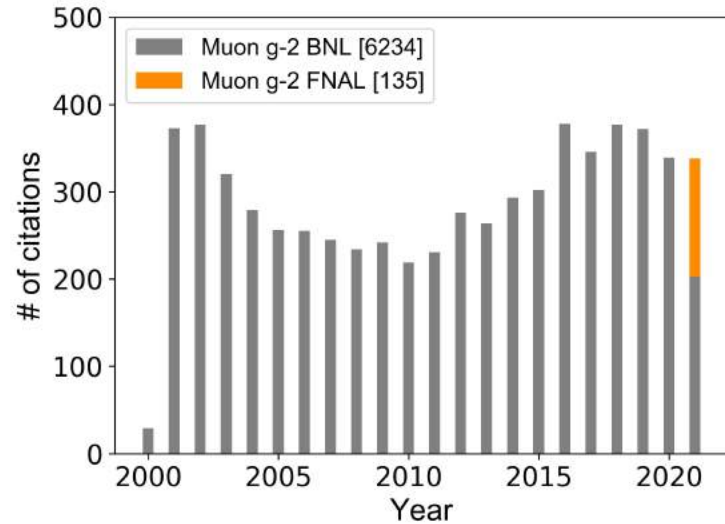
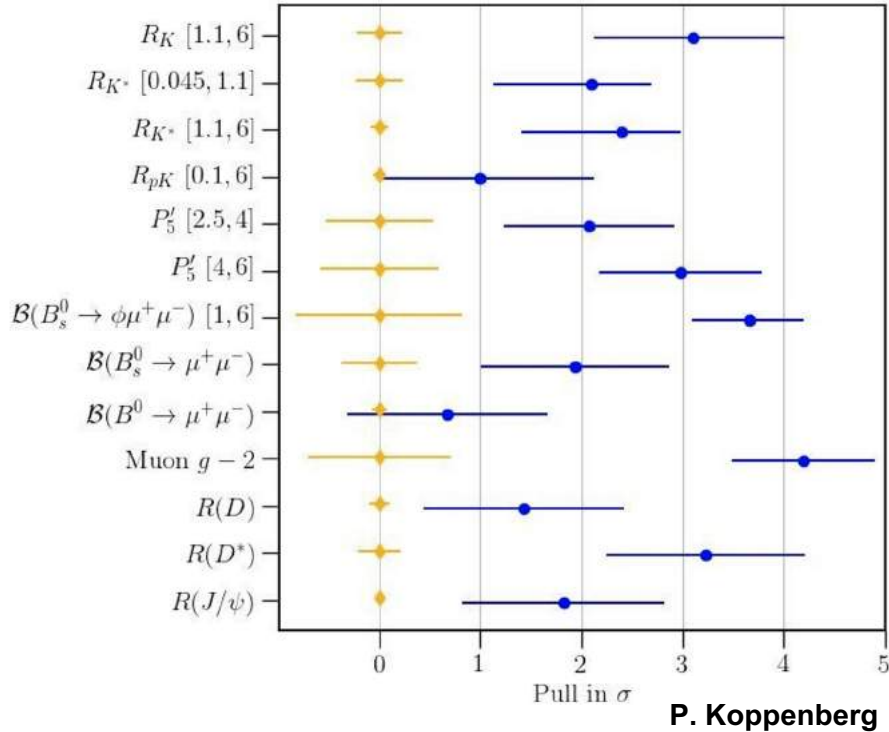
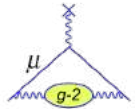


Unblinded result

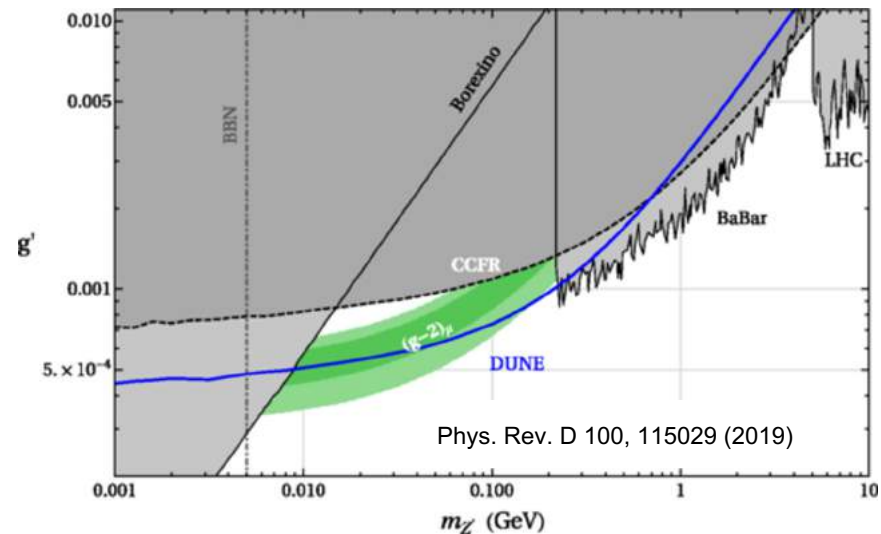


Interpretation

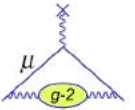
Needs more precision



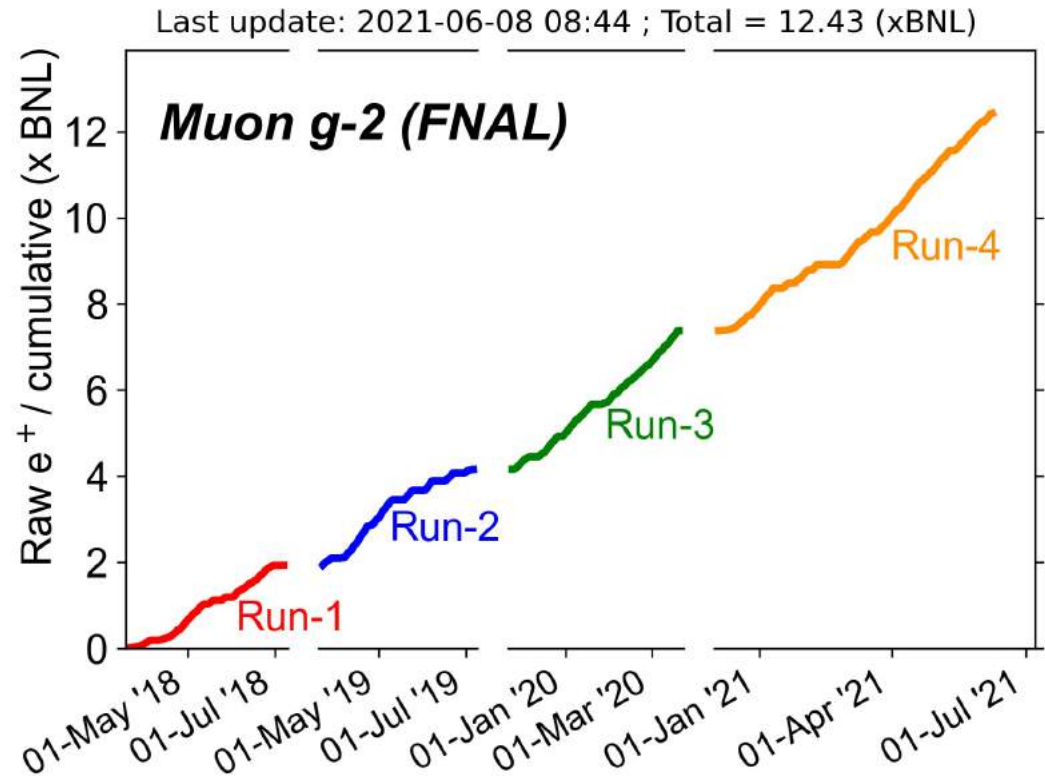
TeV Leptoquarks
 Z', ALPs
 LHC evading SUSY
 Tweaked Higgs extensions ...



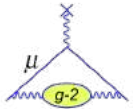
Conclusions



- The analysis of the Run-1 data produced a result with 460 ppb precision
- 4.2σ tension with the theoretical prediction
- There is a lot more data to analyse - expect a factor 2 improvement for Run-2/3 analysis



Thank you

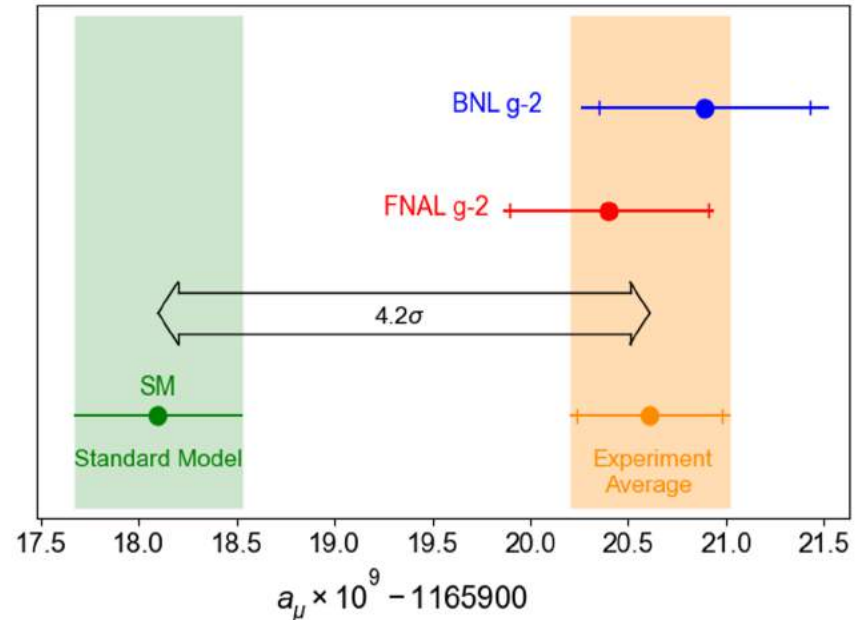


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Horizon 2020

- FNAL Main: [Phys.Rev.Lett. 126 \(2021\) 141801](#)
- FNAL ω_a : [Phys.Rev.D 103 \(2021\) 072002](#)
- FNAL Field: [Phys.Rev.A 103 \(2021\) 042208](#)
- FNAL Beam Dynamics: [arXiv:2104.03240 \(2021\)](#)



- Muon g-2 Theory Initiative (all contributions within): [Phys.Rept. 887 \(2020\) 1-166](#), <https://muon-gm2-theory.illinois.edu/white-paper/>
- HVP/HLbL Plots: [Aida X. El-Khadra, First results from the Muon g-2 experiment at Fermilab \(2021\)](#)
- BMW Lattice HVP (2021): [Nature \(2021\)](#)
- Mainz HLbL: [arXiv:2104.02632 \(2021\)](#)
- BNL Final: [Phys.Rev.D 73 \(2006\) 072003](#)
- Dune/g-2 Z' sensitivity: [Phys. Rev. D 100 \(2019\) 115029](#)
- BSM g-2: [arXiv:2104.03691 \(2021\)](#)