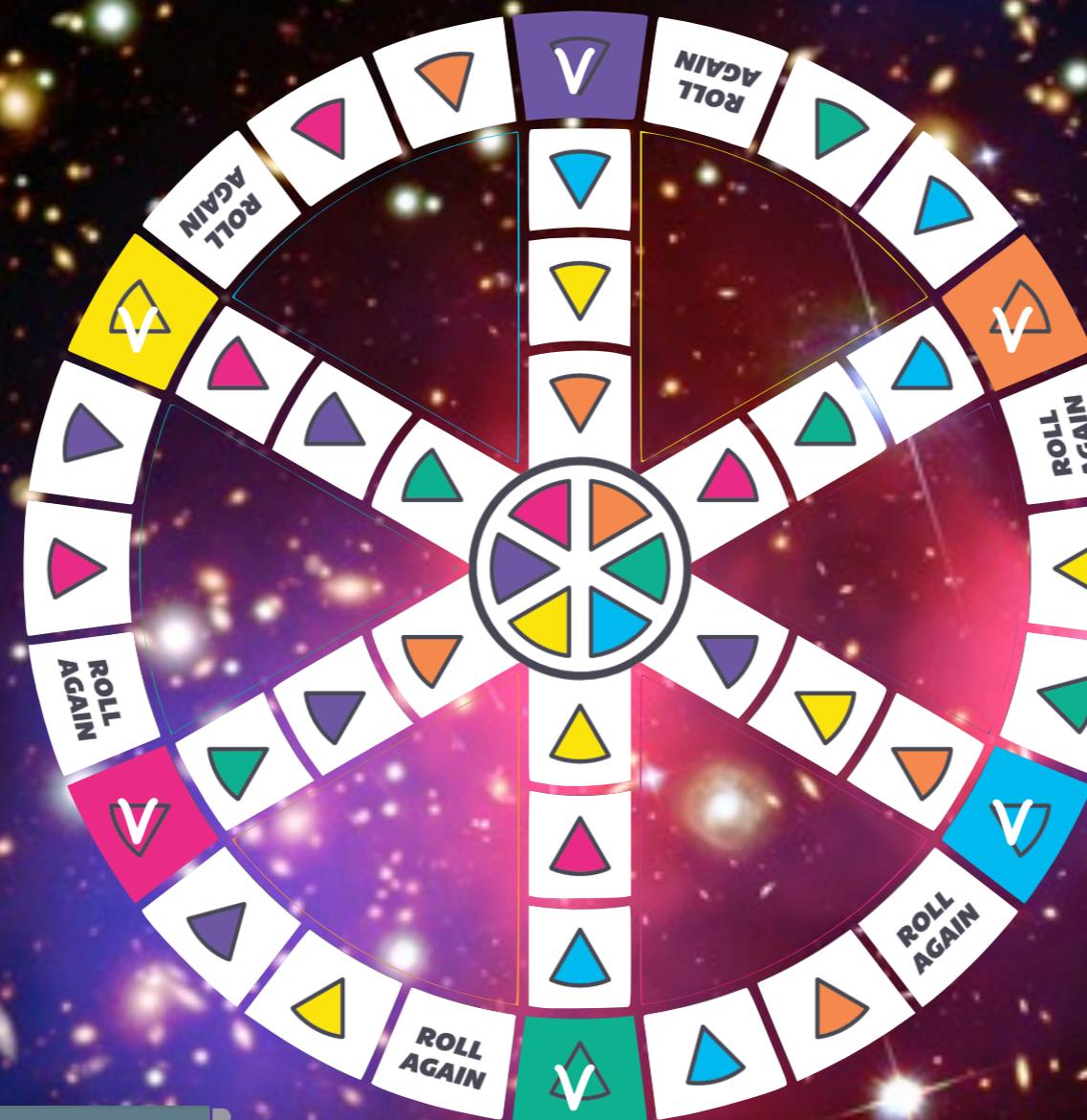


Olga Mena

IFIC-CSIC/UV Valencia (Spain)



The Cosmic NeuTRIVIA game

Cosmic neuTRIVIA game steps

1. Familiarize yourself with the board's layout:

- The Λ CDM trivia: the players!
- The neutrino pie piece: decoupling in the early universe

2. Roll the dice and get:

- Number of neutrinos and Big-Bang Nucleosynthesis
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- Neutrino masses and Cosmic Microwave Background Radiation?
- Neutrino masses and structure formation in the universe?

3. Is anyone cheating? Neutrinos and Tensions

4. Final score:

- Take home messages

Λ CDM NeuTRIVIA: PLAYERS

$$\Omega_i \equiv \frac{\rho_i}{\rho_c}$$

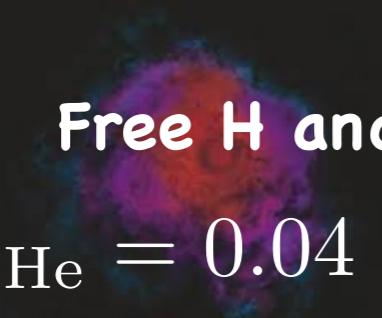


Λ CDM NeuTRIVIA: PLAYERS

$$\Omega_i \equiv \frac{\rho_i}{\rho_c}$$

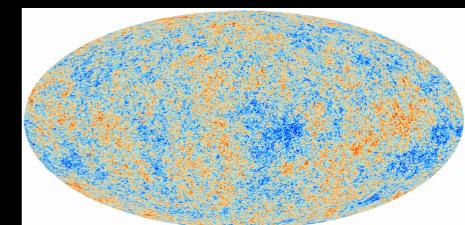
 Heavy elements
 $\Omega_{\text{Heavy Elements}} = 0.0003$

 Stars
 $\Omega_{\text{Stars}} = 0.005$

 Free H and He
 $\Omega_{\text{H, He}} = 0.04$



Dark energy
 $\Omega_\Lambda = 0.7$
v
Neutrinos!



$$\Omega_\gamma \sim 10^{-5}$$

$$\Omega_{\text{CDM}} = 0.25$$

How large is the MASSIVE neuTRIVIA pie piece?



Cosmology tells us

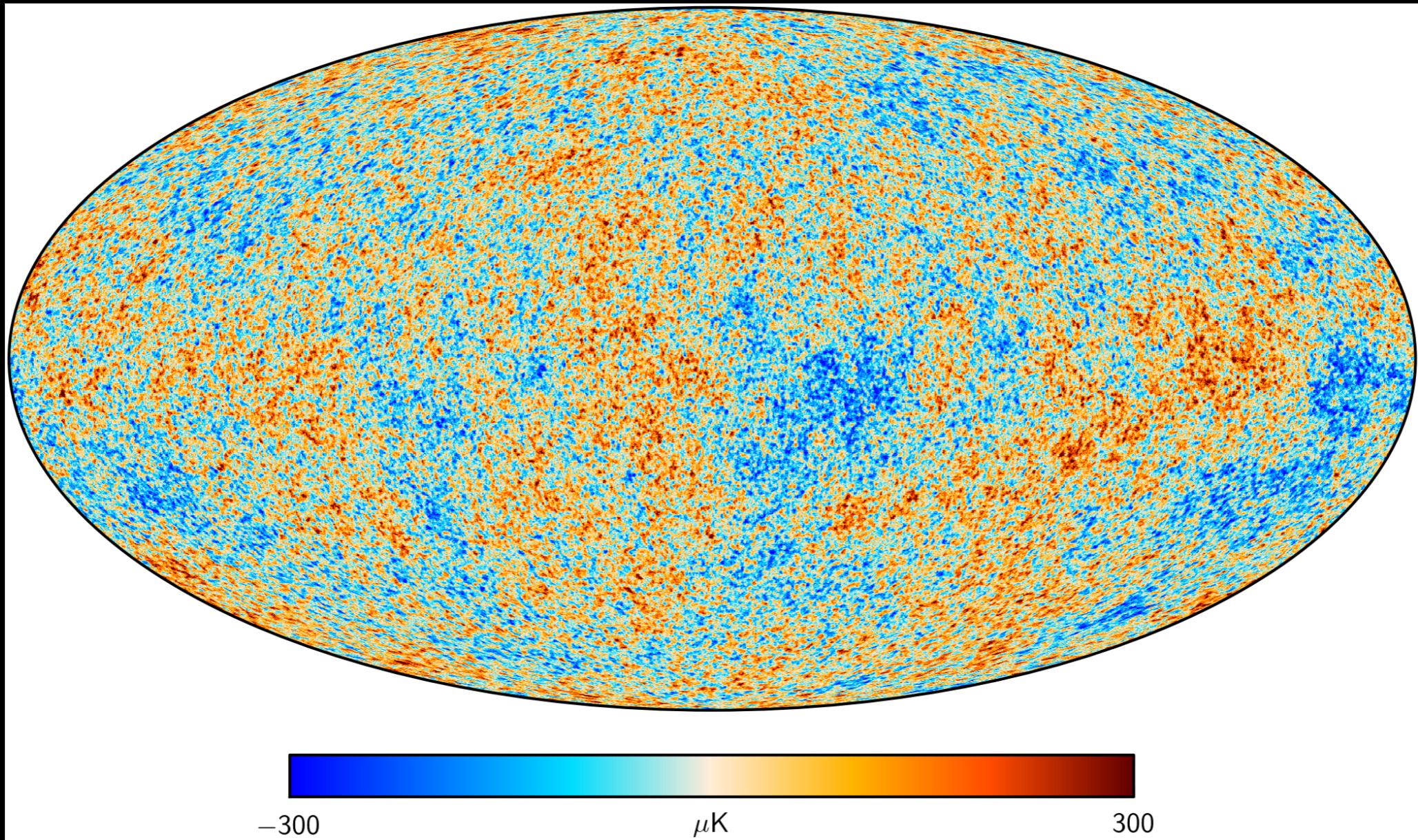
$$\Omega_\nu \lesssim 0.0024 \text{ 95% CL}$$

Neutrino
oscillations tell us

$$\Omega_\nu \gtrsim 0.0012 \text{ 95% CL}$$

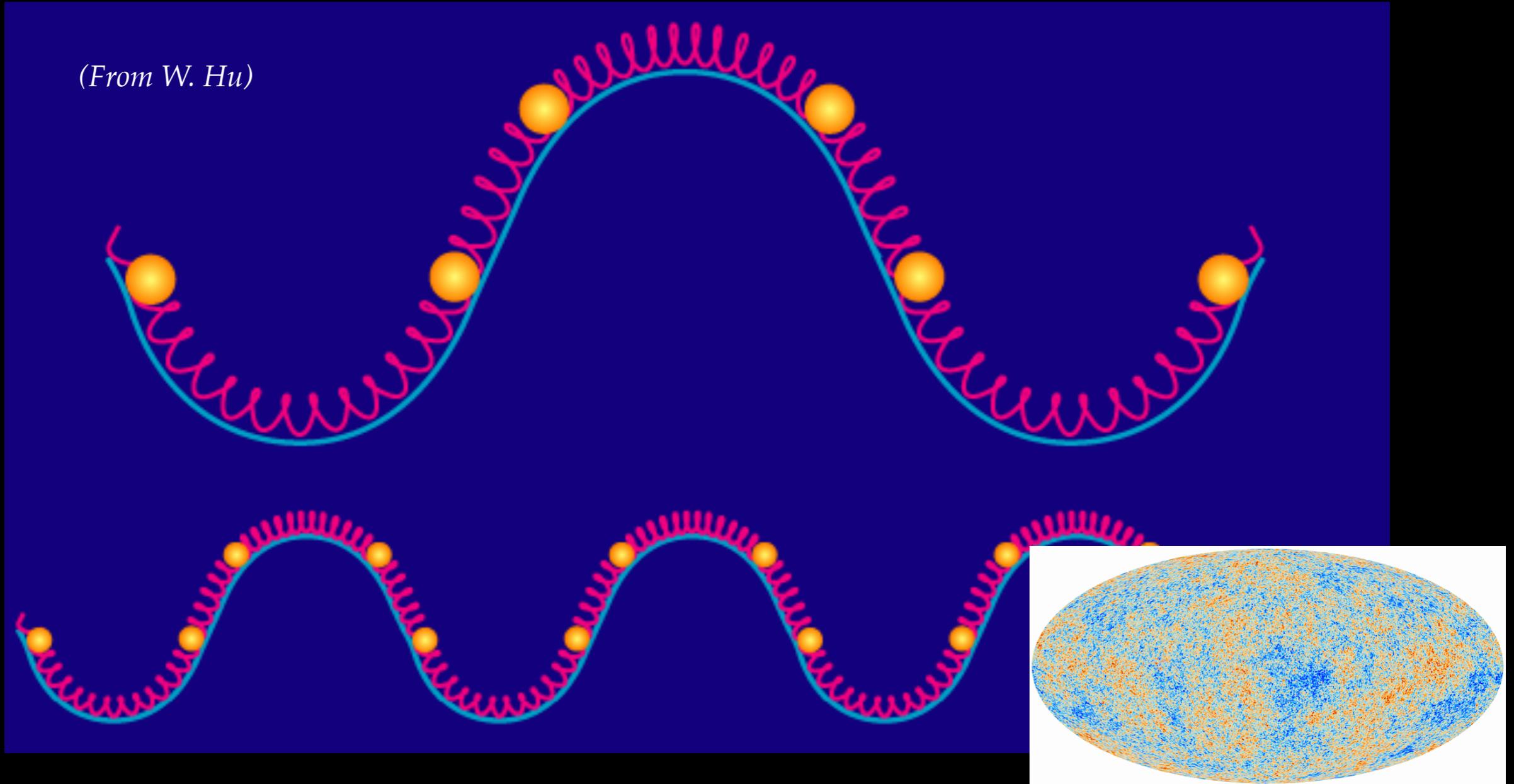
I guess all you know about dark matter/
what about dark radiation?

But radiation is visible, and it has a mean $T \approx 2.725$ K!



This map is just telling us how the CMB temperature fluctuations vary with the angular size of patches in the sky...

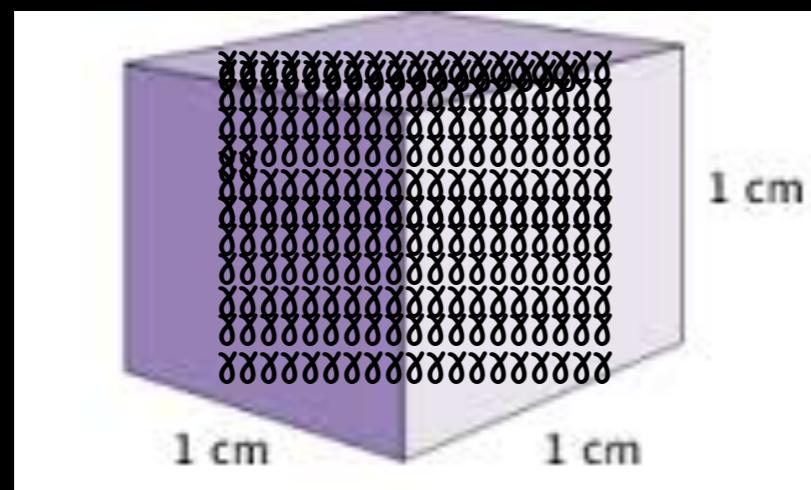
The CMB fluctuations are due to the acoustic oscillations in the baryon-photon fluid before recombination.



Potential wells \longrightarrow High density \longrightarrow COLD SPOTS in CMB maps
Potential hills \longrightarrow Low density \longrightarrow HOT SPOTS in CMB maps

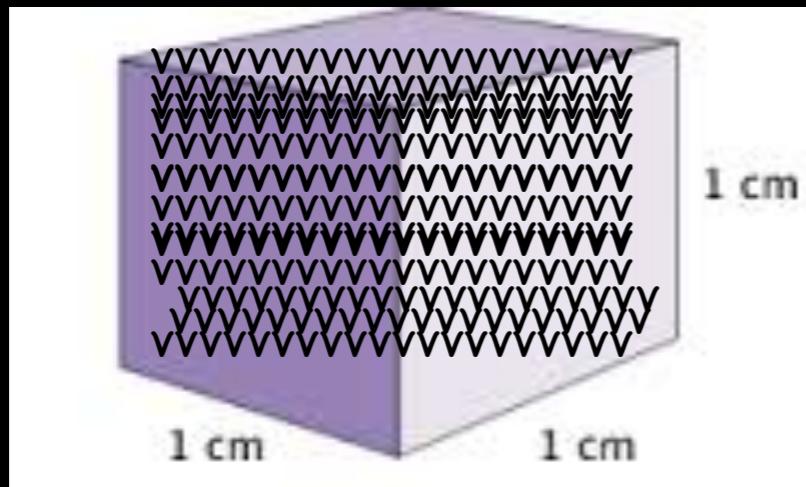
I guess all you know about dark matter/
what about dark radiation?
But radiation is visible!

410 photons/cm³



According to standard cosmology,
there is a cosmic neutrino background,
equivalent to the CMB photon background, albeit slightly colder $T \approx 1.94$ K

340 neutrinos/cm³

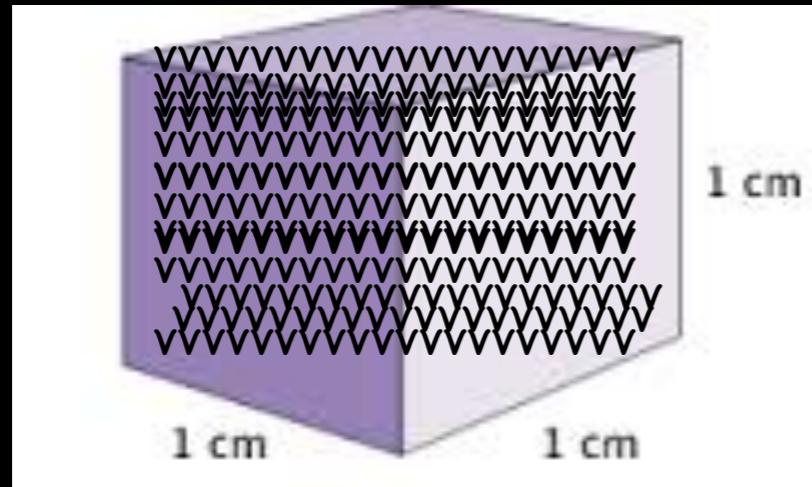


This cosmic relic neutrino background has never been detected directly.

The universe is filled with a dense flux of "relic neutrinos"
created in the Big Bang.

This makes neutrinos the most abundant KNOWN form of...

340 neutrinos/cm³



HOT dark matter!

According to standard cosmology,
 there are three active Dirac or Majorana neutrinos,
 which decouple from the thermal bath when their
 scattering rate is smaller than
 the expansion rate of the universe: $\Gamma_\nu \lesssim H$

- Neutrinos only interact via weak interactions, with a rate:

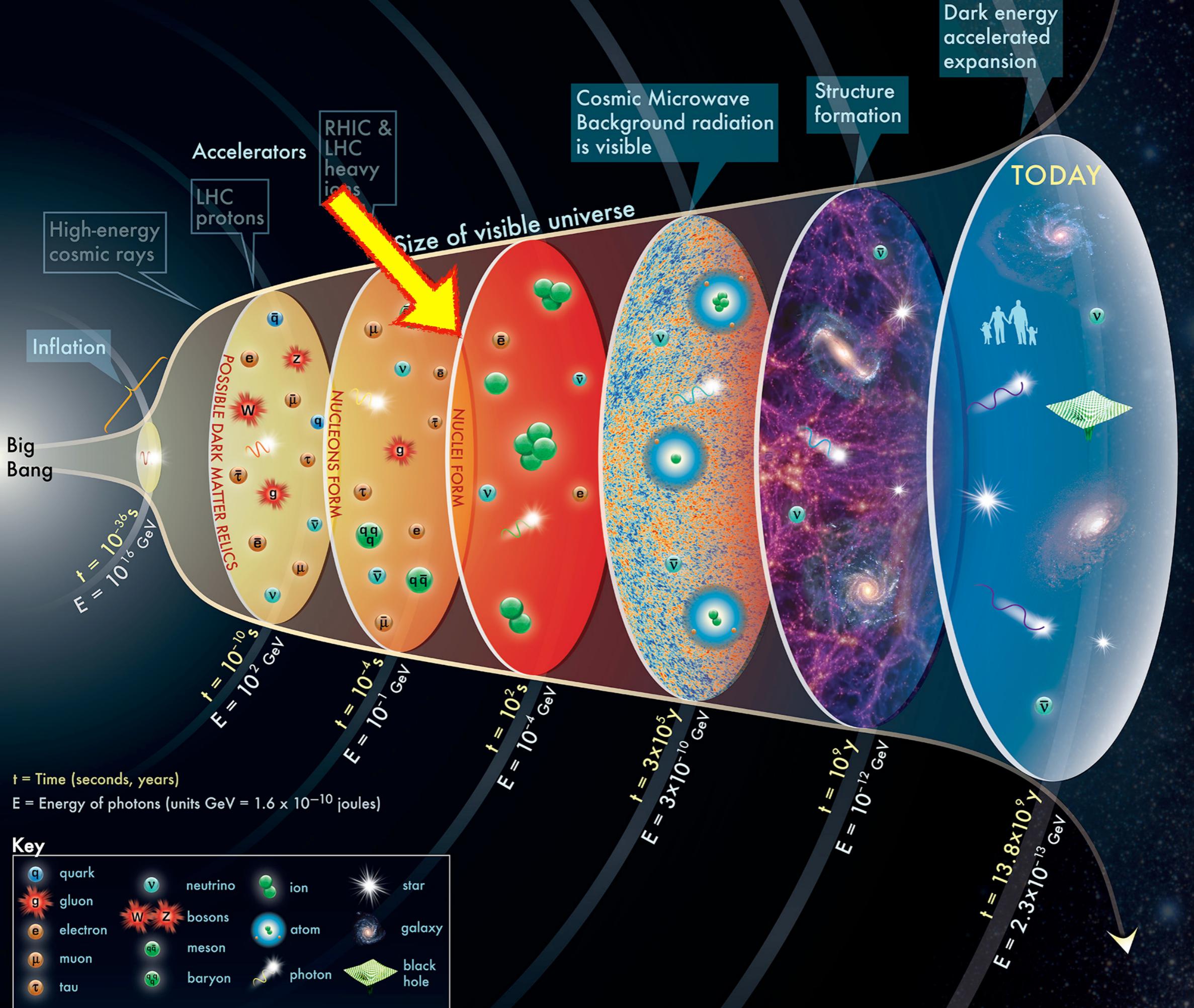
$$\Gamma_\nu = n\sigma v \simeq T^3 G_F^2 T^2 \sim G_F^2 T^5$$

- While the expansion rate of the universe is given by the Hubble factor:

$$H^2 = \frac{8\pi G}{3} \rho \sim T^4 / m_{pl}^2$$

$$\Gamma_\nu / H \sim \left(\frac{T}{1 \text{ MeV}} \right)^3$$

- Therefore neutrinos decouple from the thermal bath around 1 MeV.



The concept for the above figure originated in a 1986 paper by Michael Turner.

Event	Time	Redshift	Temperature
Baryogenesis	?	?	?
EW phase transition	$2 \times 10^{-11} s$	10^{15}	$100 GeV$
QCD phase transition	$2 \times 10^{-5} s$	10^{12}	$150 MeV$
Neutrino decoupling	1s	6×10^9	$1 MeV$
Electron-positron annihilation	6s	2×10^9	$500 keV$
Big bang nucleosynthesis	3min	4×10^8	$100 keV$
Matter-radiation equality	$6 \times 10^4 yrs$	3400	.75eV
Recombination	$2.6 - 3.8 \times 10^5 yrs$	1100-1400	.26 – .33eV
CMB	$3.8 \times 10^5 yrs$	1100	.26eV

Relic neutrinos do not inherit any of the energy associated to $e^+ e^-$ annihilations, being colder than photons:

$$T_{\nu 0} = \left(\frac{4}{11}\right)^{1/3} T_0 = 1.945 \text{ K} \sim 1.697 \times 10^{-4} \text{ eV}$$

If these neutrinos are massive, their energy density, at $T \ll m$ is

$$\rho_\nu = m_\nu n_\nu \quad n_{\nu_i}(T_{\nu 0}) \approx 56 \text{ cm}^{-3} \quad \Omega_\nu h^2 = \frac{\sum m_\nu}{93 \text{ eV}}$$

Then, demanding that massive neutrinos do not over-close the universe, $\sum m_\nu \lesssim 45 \text{ eV}$

Their thermal motion is: $\langle v_{\text{thermal}} \rangle \simeq 81(1+z) \left(\frac{\text{eV}}{m_\nu}\right) \text{ km s}^{-1}$

For a 1 eV neutrino, thermal motion is comparable to the typical velocity dispersion of a galaxy.

For dwarf galaxies,
the velocity dispersion is smaller, 10 km/s



oooooooooooooo



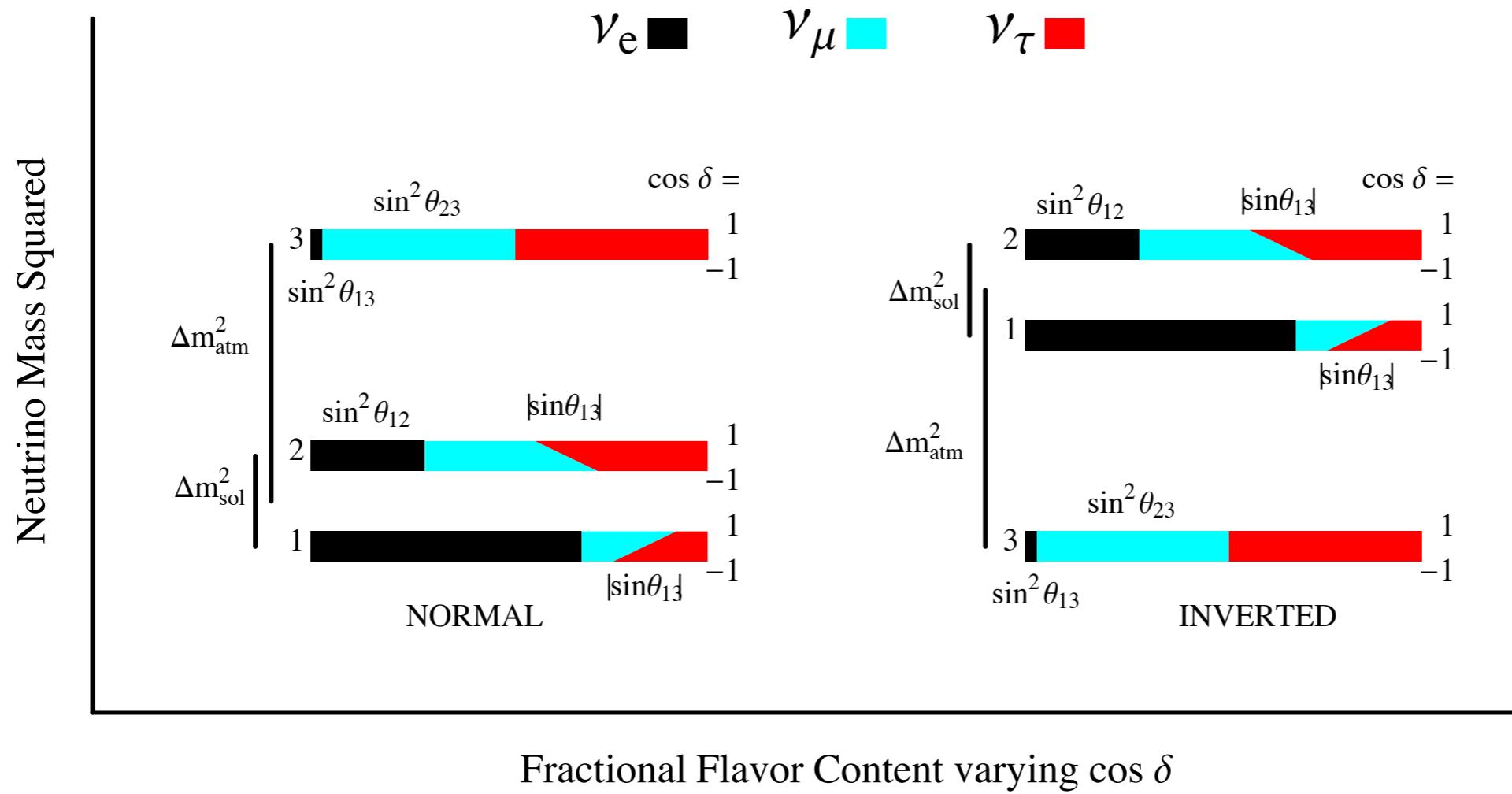
Too much thermal energy to be squeezed into small volumes to form the smaller structures we observe today! ¹⁴

According to **neutrino oscillation physics**,
we know that there are at least two Dirac or
Majorana **massive neutrinos**:

$$\Delta m_{12}^2 = (7.05 - 8.14) \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{13}^2 = (2.41 - 2.60) \times 10^{-3} \text{ eV}^2$$

$$\Delta m_{13}^2 = -(2.31 - 2.51) \times 10^{-3} \text{ eV}^2$$



(Mena, Parke, PRD'04)

According to **neutrino oscillation physics**, we know that there are at least two Dirac or Majorana **massive neutrinos**:

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We are sure then that **two neutrinos have a mass above**:

$$\sqrt{\Delta m_{12}^2} \simeq 0.008 \text{ eV}$$

and that **at least one of these neutrinos has a mass larger than**

$$\sqrt{|\Delta m_{13}^2|} \simeq 0.05 \text{ eV}$$

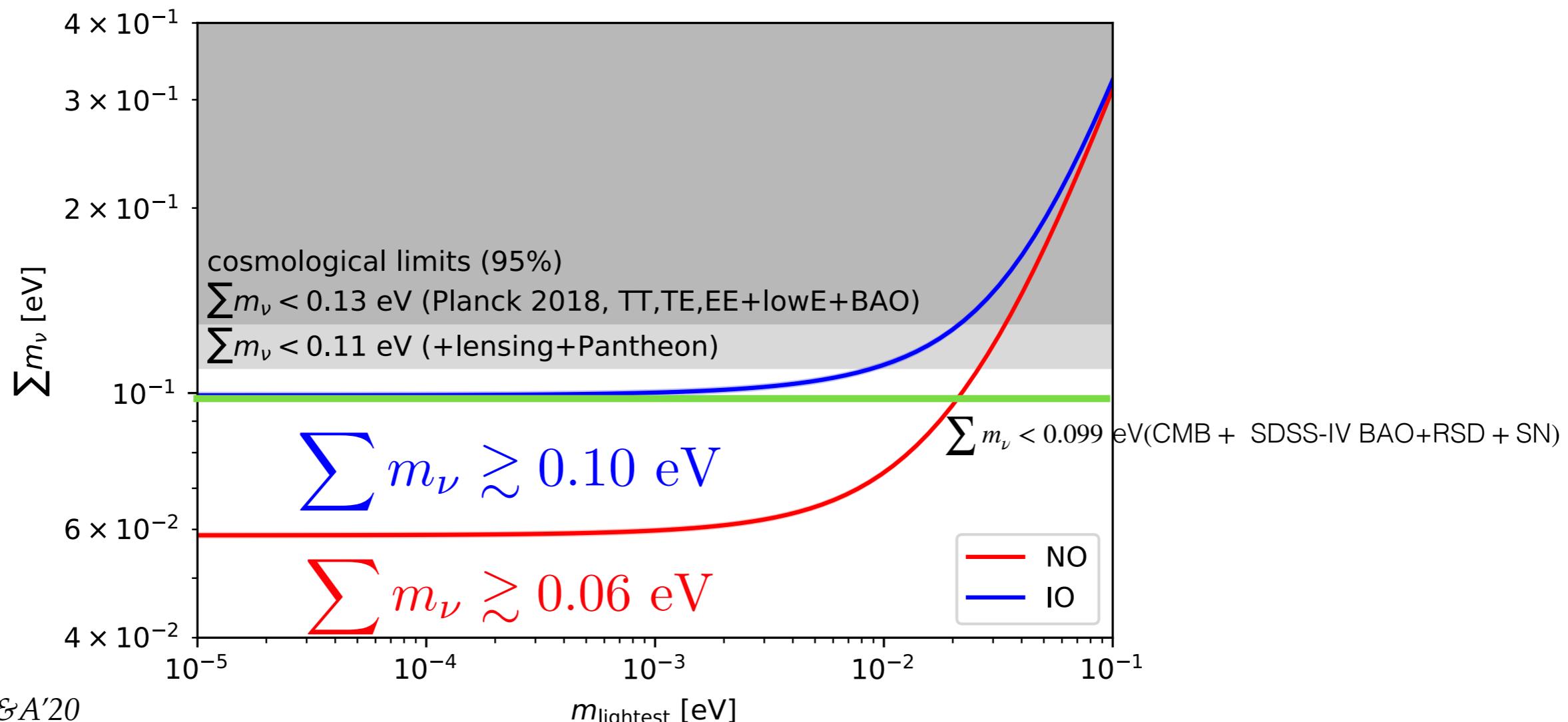
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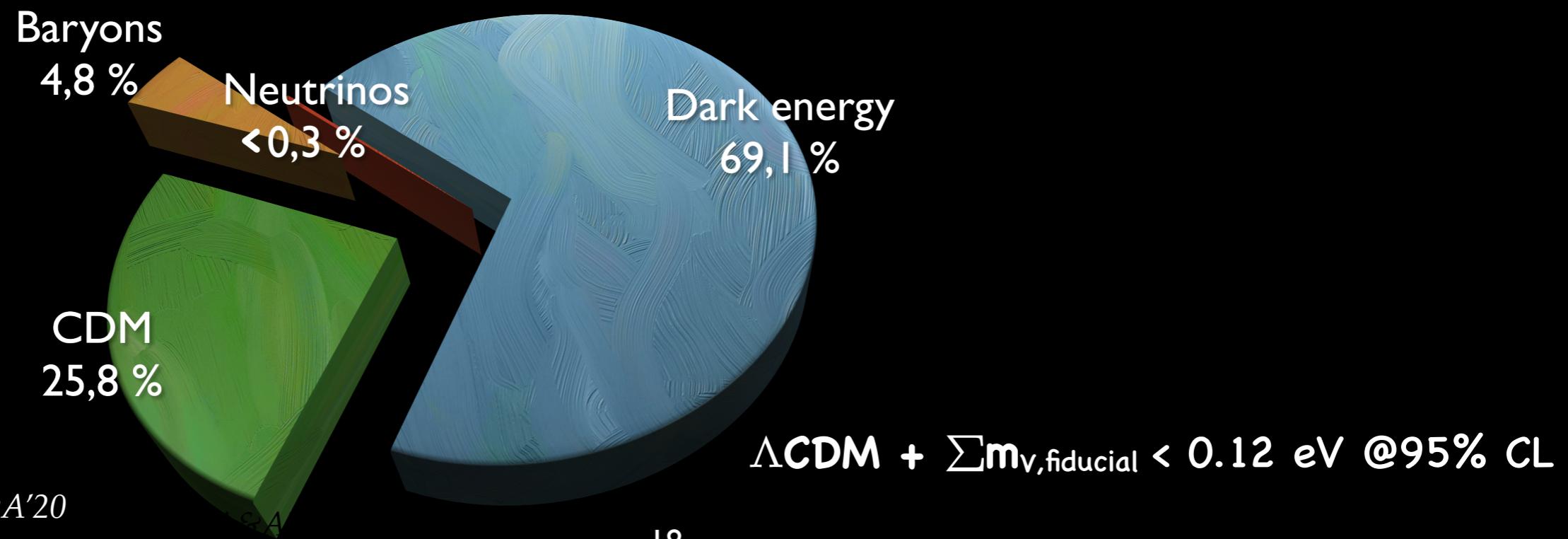
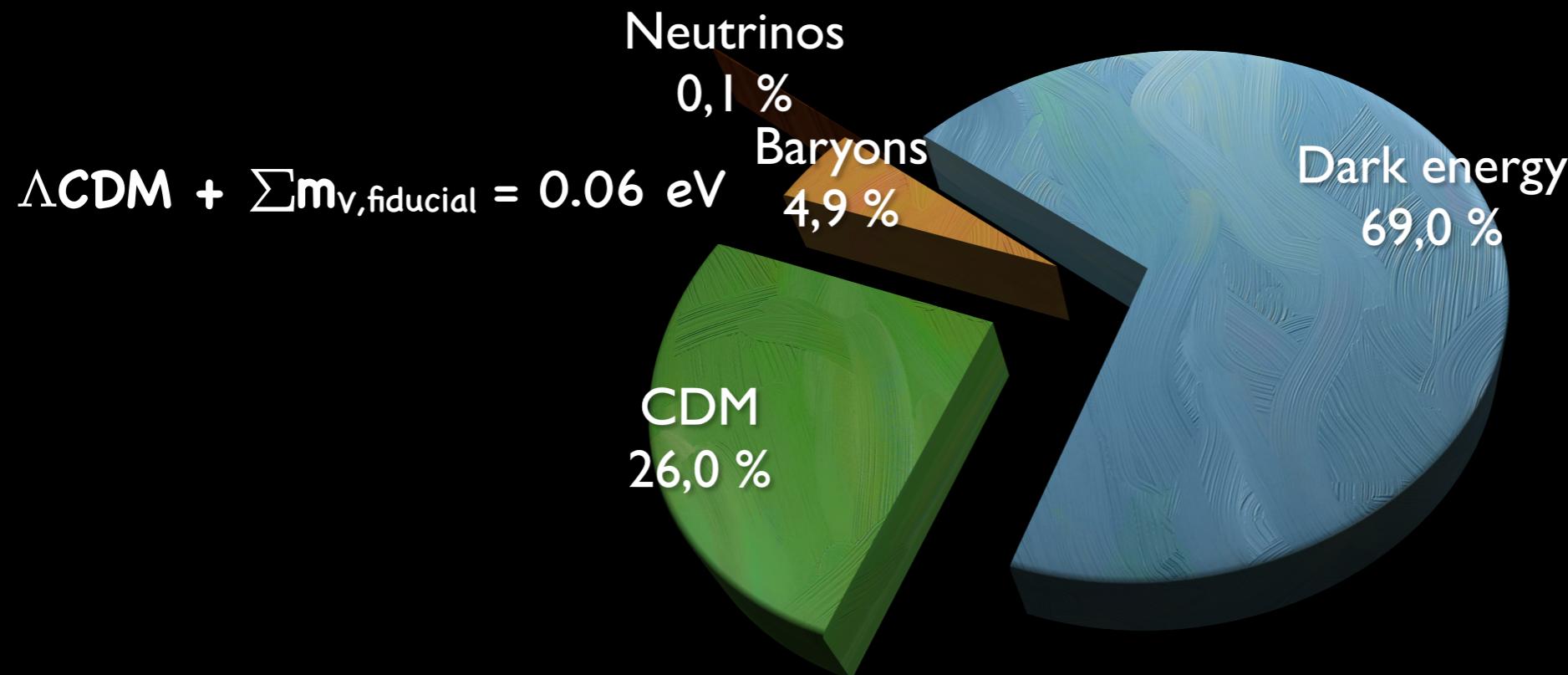
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which translates into a lower bound on the total neutrino mass, depending on the ordering:



Planck 2018 Cosmic Trivia



Cosmic neuTRIVIA game steps

1. Familiarize yourself with the board's layout:

- The Λ CDM trivia: the players!
- The neutrino pie piece: decoupling in the early universe

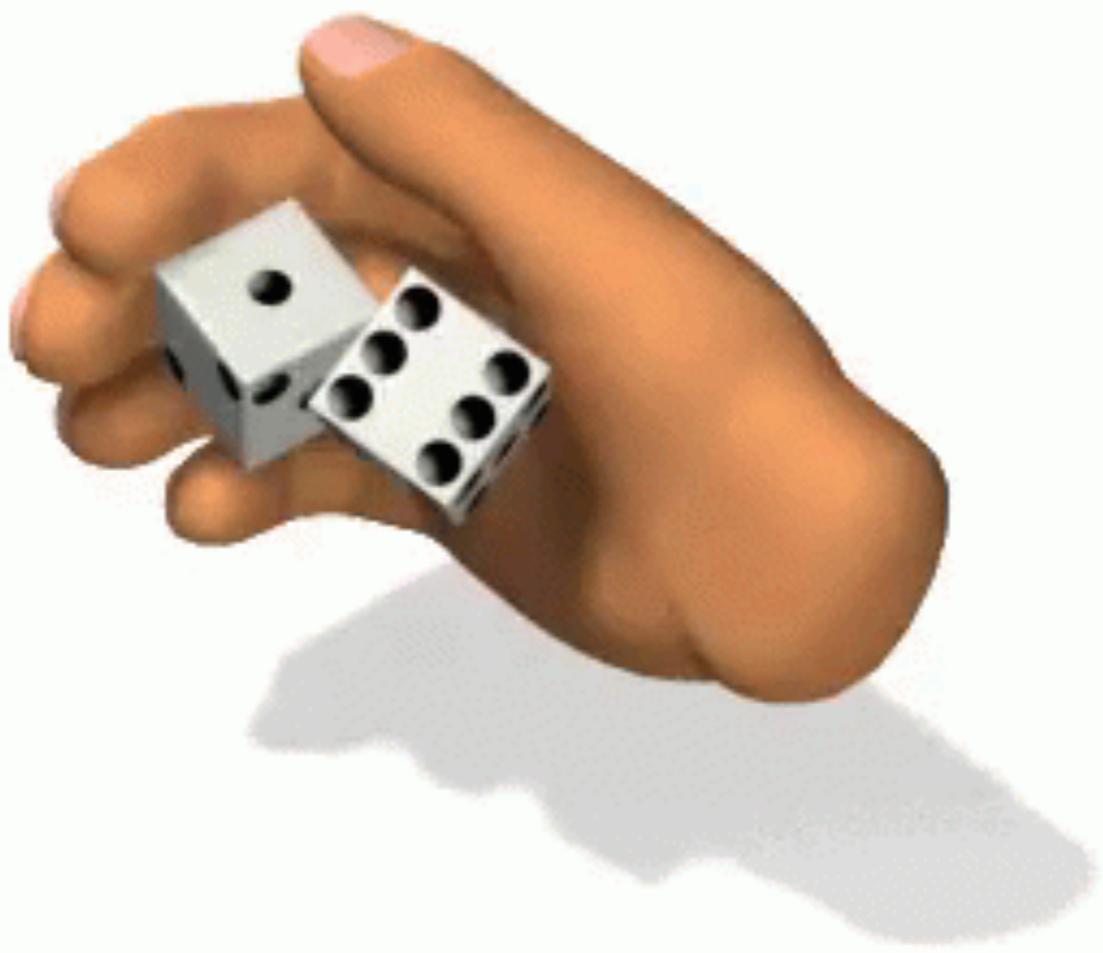
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3. Is anyone cheating? Neutrinos and Tensions

4. Final score:

- Take home messages



Number of neutrinos: N_{eff}

The total radiation in the universe can be written as:

$$\Omega_r h^2 = \left(1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right) \Omega_\gamma h^2$$

Bennett et al, 2012.02726

$N_{\text{eff}} = 3.0440 \pm 0.0002$ standard scenario: electron, muon and tau neutrinos

$N_{\text{eff}} < 3.044$ (less neutrinos): Neutrino decays ?

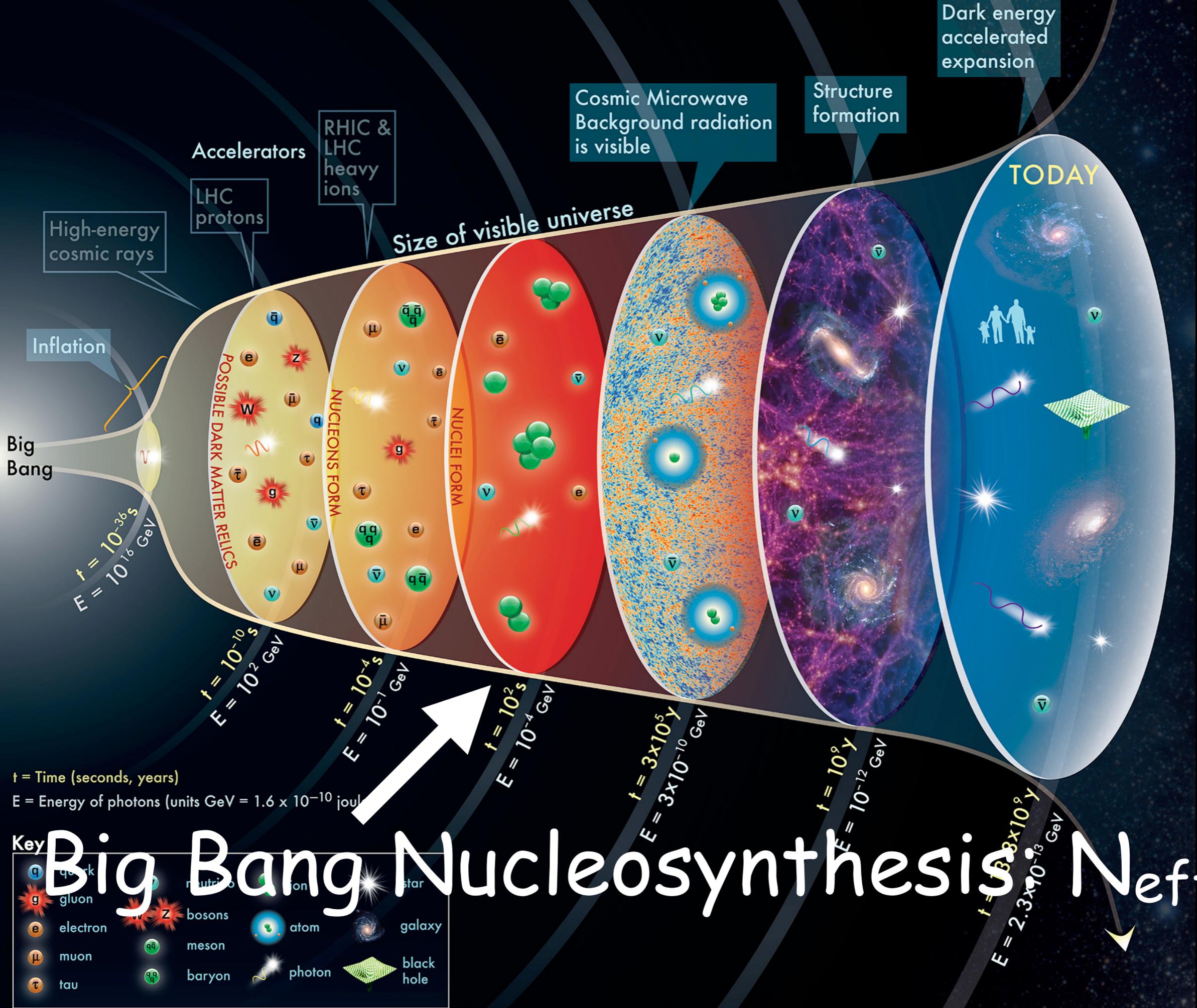
$N_{\text{eff}} > 3.044$ (more neutrinos): Sterile neutrino species ?

But....if they are sterile, and do not interact with other particles, how cosmologists measure them?



That's the dark side of the GRAVITATIONAL FORCE...





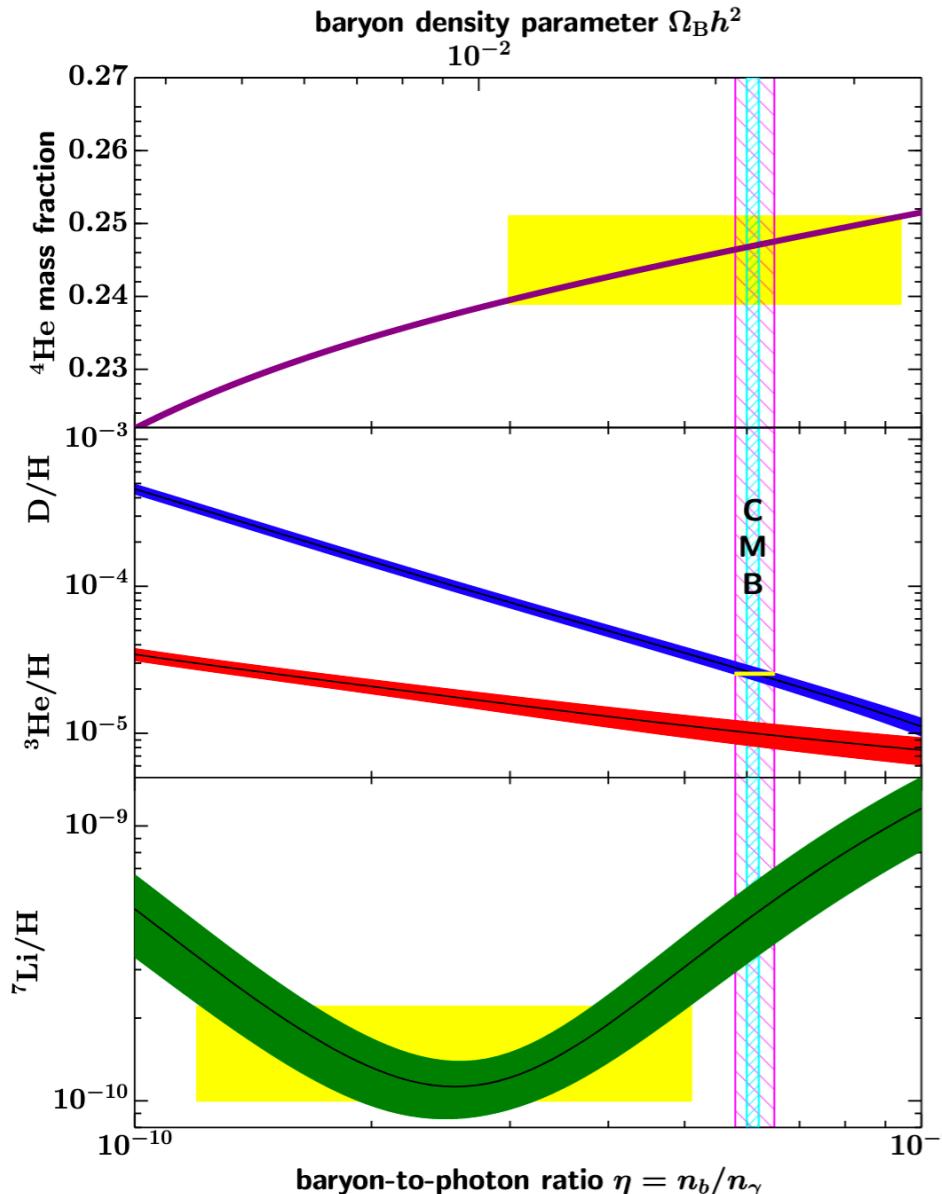
Event	Time	Redshift	Temperature
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Big Bang Nucleosynthesis: N_{eff}

BBN theory predicts the abundances of D, ^3He , ^4He and ^7Li which are fixed by $t \approx 180$ s. They are observed at late times: low metallicity sites with little evolution are “ideal”.

Big Bang Nucleosynthesis: Neff

BBN theory predicts the abundances of D, ^3He , ^4He and ^7Li which are fixed by $t \approx 180$ s. They are observed at late times: low metallicity sites with little evolution are “ideal”.



Low metallicity extragalactic HII regions.



Produced in stars.



High z QSO absorption lines.



Destroyed in stars.



Solar system and high metallicity HII galactic regions.

^3He not used for cosmological constraints.



Metal poor stars in our galaxy.



Destroyed in stars and produced by galactic cosmic ray interactions.



Figure 24.1: The primordial abundances of ^4He , D, ^3He , and ^7Li as predicted by the standard model of Big-Bang nucleosynthesis — the bands show the 95% CL range [47]. Boxes indicate the observed light element abundances. The narrow vertical band indicates the CMB measure of the cosmic baryon density, while the wider band indicates the BBN D+ ^4He concordance range (both at 95% CL).

Big Bang Nucleosynthesis: N_{eff}

N_{eff} changes the freeze out temperature of weak interactions:

$$\Gamma_{n \leftrightarrow p} \sim H$$

MORE NEUTRINOS:

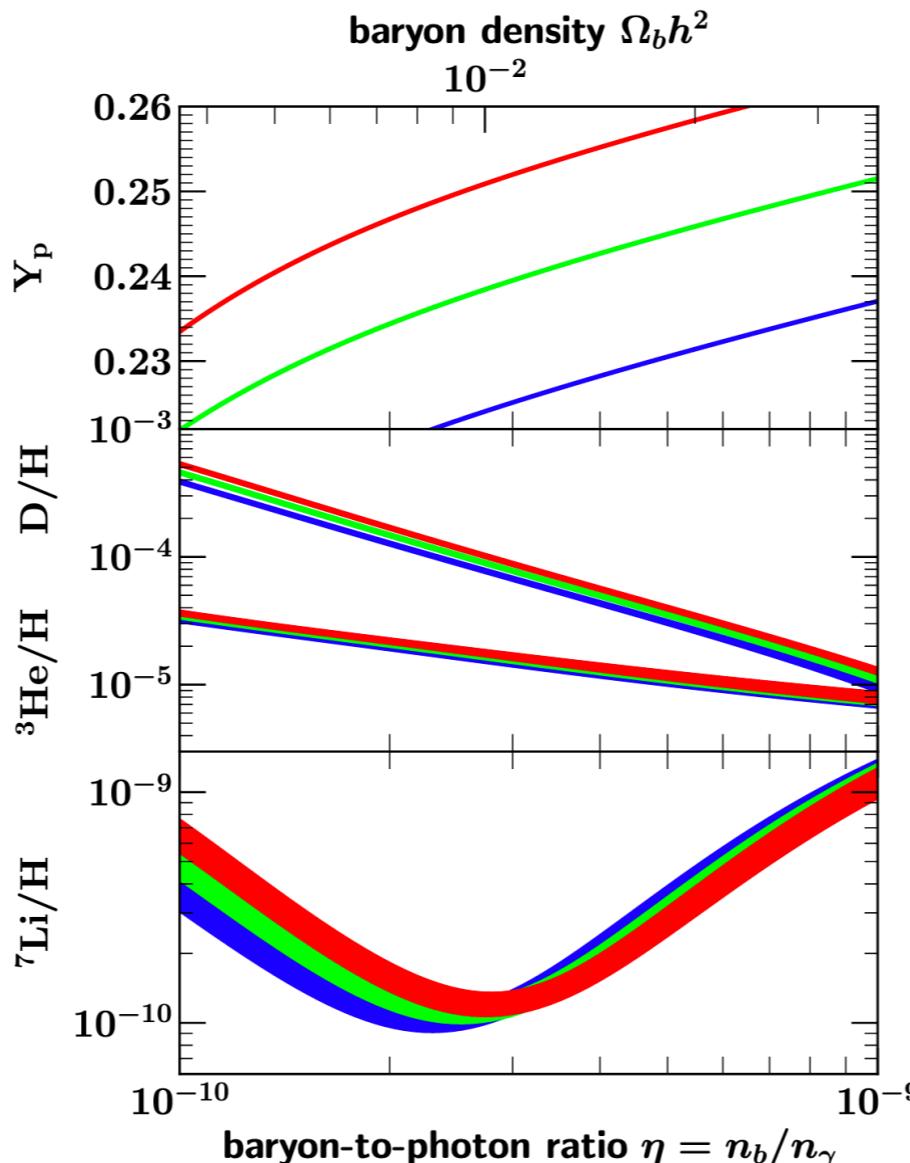
Higher N_{eff} : larger expansion rate & freeze out temperature, MORE HELIUM 4

$$n/p \simeq e^{-\frac{m_n - m_p}{T_{\text{freeze}}}} \quad Y_p = \frac{2(n/p)}{1 + n/p}$$

$$N_{\text{eff}} = 2$$

$$N_{\text{eff}} = 3$$

$$N_{\text{eff}} = 4$$



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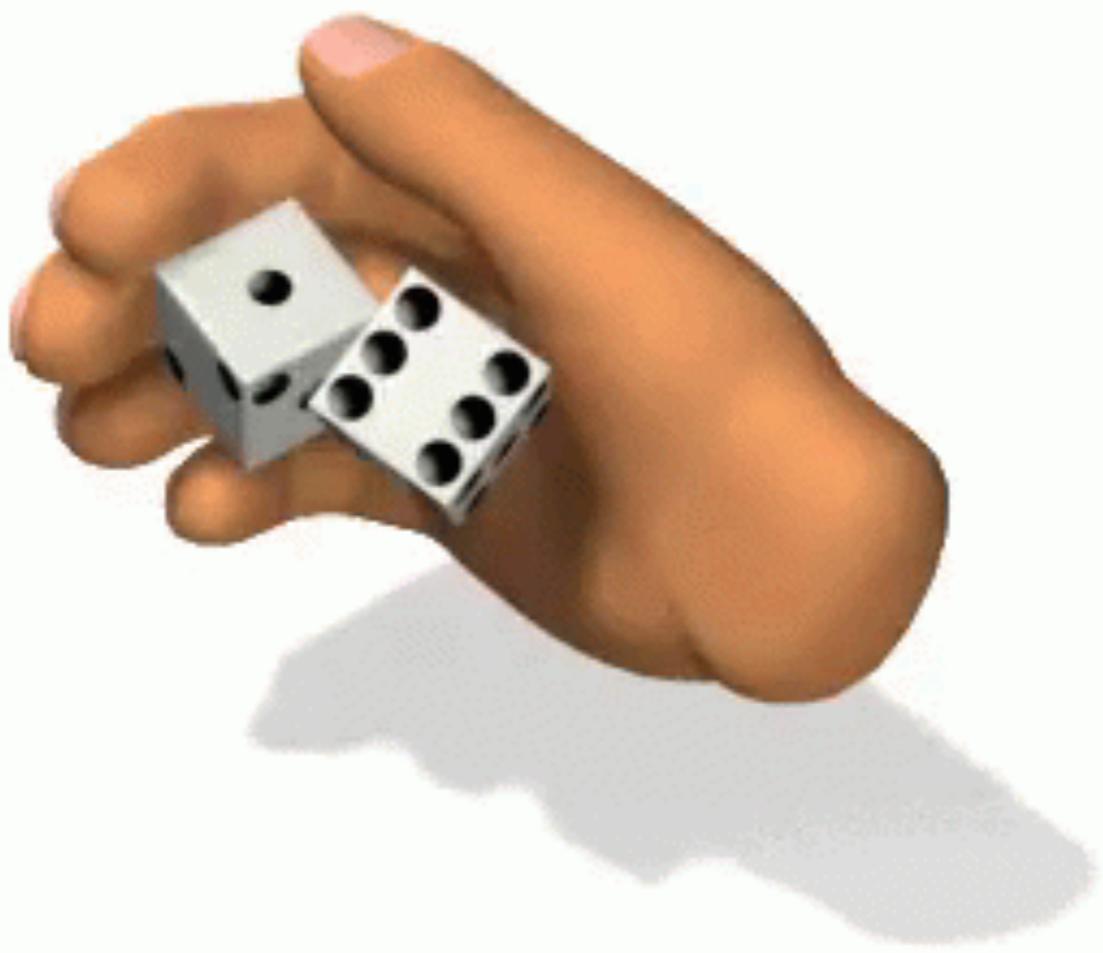
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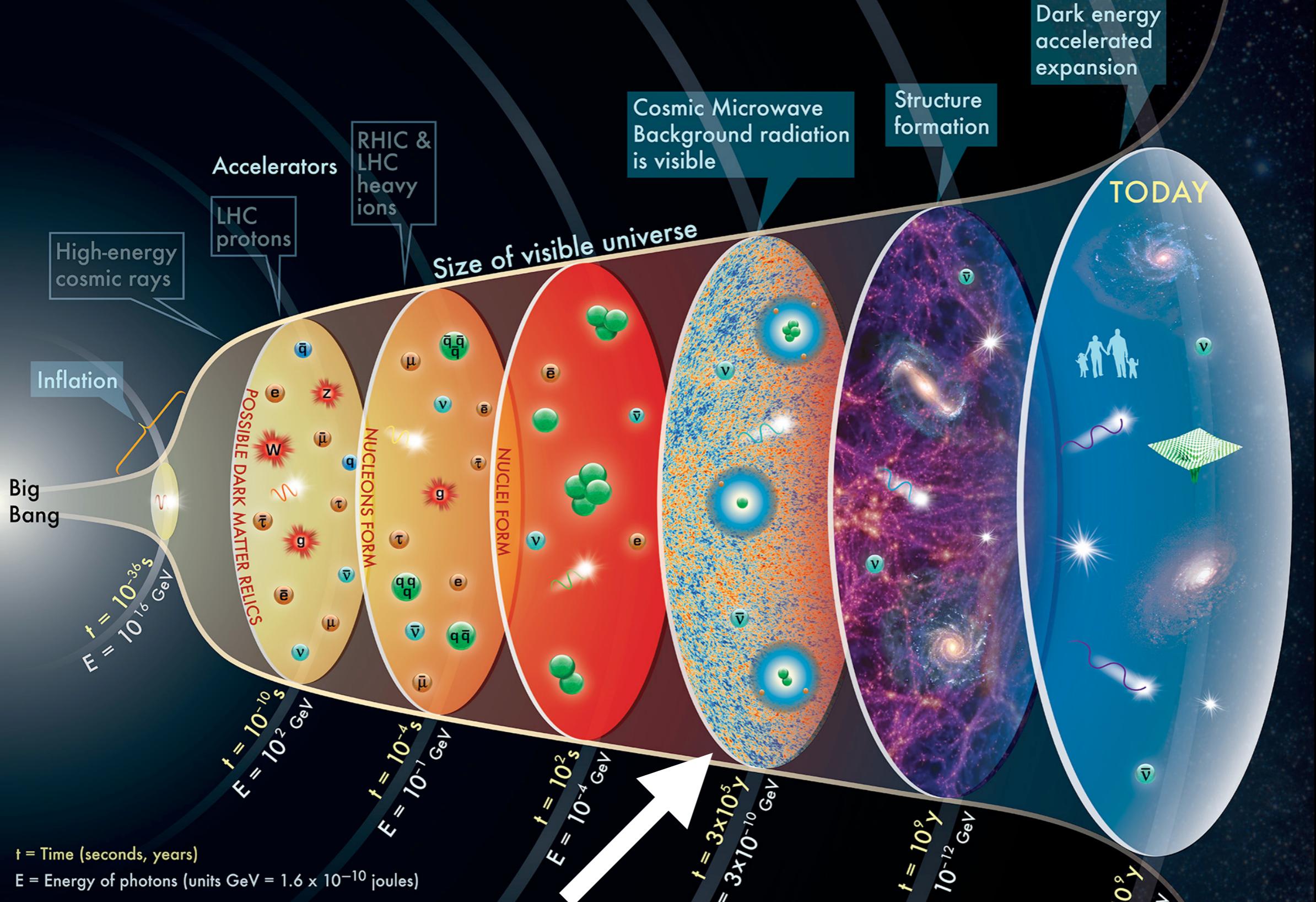
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- Neutrino masses and Cosmic Microwave Background Radiation?
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3. Is anyone cheating? Neutrinos and Tensions

4. Final score:

- Take home messages





Cosmic Background Radiation: N_{eff}

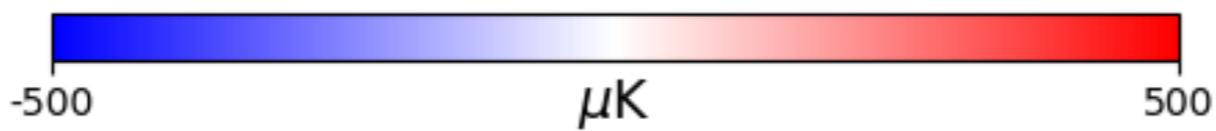
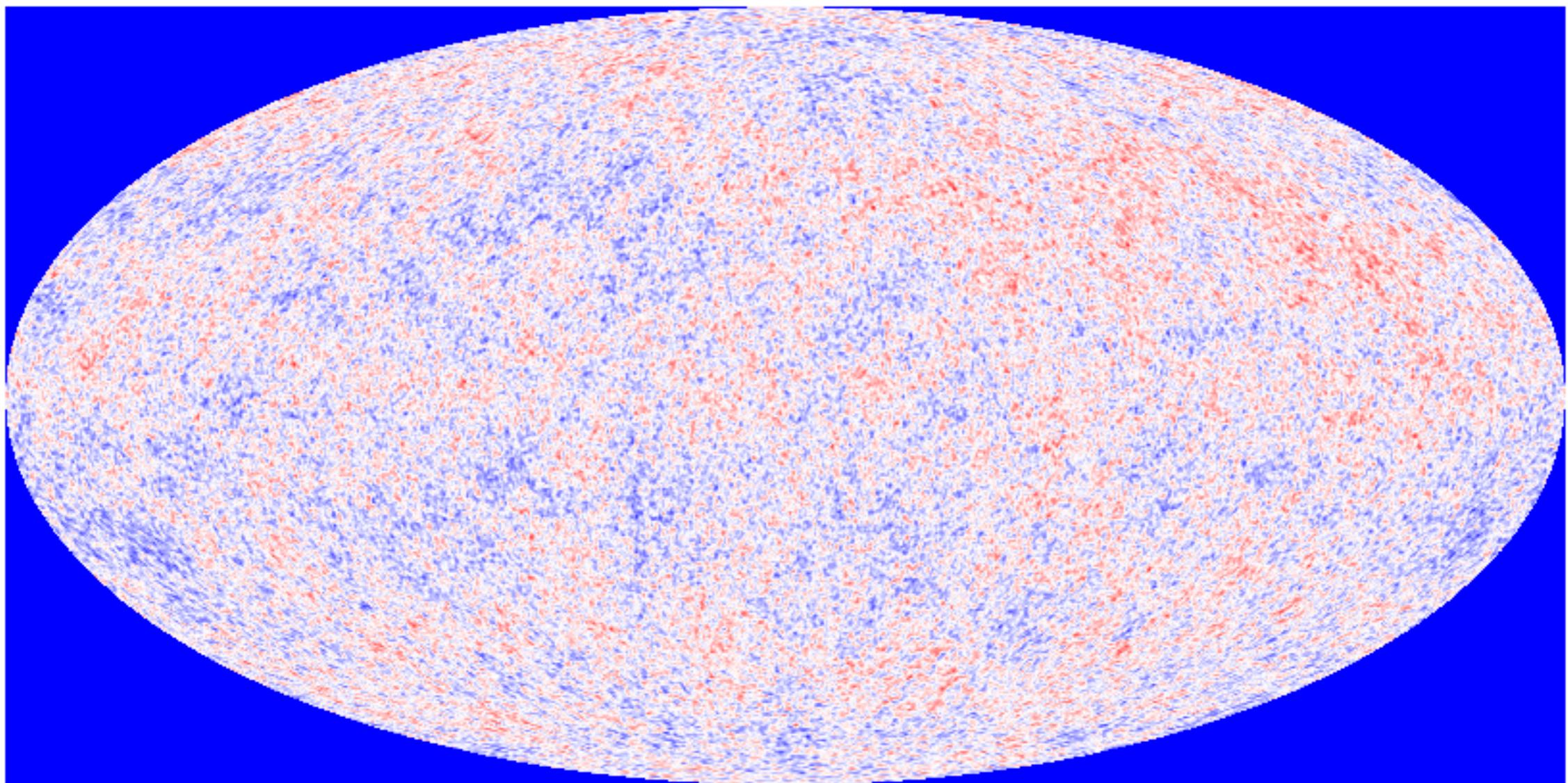
The concept for the above figure originated in a 1986 paper by Michael Turner.

Event	Time	Redshift	Temperature
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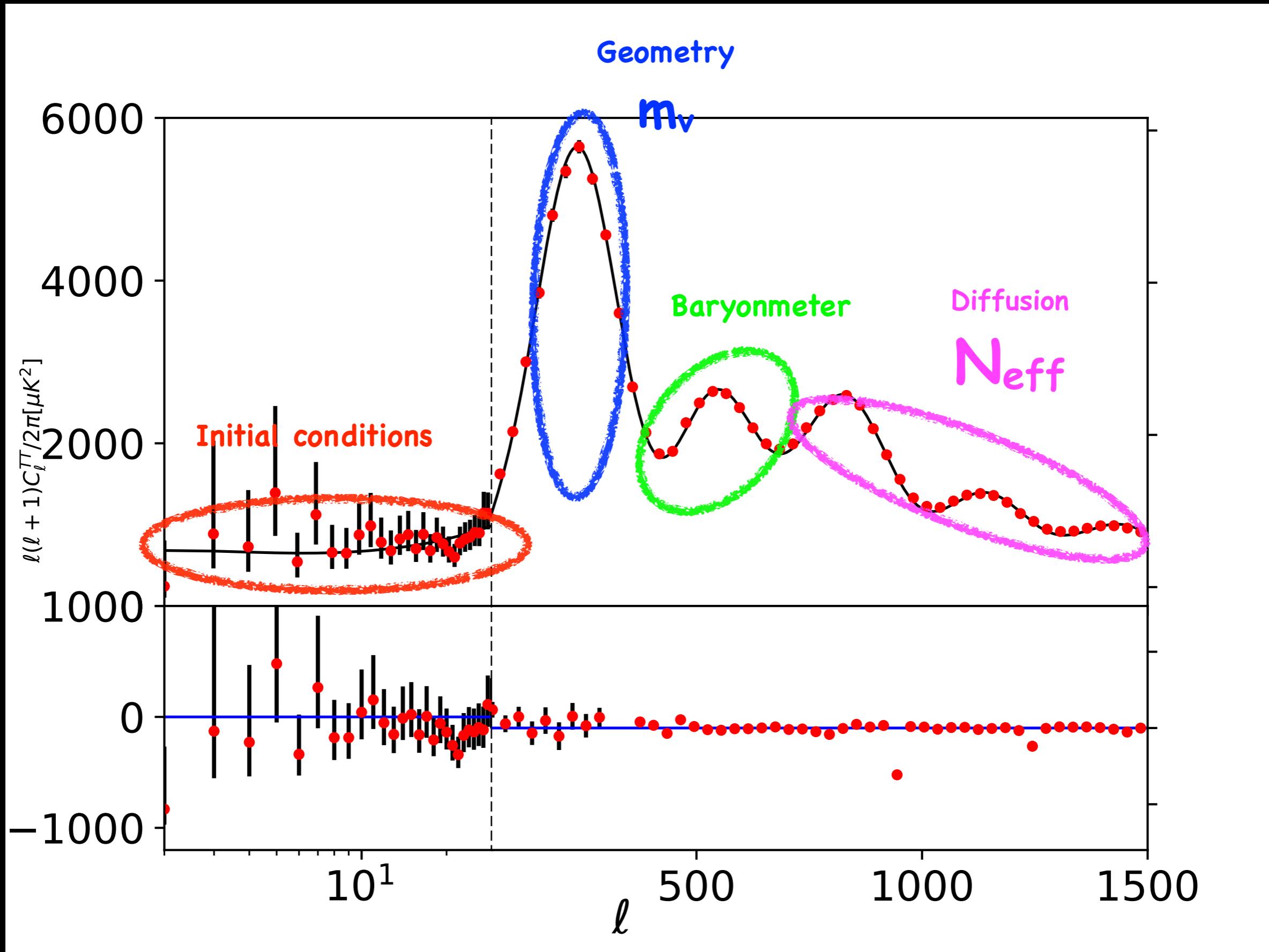
Also known as “photon decoupling”, as photons started freely travel through the universe without interacting with matter and the CMB is “frozen”

CMB: Neff

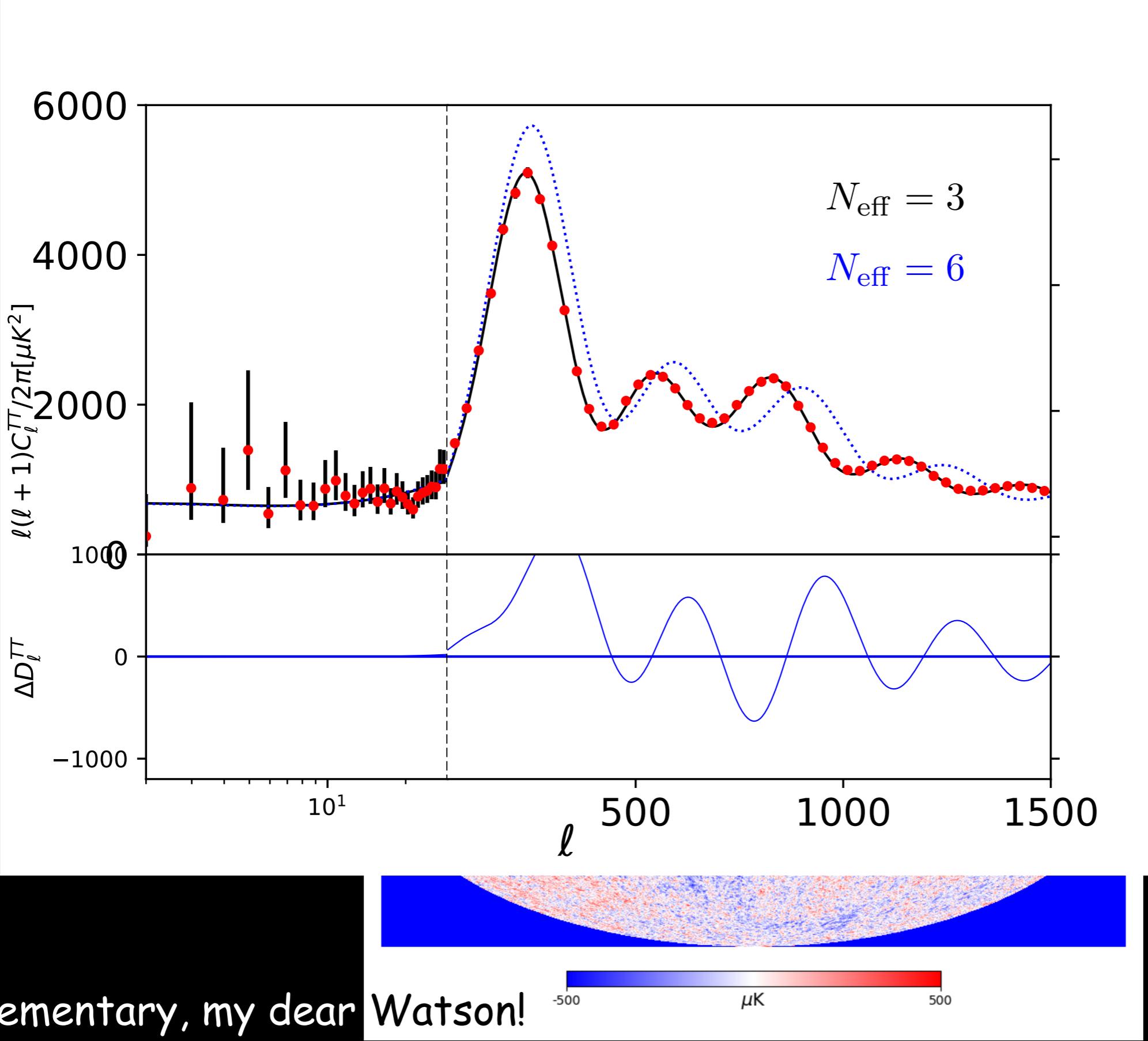
67.31_0.02222_0.1197_0.078_3.089_0.9655_3.046_0 map



CMB: a lot to learn about....

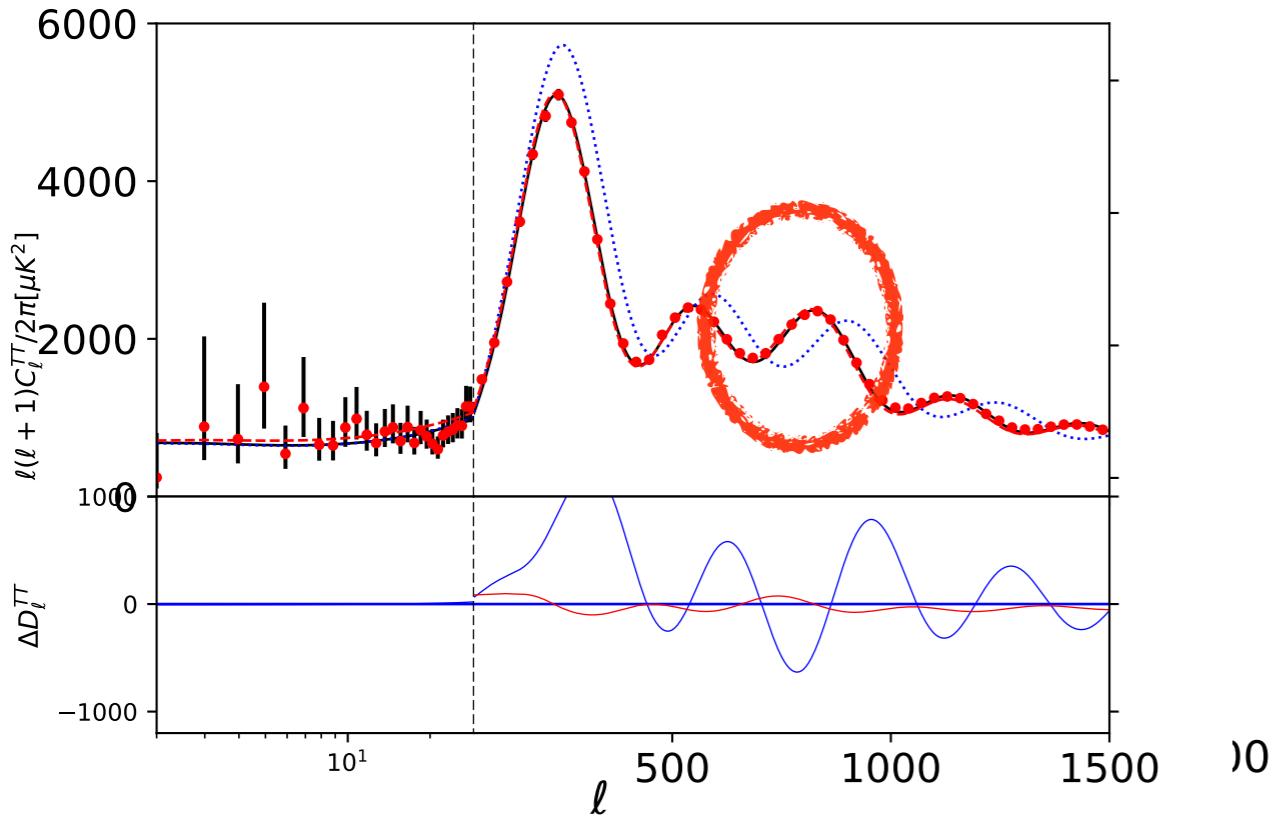


CMB: N_{eff}



CMB: N_{eff}

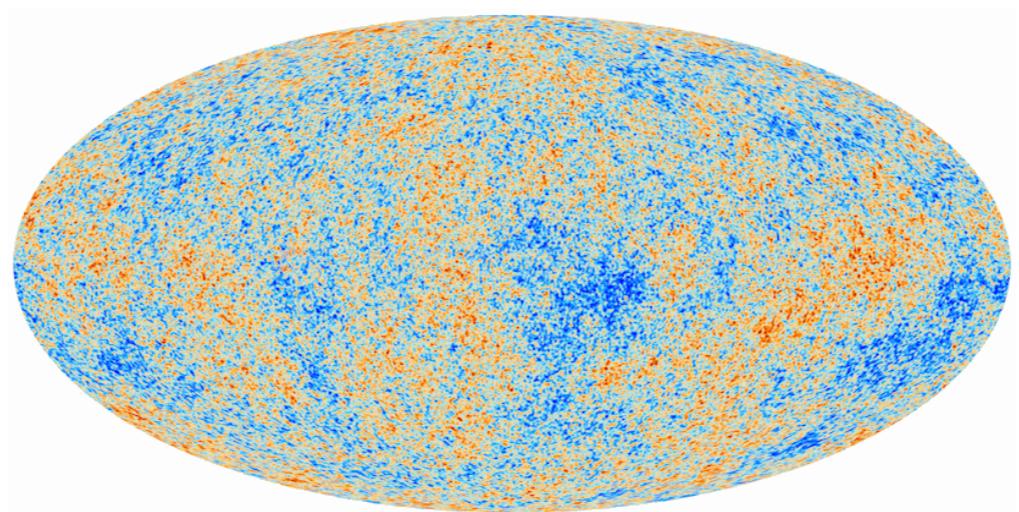
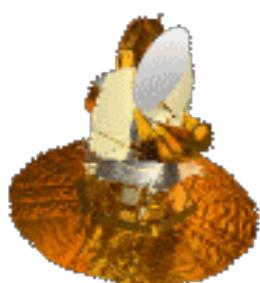
$$N_{\text{eff}} = 6 \quad N_{\text{eff}} = 3 \quad N_{\text{eff}} = 6$$



$(\omega_b, \omega_m, h, A_s, n_s, \tau, N_{\text{eff}})$
Warning!

It is elementary, Sherlock Holmes!

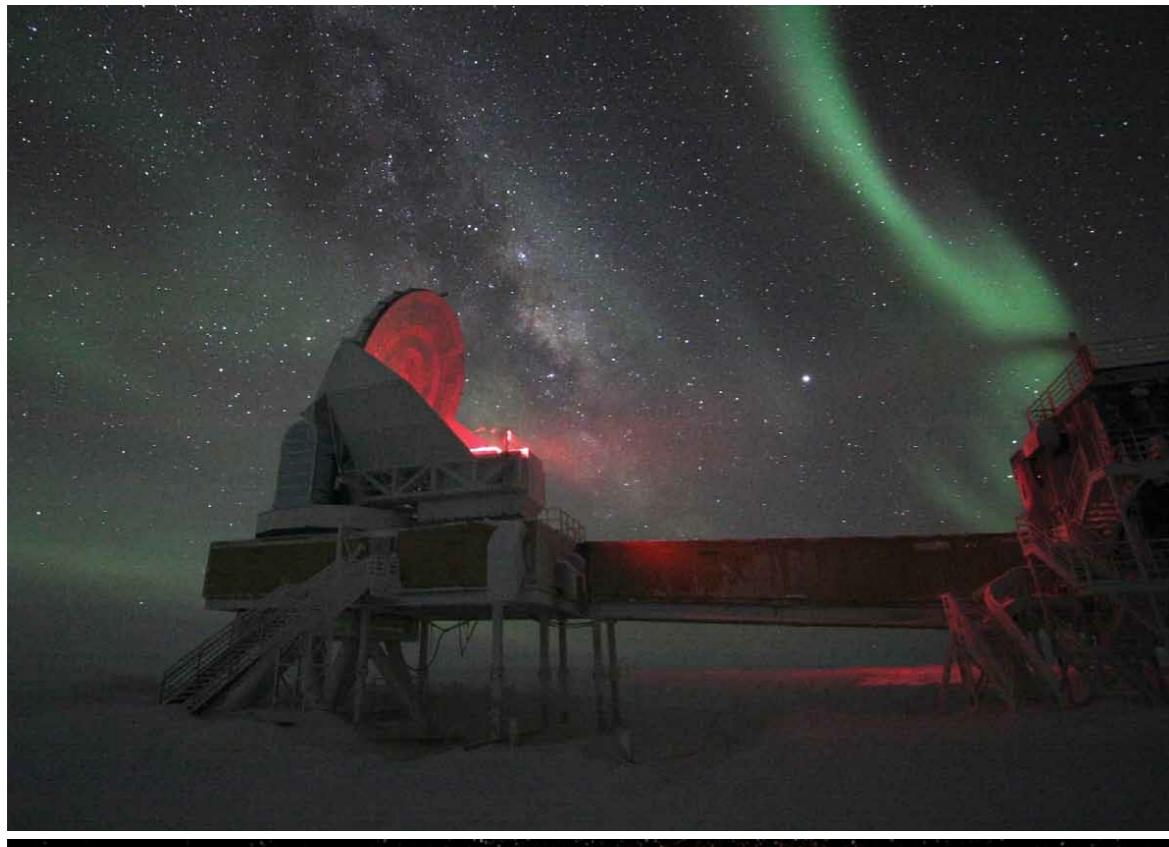
Only effect at $l < 1000$ that can not be mimicked by others: anisotropic stress, around 3rd peak



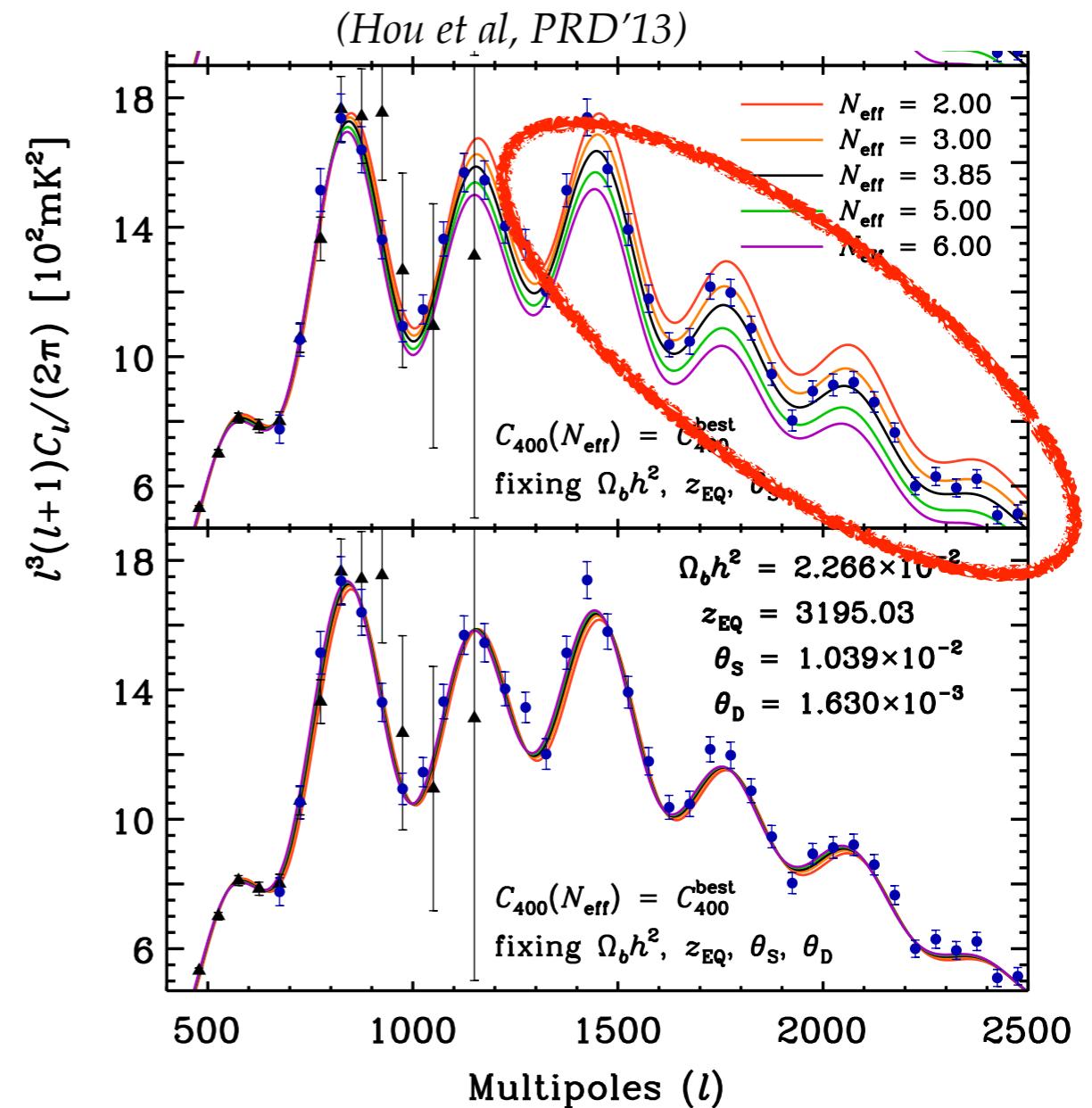
Neutrinos are free-streaming particles propagating at the speed of light, faster than the sound speed in the photon fluid, suppressing the oscillation amplitude of CMB modes that entered the horizon in the radiation epoch.

CMB: N_{eff}

@Cosmic Microwave Background in the damping tail,
 measured by SPT, ACT & Planck:
 Higher N_{eff} will increase the expansion rate AND
 the damping at high multipoles.



$$r_d^2 \propto \int_0^{a_*} \frac{da}{a^3 \sigma_T n_e H}$$



N_{eff}

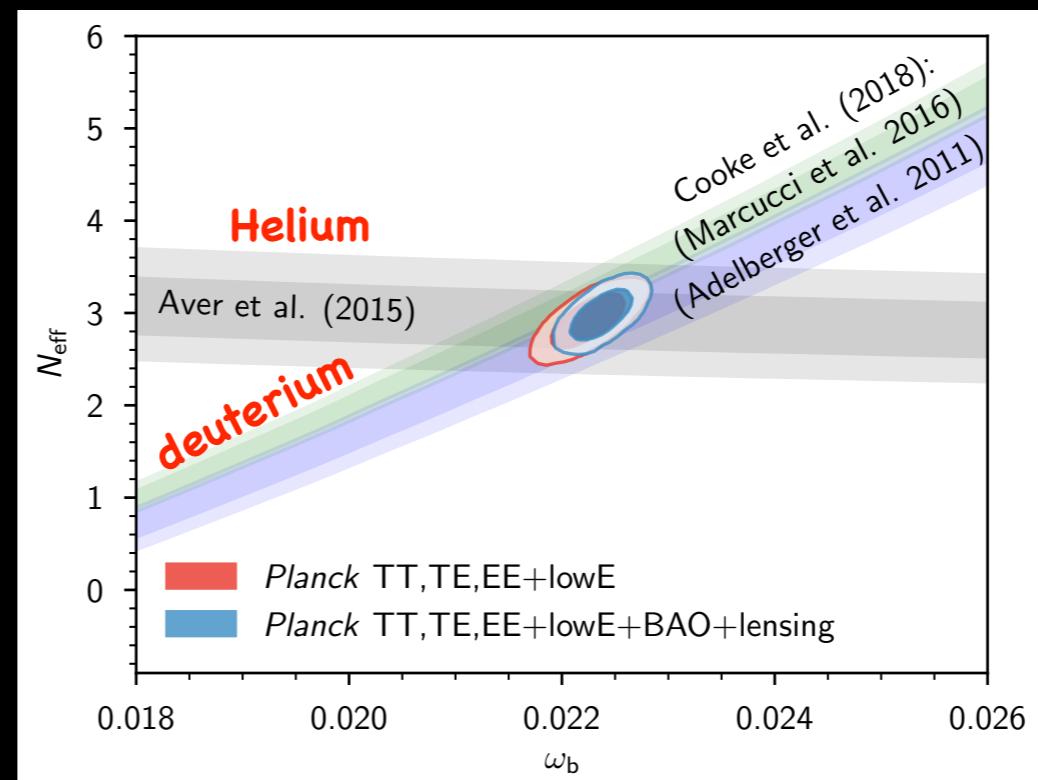
- Planck 2018 CMB temperature polarization and lensing potential data:

$$N_{\text{eff}} = 2.89^{+0.36}_{-0.38} \text{ 95%CL}$$

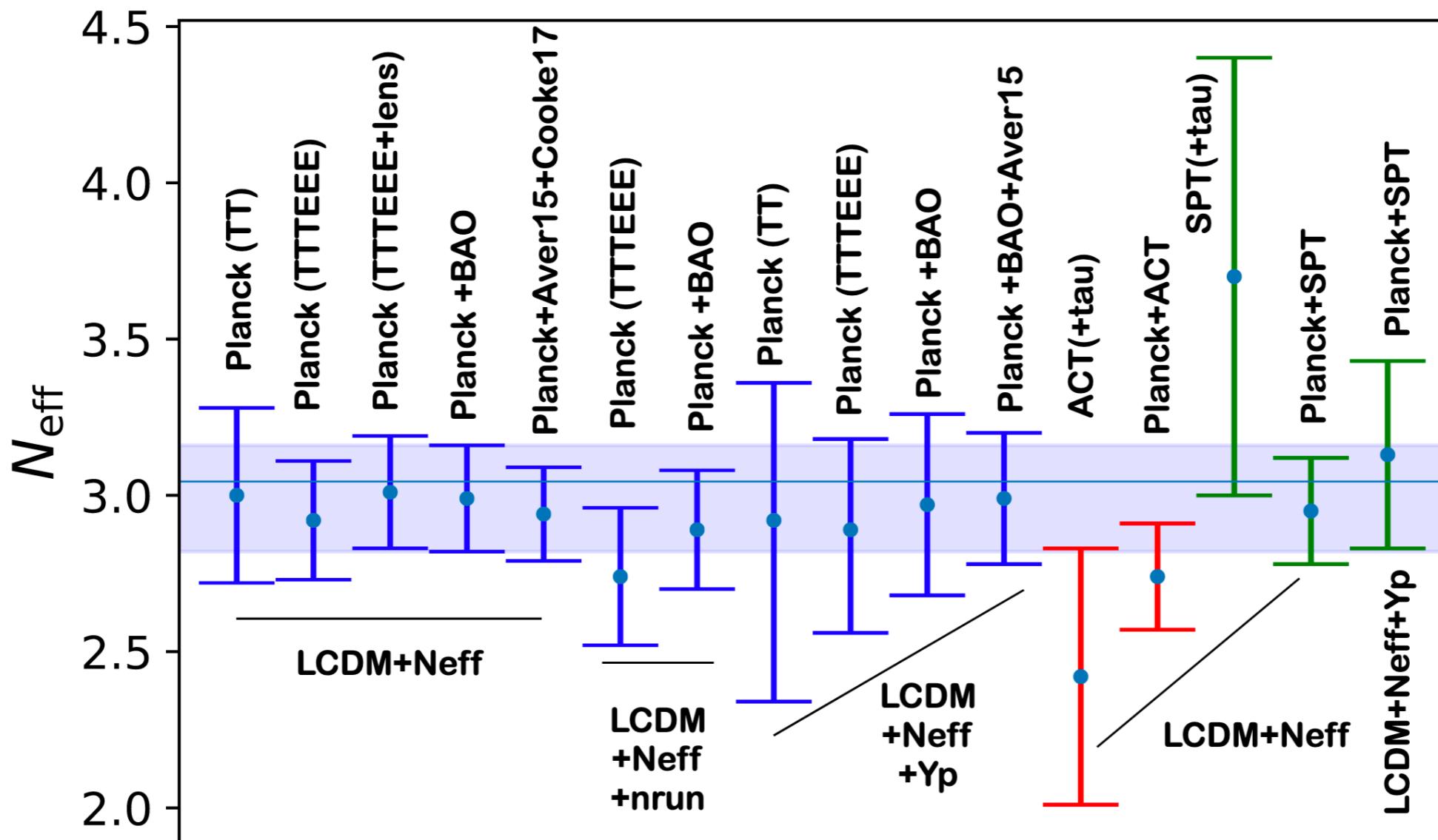
- If we add large scale structure information in the BAO shape form:

$$N_{\text{eff}} = 2.99^{+0.34}_{-0.33} \text{ 95%CL}$$

- Perfectly consistent with BBN estimates:



Neff: current status



Planck collaboration, VI 2018
 ACT Collaboration (Aiola+), 2020
 SPT Collaboration (Dutcher+, Balkenhol+), 2021



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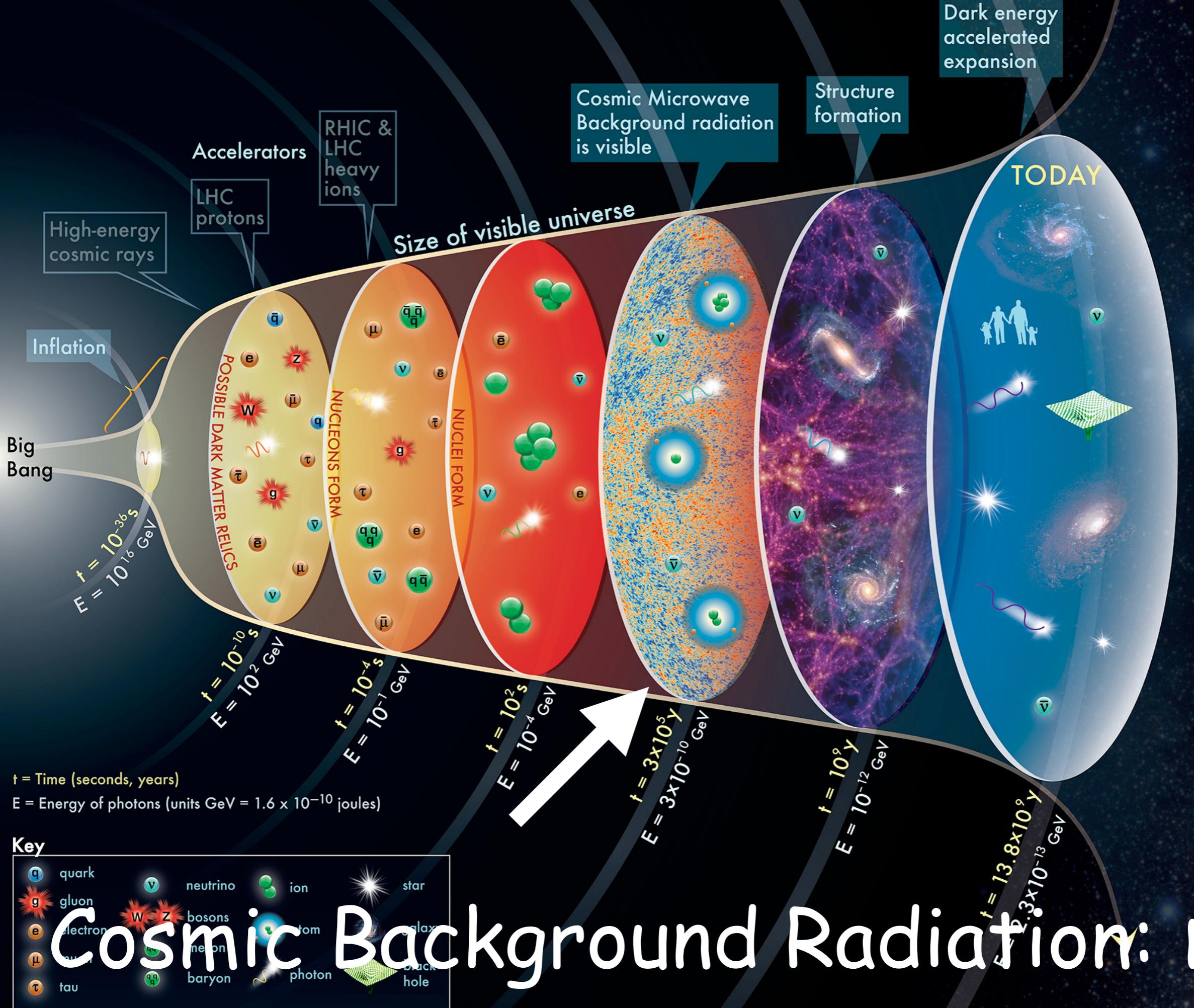
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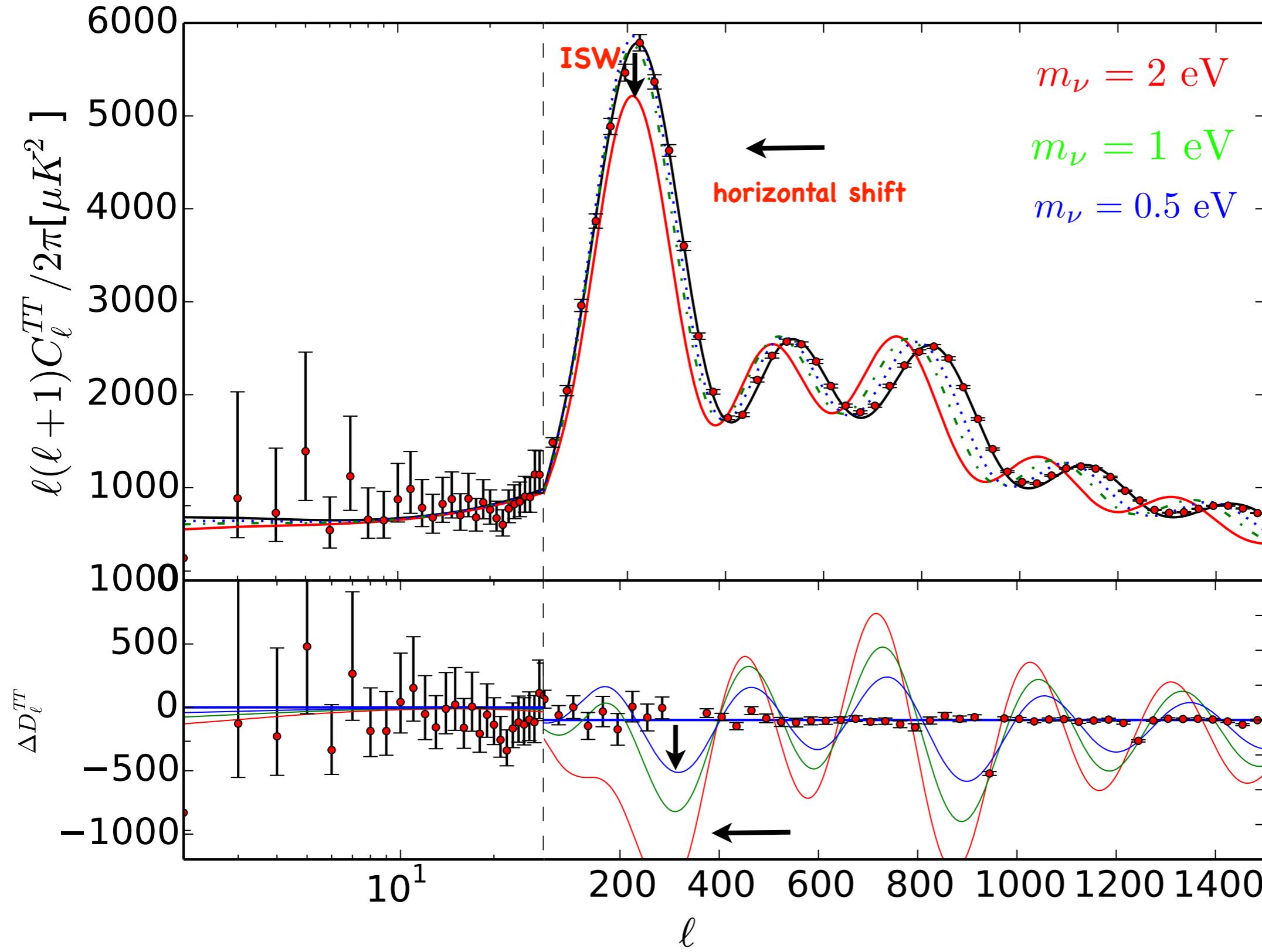
- Take home messages



CMB: $\sum m_\nu$

@ CMB: Early Integrated Sachs Wolfe effect (ISW).

Shift in the angular position of the peaks.



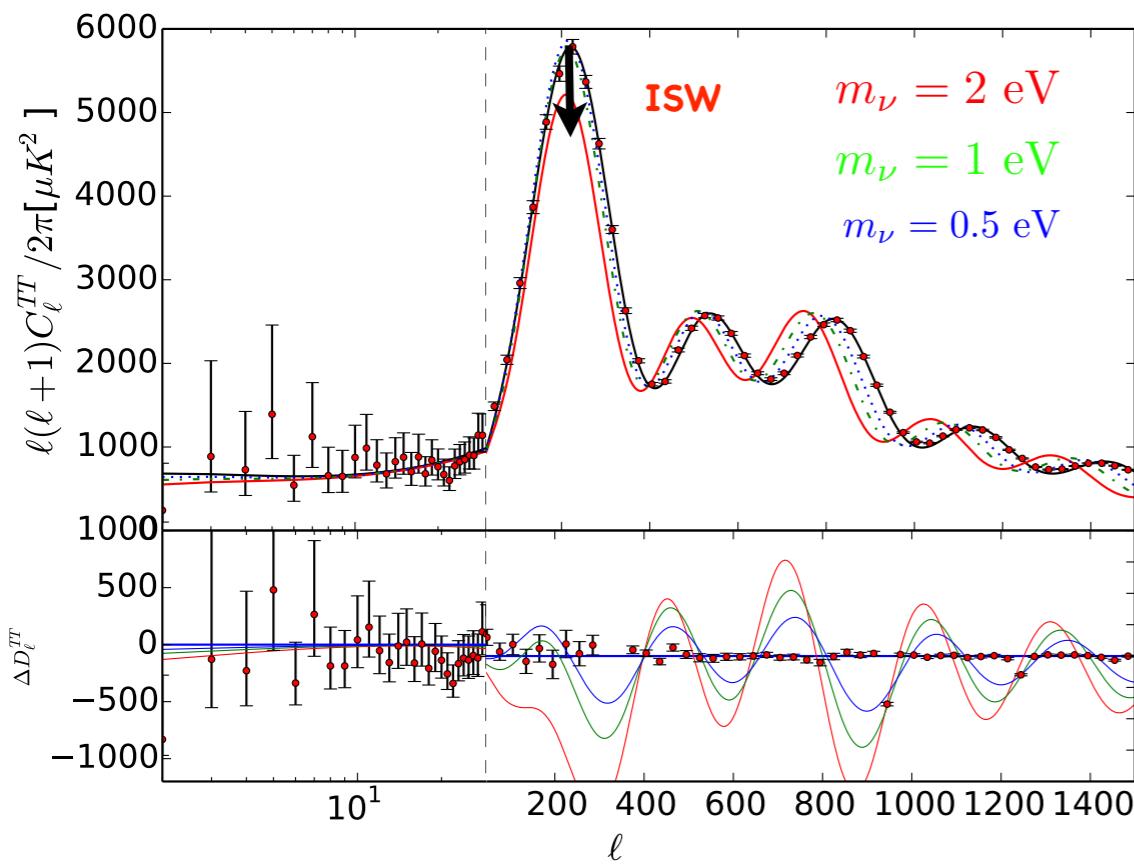
CMB: Σm_ν

@ CMB: Early Integrated Sachs Wolfe effect (ISW)

$$\Theta(\hat{n}) = \frac{\delta T}{T}(\hat{n}) \simeq \Theta_0 + \Psi + \hat{n}(\hat{v}_e - v) + \int \dot{\Psi} + \dot{\Phi} d\eta$$

In matter domination, the gravitational potential is constant: NO ISW effect!

The transition from the relativistic to the non relativistic neutrino regime gets imprinted in the decays of the gravitational potentials near the recombination period, contributing to the ISW effect!



This early ISW effect leads to a depletion of:

$$\frac{\Delta C_\ell}{C_\ell} = -\left(\sum m_\nu / 0.1 \text{ eV}\right)\%$$

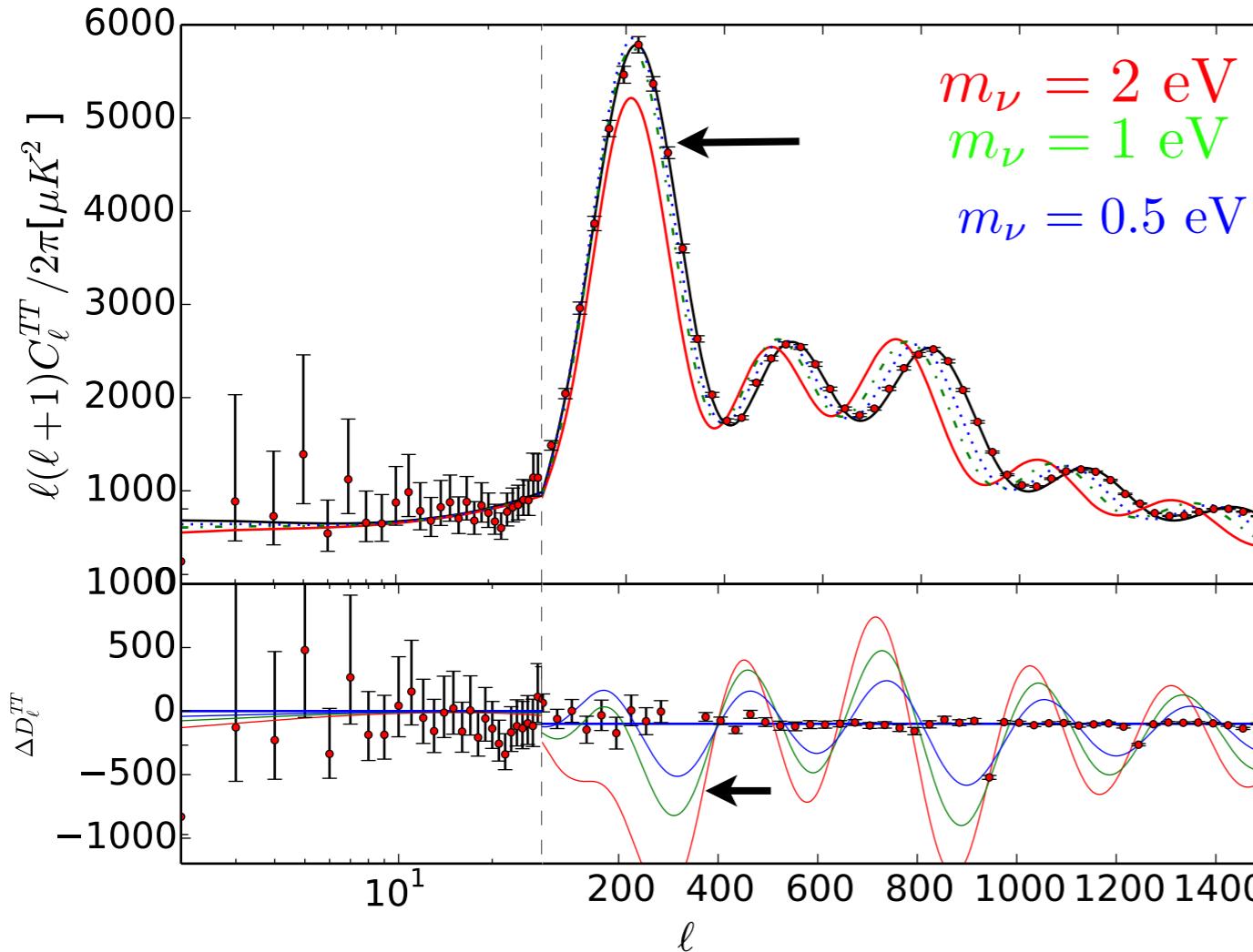
on multipoles:

$$20 < \ell < 200$$

CMB: Σm_ν

@ CMB: Early Integrated Sachs Wolfe effect (ISW).

Shift in the angular position of the peaks.



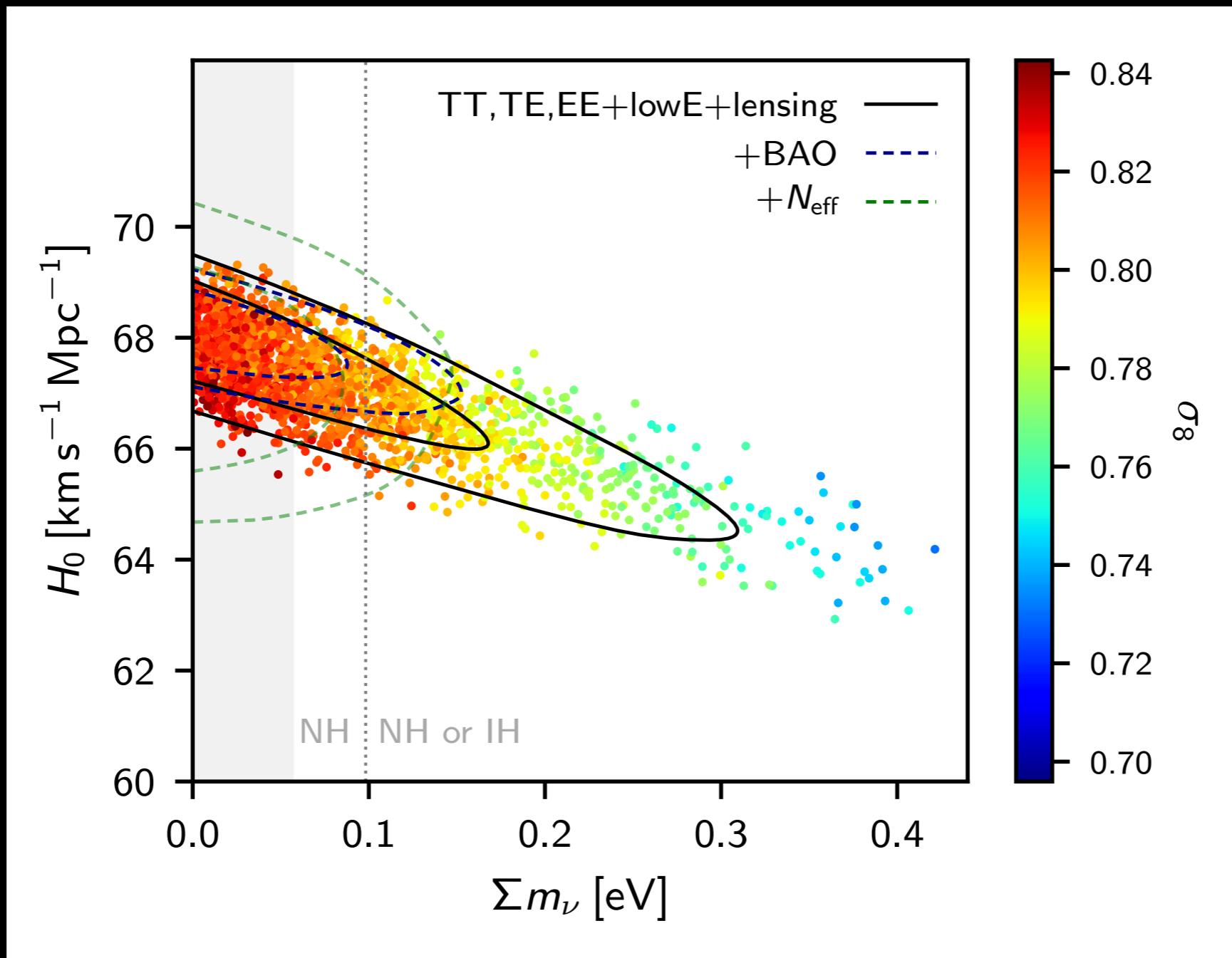
$$\theta_s = \frac{r_s}{D_A}$$

$$r_s = \int_0^{t(z_d)} c_s (1+z) dt = \frac{2}{3k_{\text{eq}}} \sqrt{\frac{6}{R_{\text{eq}}}} \ln \frac{\sqrt{1+R_d} + \sqrt{R_d + R_{\text{eq}}}}{1 + \sqrt{R_{\text{eq}}}}$$

$$D_A = \int_0^{z_{\text{rec}}} \frac{dz}{H(z)}$$

The higher the neutrino mass, the lower the angular diameter distance.
 Peaks shift to lower multipoles. But this effect can be compensated with a lower Hubble constant:

Strong degeneracy between Σm_ν and the Hubble constant H_0 !



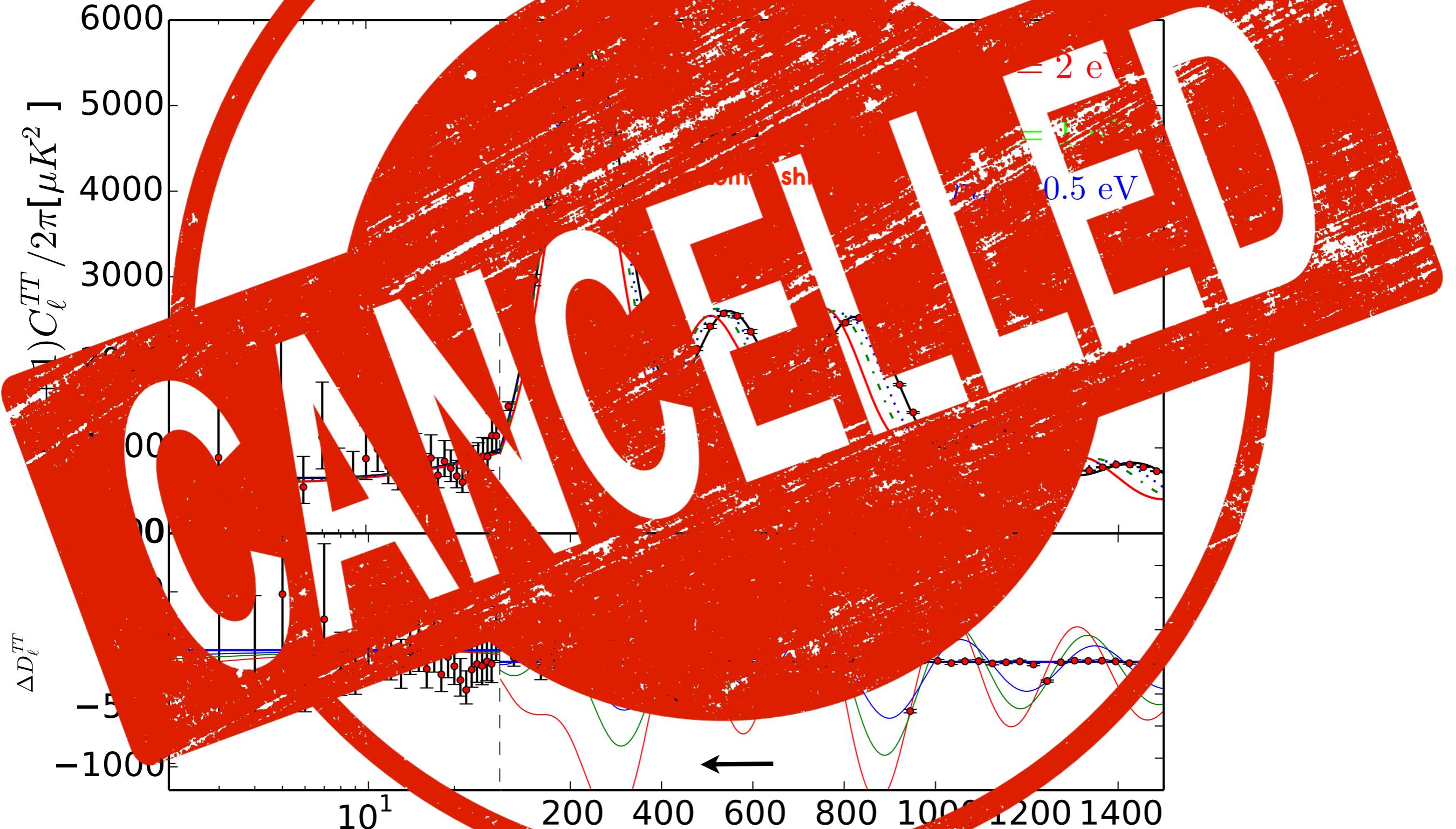
Planck Coll. A&A'20

Strong degeneracy between Σm_ν and the Hubble constant H_0 !

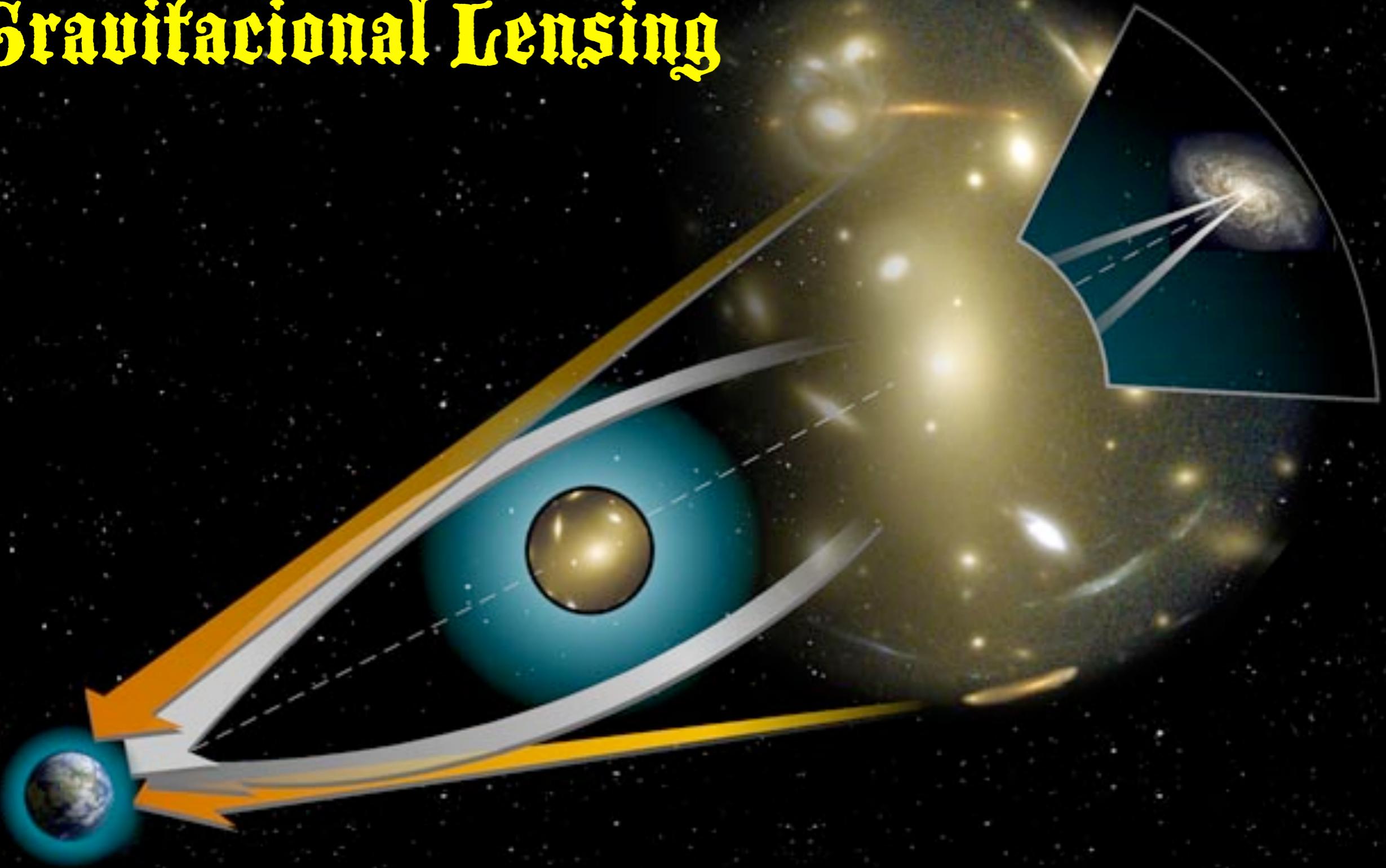
CMB: Σm_ν

@ CMB: Early Integrated Sachs Wolfe effect (ISW).

Shift in the angular position of the peaks.



Gravitacional Lensing



Einstein's relativity predicts that the presence of a massive body will curve space time, distorting the light trajectory. The shape of the background objects will change/ multiplied by the presence of intervening galaxies.

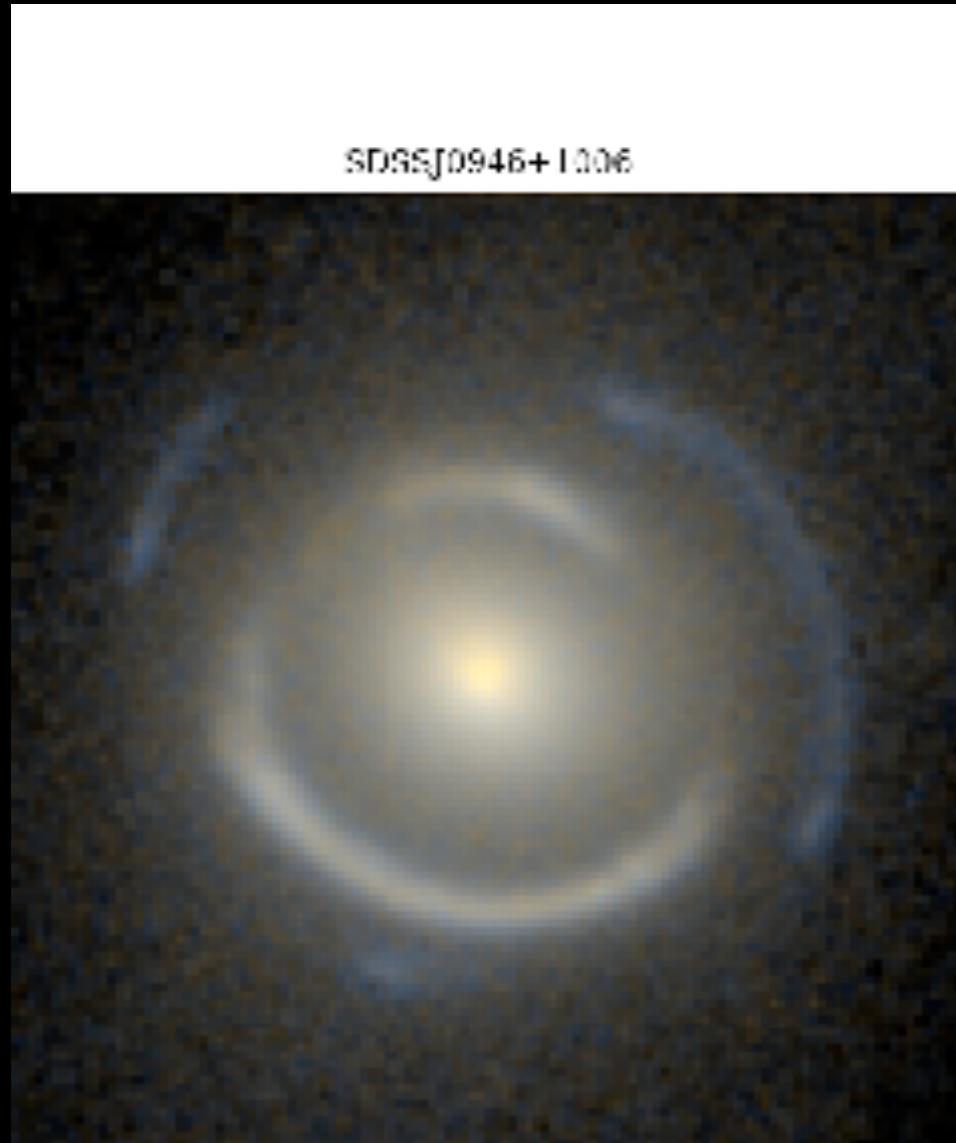
Einstein rings: Perfect alignment: Syzygy!

Lensing Galaxy



This movie shows a spiral galaxy acting as a lens of a background quasar (Quasi-stellar radio source) moving behind the galaxy. When the alignment source-lens-observer is perfect, we see the formation of the Einstein ring!

Gravitacional Lensing



Double Einstein ring! 3 perfectly aligned galaxies (probably less than 100 cases in all the universe, and we have seen one!)

CMB Lensing: Σm_ν

Lensing remaps the CMB fluctuations:

$$\Theta_{\text{lensed}}(\hat{n}) = \Theta(\hat{n} + \nabla\phi(\hat{n}))$$

Lensing potential ϕ is a measure of the integrated mass distribution back to the last scattering surface

$$\phi(\hat{n}) = -2 \int_0^{z_{\text{rec}}} \frac{dz}{H(z)} \boxed{\Psi(z, D(z)\hat{n})} \left(\frac{D(z_{\text{rec}}) - D(z)}{D(z_{\text{rec}})D(z)} \right)$$

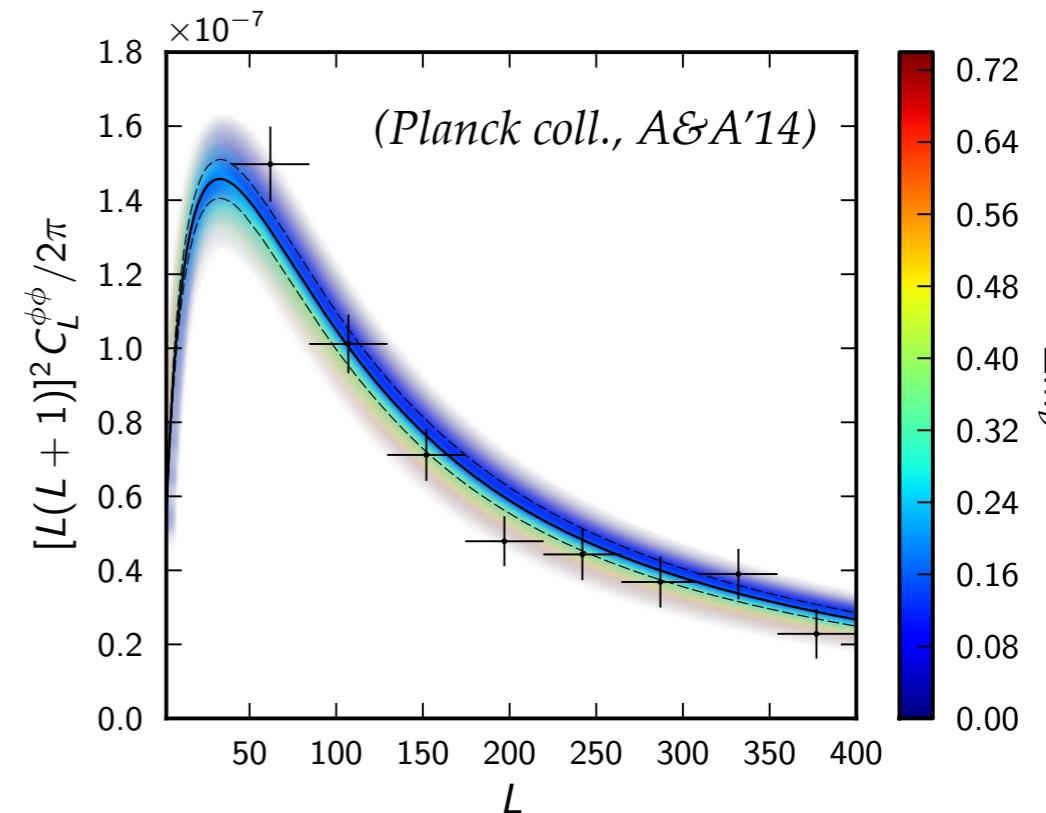
Matter distribution

Geometry

$$C_L^{\phi\phi} = \frac{8\pi^2}{L^3} \int_0^{z_{\text{rec}}} \frac{dz}{H(z)} D(z) \left(\frac{D(z_{\text{rec}}) - D(z)}{D(z_{\text{rec}})D(z)} \right)^2 P_\Psi(z, k = L/D(z))$$

Neutrinos are hot relics with large thermal velocities, implying less clustering on small scales, reducing therefore CMB lensing!

(Kaplinghat et al PRL'03, Lesgourges et al, PRD'06)



CMB: $\sum m_\nu$

Planck TTTEEE+lowT+lowE+lensing

Planck Coll. A&A'20

$$\sum m_\nu < 0.24 \text{ eV } 95\% \text{CL}$$

Cosmic neuTRIVIA game steps

1. Familiarize yourself with the board's layout:

- The Λ CDM trivia: the players!
- The neutrino pie piece: decoupling in the early universe

2. Roll the dice and get:

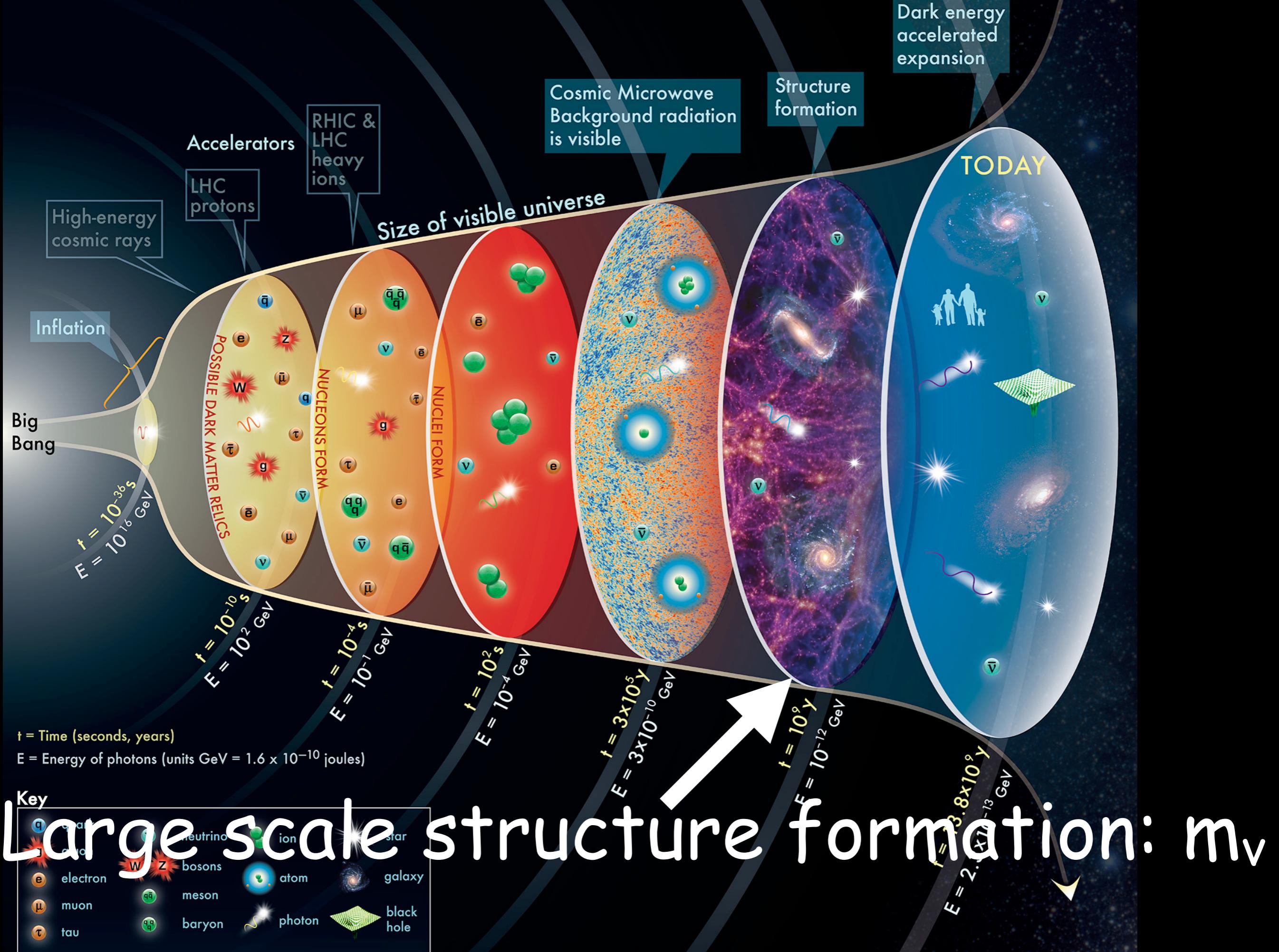
- Number of neutrinos and Big-Bang Nucleosynthesis
- Number of neutrinos and Cosmic Microwave Background Radiation
- Neutrino masses and Cosmic Microwave Background Radiation?
- Neutrino masses and structure formation in the universe?

3. Is anyone cheating? Neutrinos and Tensions

4. Final score:

- Take home messages



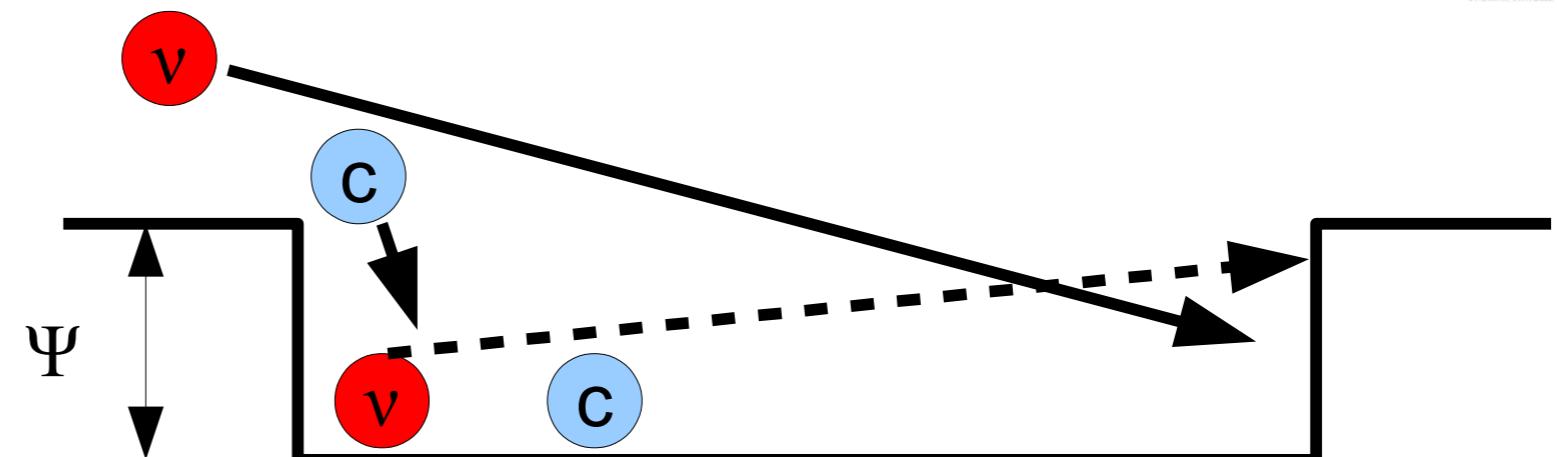
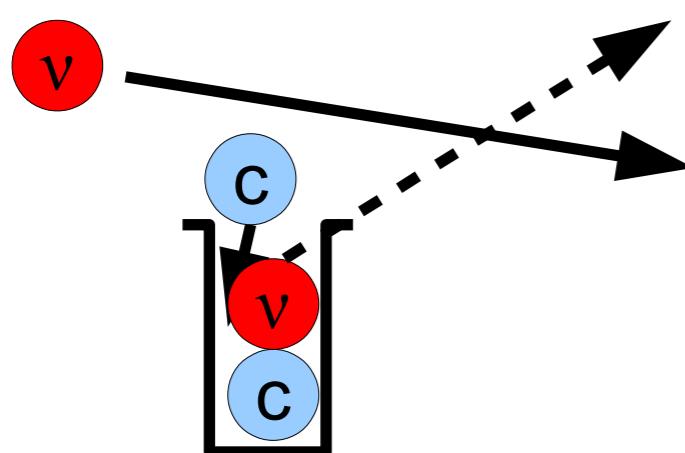


Large scale structure: m_ν

Neutrino masses suppress structure formation on scales larger than their free streaming scale when they turn non relativistic. (Bond et al PRL'80)

Neutrinos with eV or sub-eV masses are HOT relics with LARGE thermal velocities!

Cold dark matter instead has zero velocity and therefore it clusters at any scale!



$$\lambda \ll \lambda_{fs,\nu} \rightarrow k \gg k_{fs,\nu}$$

$$\lambda \gg \lambda_{fs,\nu} \rightarrow k \ll k_{fs,\nu}$$



Large scale structure: m_V

Growth equation for a single uncoupled fluid, linear regime, with constant sound speed:

$$\ddot{\delta} + \boxed{2\frac{\dot{a}}{a}\dot{\delta}} - \boxed{c_s^2 k^2 \frac{\delta}{a^2}} = \boxed{4\pi G \rho \delta}$$

Pressure

Jeans scale:

Neutrino free streaming scale:

$$k_J \equiv \sqrt{\frac{4\pi G\rho}{c_s^2(1+z)^2}}$$

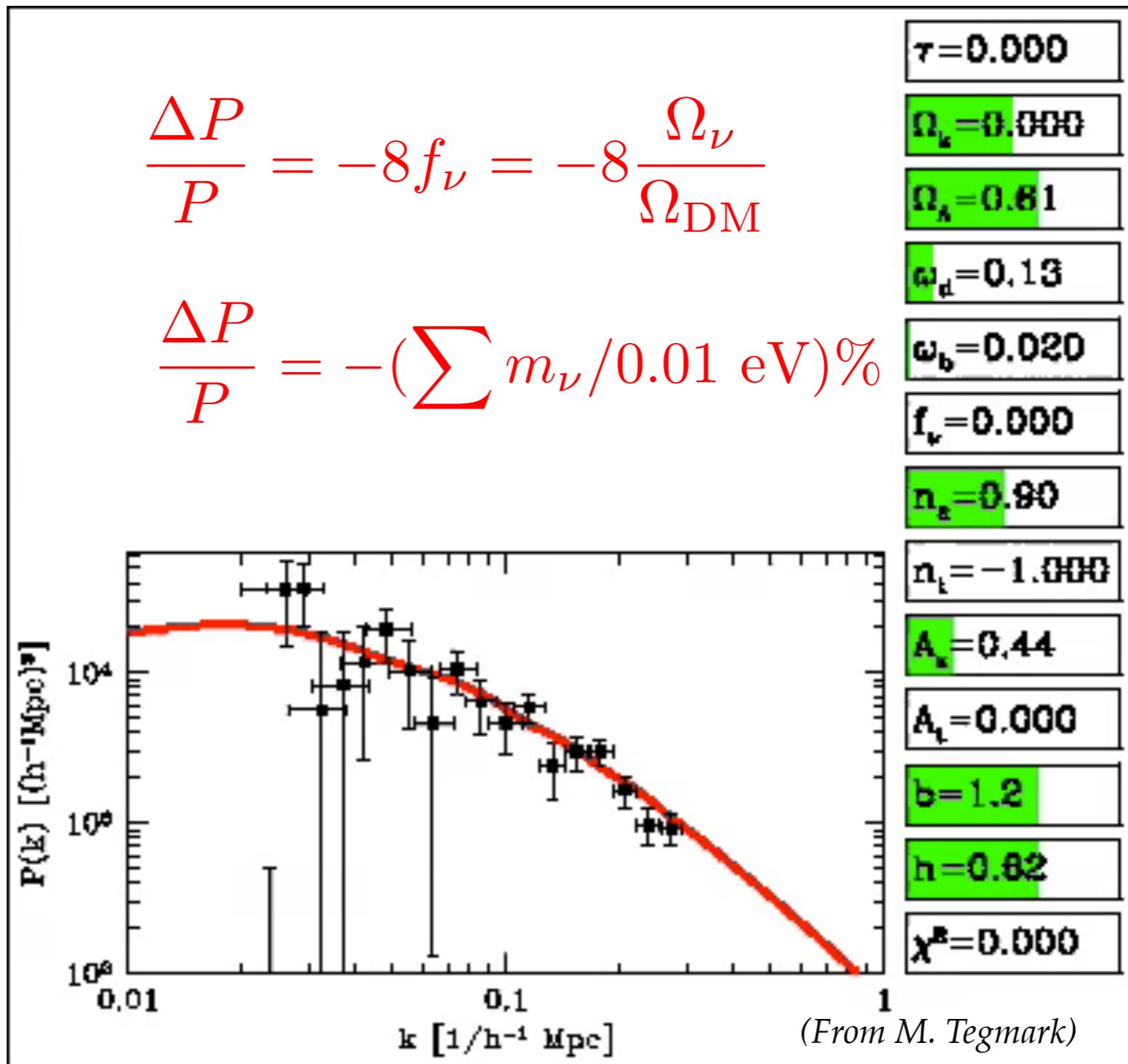
$k > k_J$ no growth can occur

$k < k_J$ density perturbations growth

$$k_{fs,\nu}(z) \equiv \sqrt{\frac{3}{2}} \frac{H(z)}{(1+z)\sigma_{v,\nu}(z)}$$

Large scale structure: m_ν

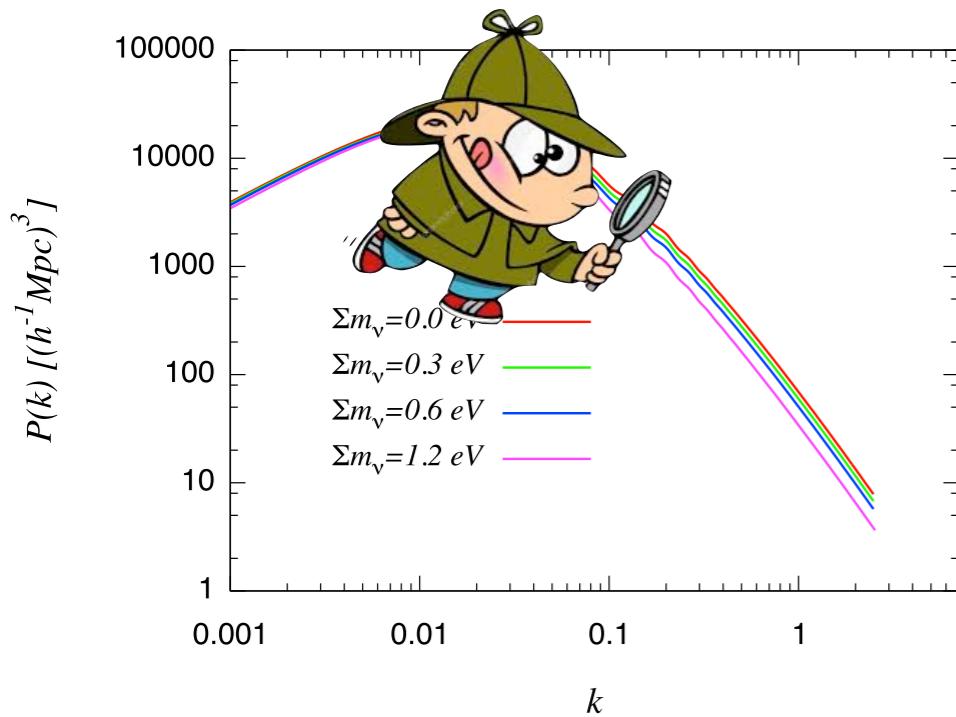
Matter power spectrum suppression:



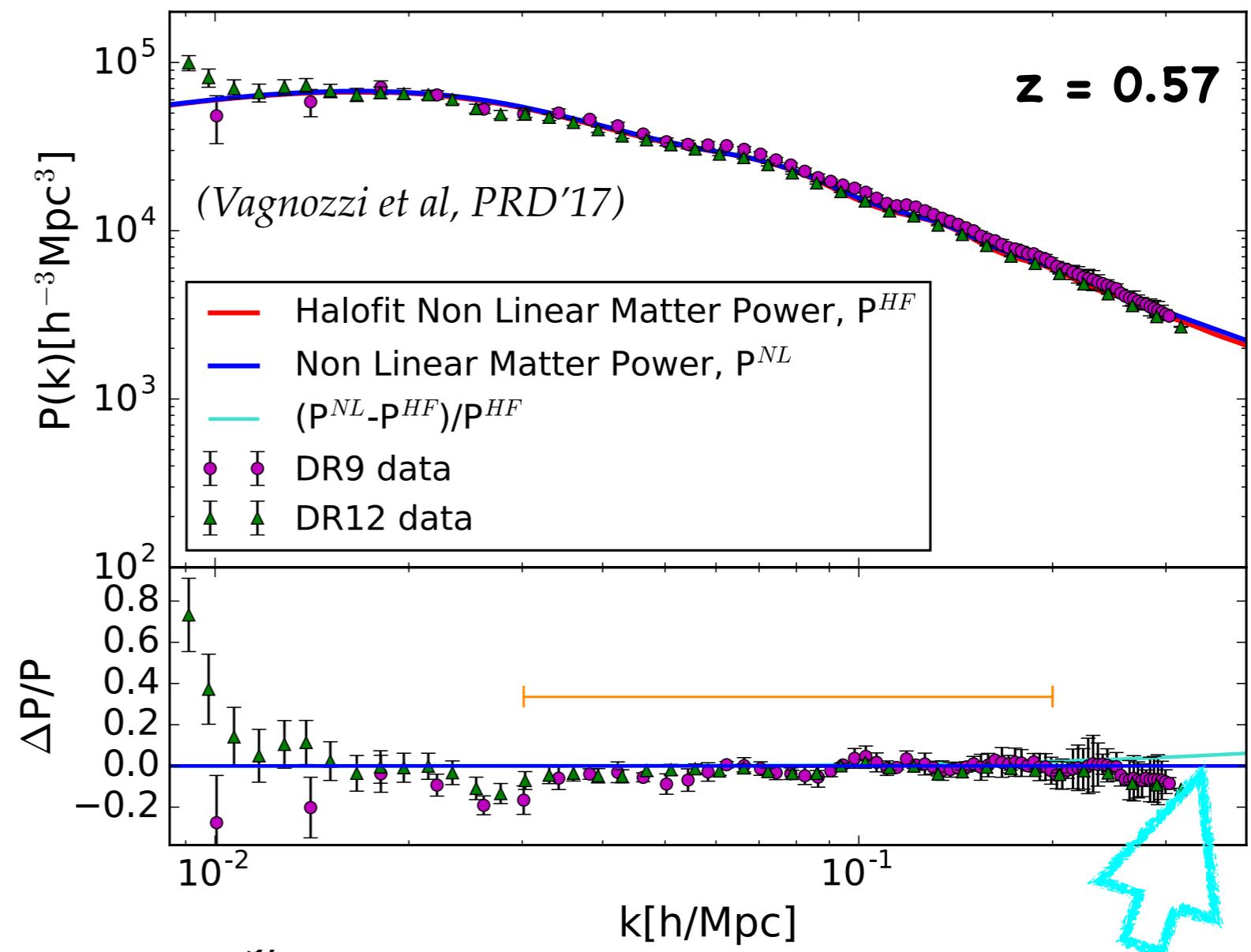
↑ Small scales
56

Large scale structure: m_ν

@LSS: Caveats, NON-LINEARITIES



Beyond a given scale k_{nl} , linear perturbation theory breaks down!



Large scale structure: m_V

@LSS: Caveats, BIAS!

$$P_{gg}(k, z) = \text{bias}^2 P(k, z)$$

Galaxies are **biased** tracers of the underlying matter density field! (Kaiser, APJ'84)

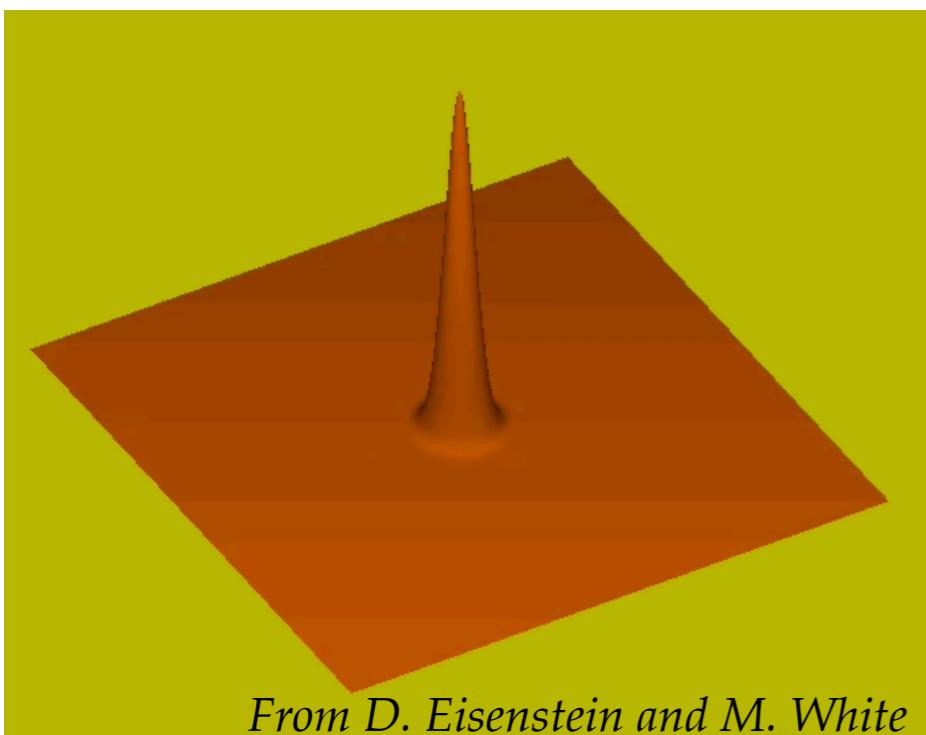
Neutrinos themselves induce a scale-dependent bias (LoVerde & Zaldarriaga; Castorina et al)

Baryon Acoustic Oscillations

Photons and baryons in the early universe behave as a tightly coupled fluid, resembling acoustic waves, generated as the baryon-photon fluid is attracted and falls onto the overdensities:

$$\ddot{\delta} + [\text{Pressure} - \text{Gravity}] \delta = 0 \quad \delta(\vec{x}) \equiv \frac{\rho(\vec{x}) - \bar{\rho}(\vec{x})}{\bar{\rho}(\vec{x})}$$

The time when the baryons are “released” from the drag of the photons is known as the drag epoch. From then on photons expand freely while the acoustic waves “freeze in” the baryons at a scale given by the size of the horizon at the drag epoch:



$$R \equiv 3\rho_b/4\rho_\gamma$$

$$r_s = \int_0^{t(z_d)} c_s (1+z) dt = \frac{2}{3k_{\text{eq}}} \sqrt{\frac{6}{R_{\text{eq}}}} \ln \frac{\sqrt{1+R_d} + \sqrt{R_d + R_{\text{eq}}}}{1 + \sqrt{R_{\text{eq}}}}$$

$$r_s = 147.09 \pm 0.26 \text{ Mpc}$$

Planck Coll. A&A'20

Baryon Acoustic Oscillations

Photons and baryons in the early universe behave as a tightly coupled fluid, resembling acoustic waves, generated as the baryon-photon fluid is attracted and falls onto the overdensities:

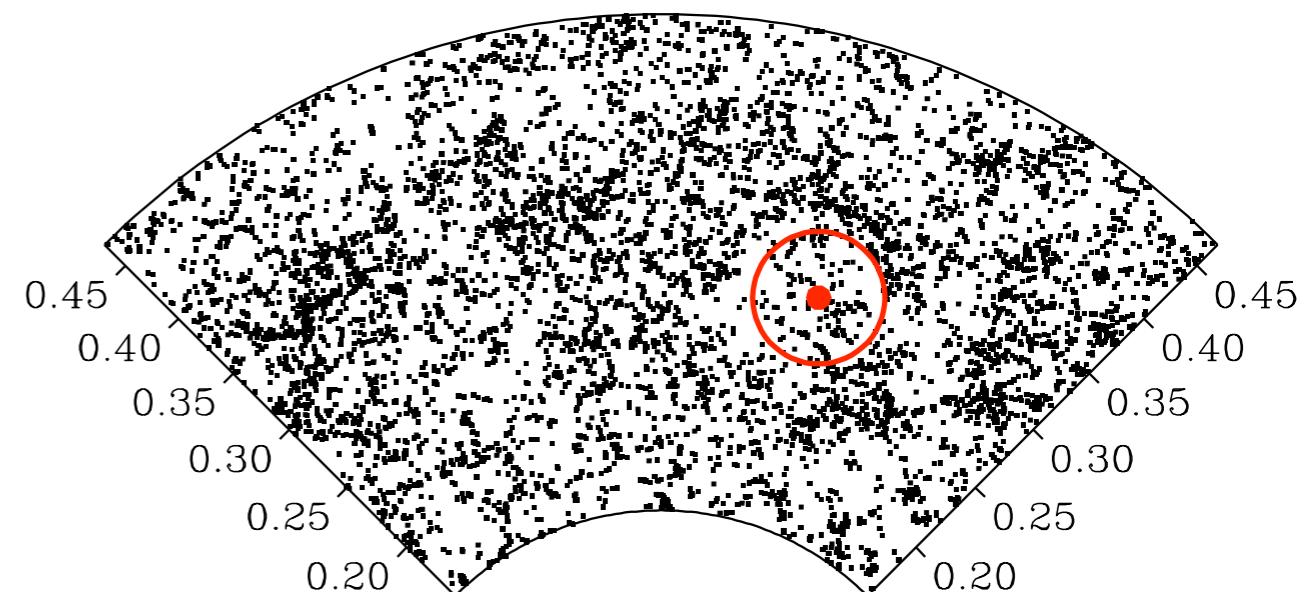
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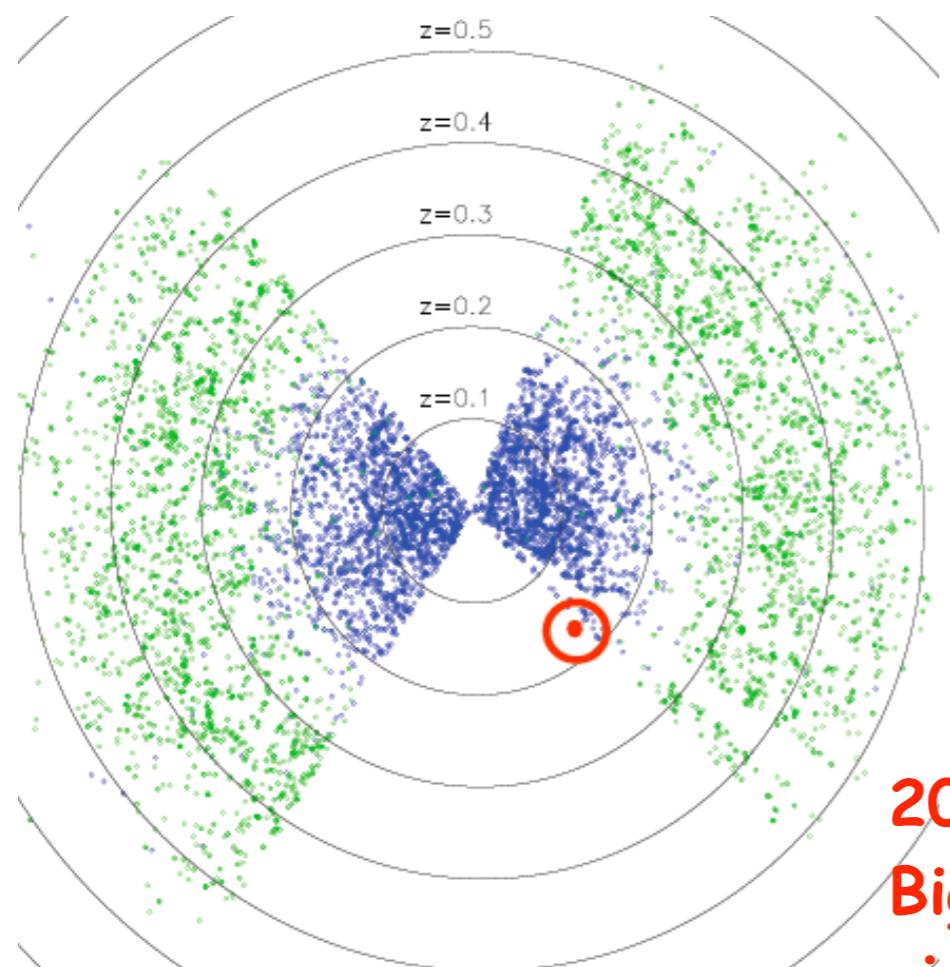
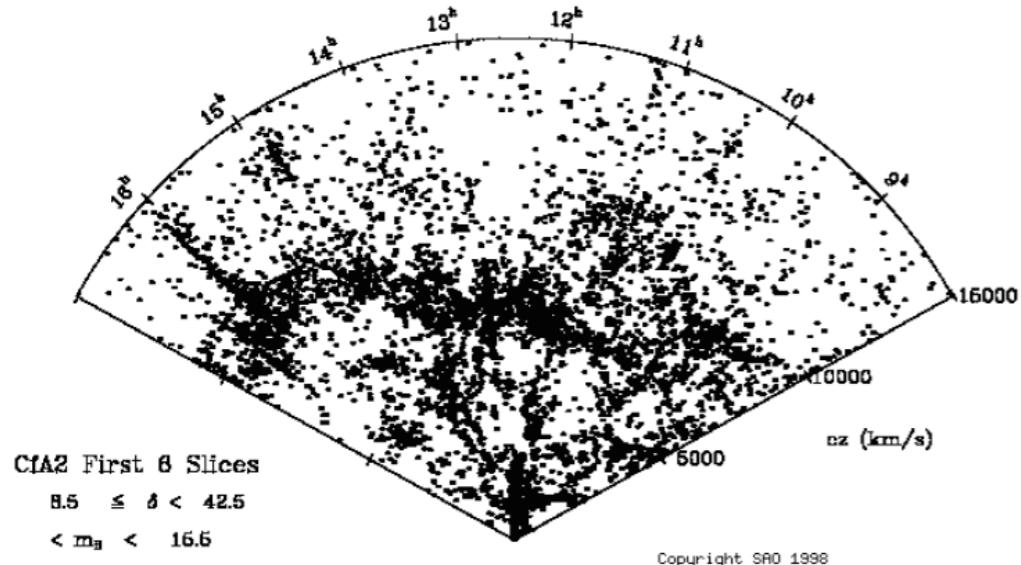
Planck Coll. A&A'20

There should be a small excess
in the two-point galaxy correlation
function around 150 Mpc!

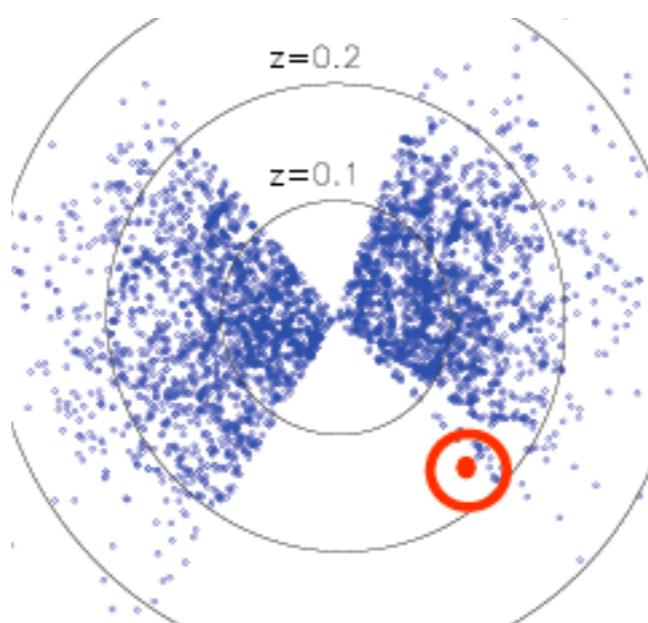


80's: Tiny surveys

Baryon Acoustic Oscillations



BAO scale

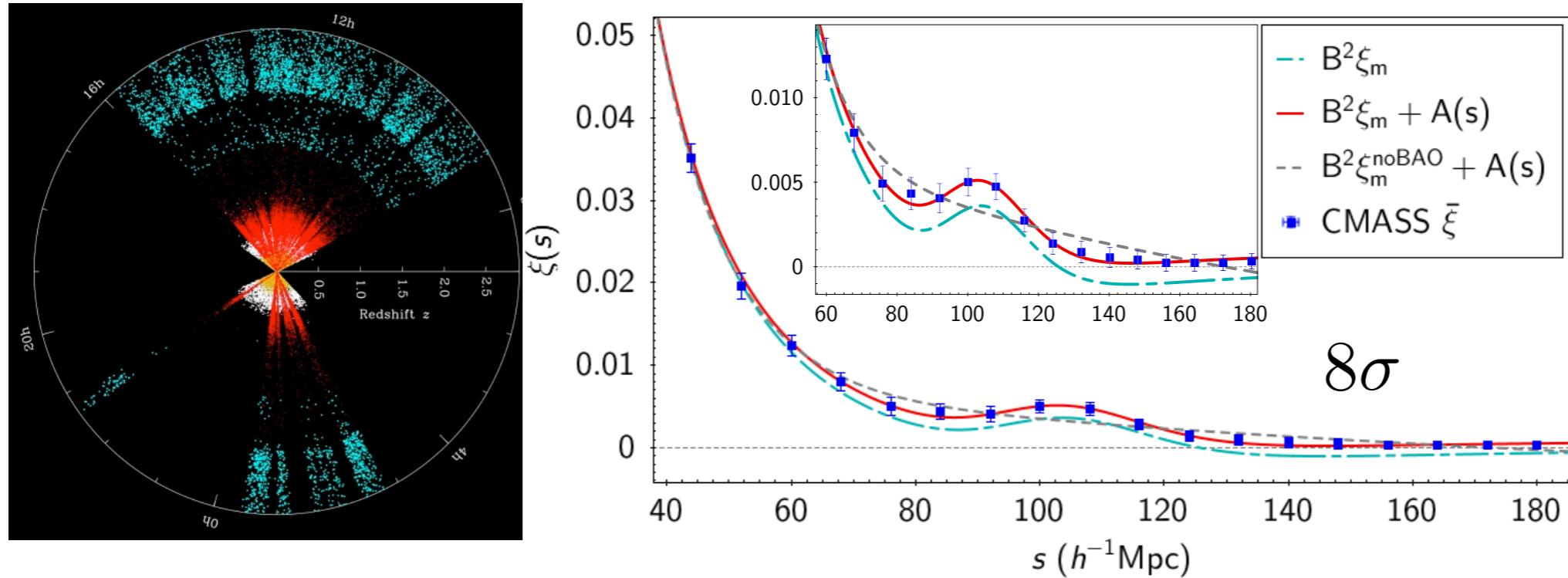


2000: Main galaxies @ SDSS.
Big number, but small volume

2005: Luminous Red Galaxies @ SDSS.
Big Volume: first detection of the BAO
signature

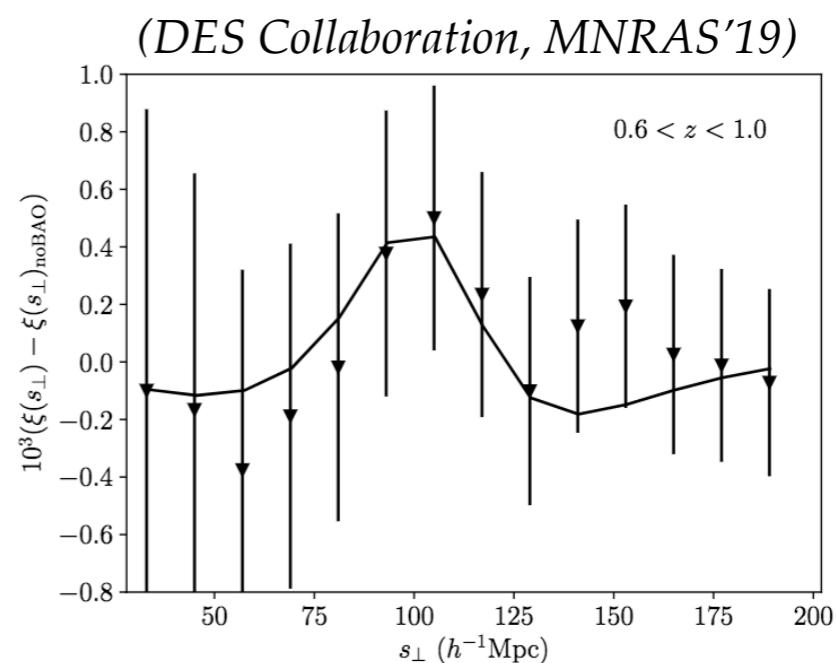
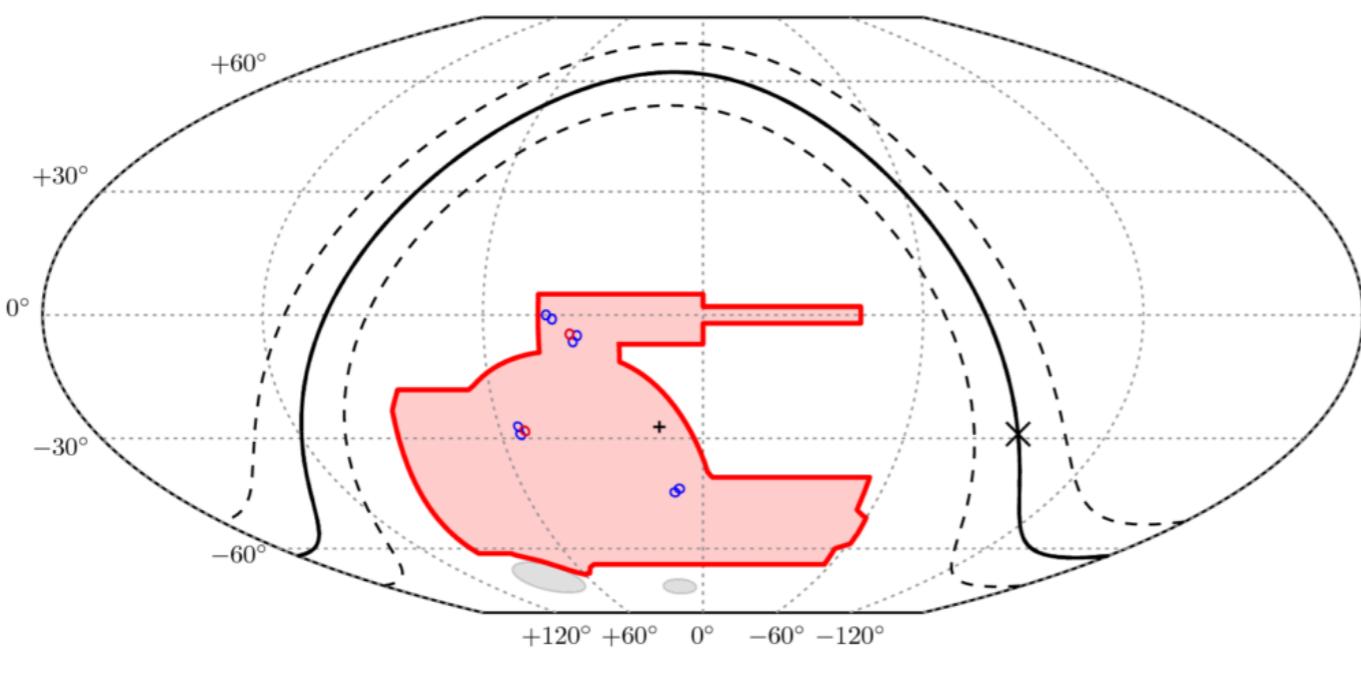
Baryon Acoustic Oscillations 2009-2014= SDSS III

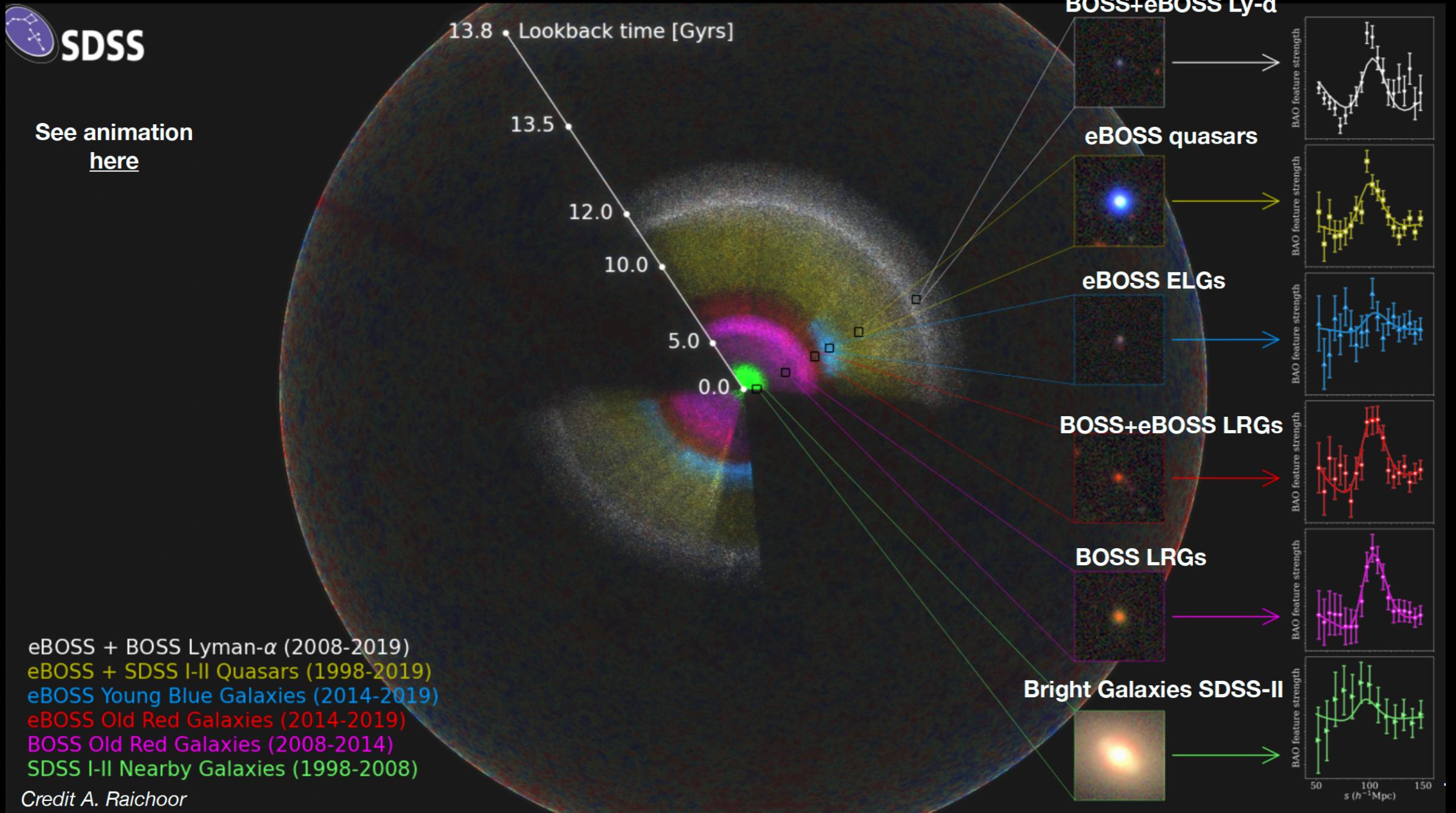
BOSS (SDSSIII) Collaboration



2013 - 2019= DES

$> 2\sigma$



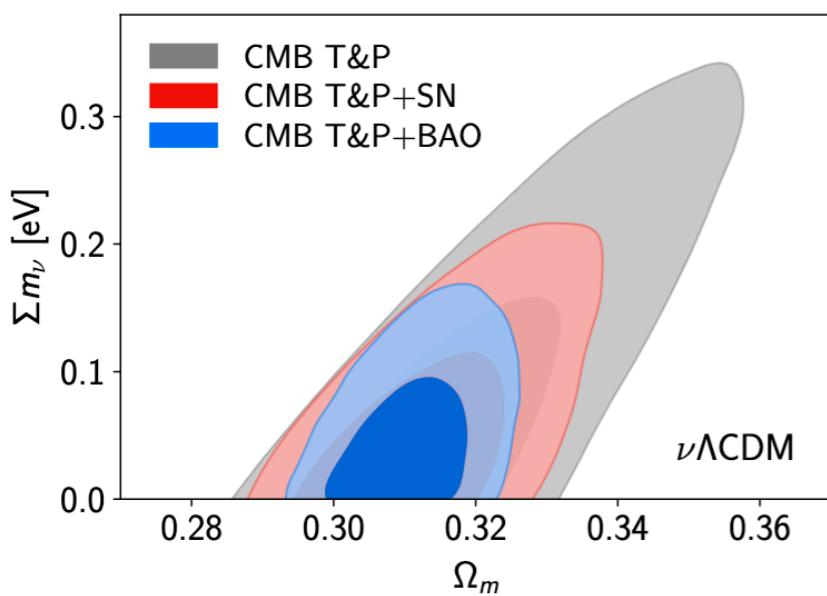


Baryon Acoustic Oscillations

2020= SDSS IV

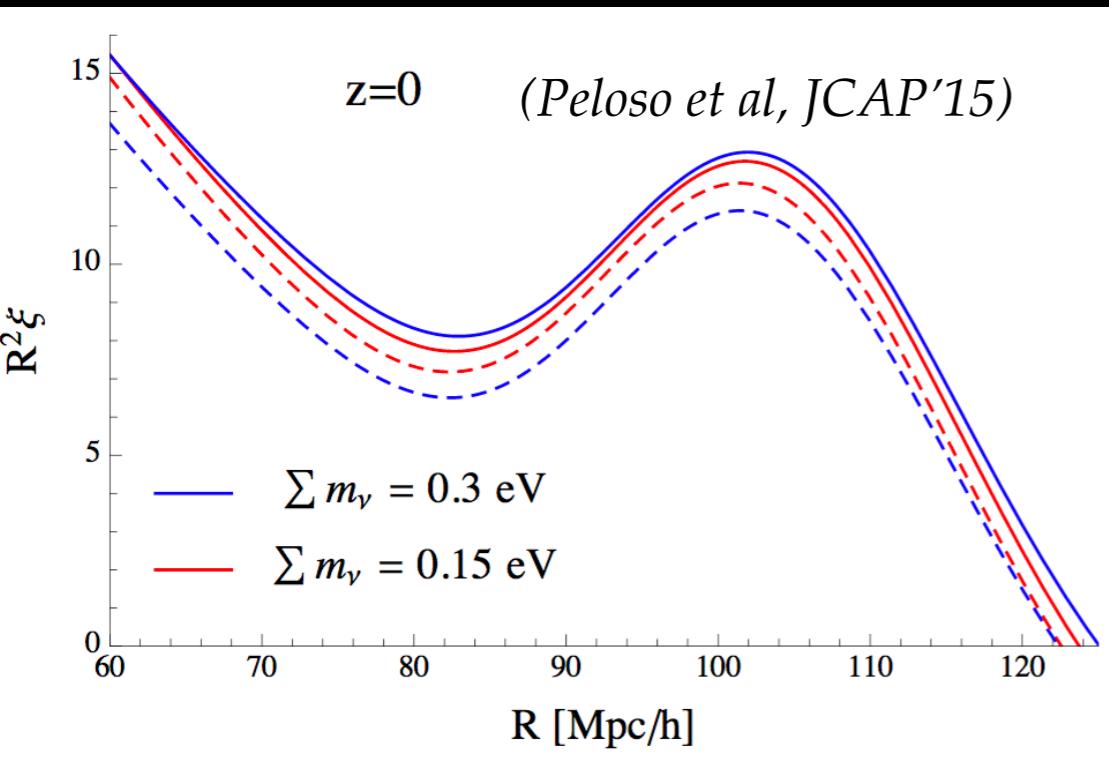
Parameter	MGS	BOSS Galaxy	BOSS Galaxy	eBOSS LRG	eBOSS ELG	eBOSS Quasar	$\text{Ly}\alpha\text{-Ly}\alpha$	$\text{Ly}\alpha\text{-Quasar}$
Sample Properties								
redshift range	$0.07 < z < 0.2$	$0.2 < z < 0.5$	$0.4 < z < 0.6$	$0.6 < z < 1.0$	$0.6 < z < 1.1$	$0.8 < z < 2.2$	$z > 2.1$	$z > 1.77$
N_{tracers}	63,163	604,001	686,370	377,458	173,736	343,708	210,005	341,468
z_{eff}	0.15	0.38	0.51	0.70	0.85	1.48	2.33	2.33
$V_{\text{eff}} (\text{Gpc}^3)$	0.24	3.7	4.2	2.7	0.6	0.6		
BAO-Only Measurements (Section 4)								
$D_V(z)/r_d$	4.47 ± 0.17				$18.33^{+0.57}_{-0.62}$			
$D_M(z)/r_d$		10.23 ± 0.17	13.36 ± 0.21	17.86 ± 0.33		30.69 ± 0.80	37.6 ± 1.9	37.3 ± 1.7
$D_H(z)/r_d$		25.00 ± 0.76	22.33 ± 0.58	19.33 ± 0.53		13.26 ± 0.55	8.93 ± 0.28	9.08 ± 0.34
RSD-Only Measurements (Section 5)								
$f\sigma_8(z)$	0.53 ± 0.16	0.500 ± 0.047	0.455 ± 0.039	0.448 ± 0.043	0.315 ± 0.095	0.462 ± 0.045		
BAO+RSD Measurements (Sections 6 and 7)								
$D_V(z)/r_d$	4.51 ± 0.14							
$D_M(z)/r_d$		10.27 ± 0.15	13.38 ± 0.18	17.65 ± 0.30	19.5 ± 1.0	30.21 ± 0.79	37.6 ± 1.9	37.3 ± 1.7
$D_H(z)/r_d$		24.89 ± 0.58	22.43 ± 0.48	19.78 ± 0.46	19.6 ± 2.1	13.23 ± 0.47	8.93 ± 0.28	9.08 ± 0.34
$f\sigma_8(z)$	0.53 ± 0.16	0.497 ± 0.045	0.459 ± 0.038	0.473 ± 0.041	0.315 ± 0.095	0.462 ± 0.045		

Alam et al, SDSS IV Coll. PRD'21



Data	95% upper limit [eV]
<i>Planck</i>	0.252
<i>Planck</i> + BAO	0.129
<i>Planck</i> + BAO + RSD	0.102
<i>Planck</i> + SN	0.170
<i>Planck</i> + BAO + RSD + SN	0.099
<i>Planck</i> + BAO + RSD + SN + DES	0.111
<i>Planck</i> + BAO + RSD + SN (νw CDM)	0.139
<i>Planck</i> + BAO + RSD + SN + DES (νw CDM)	0.161

Large scale structure: m_ν

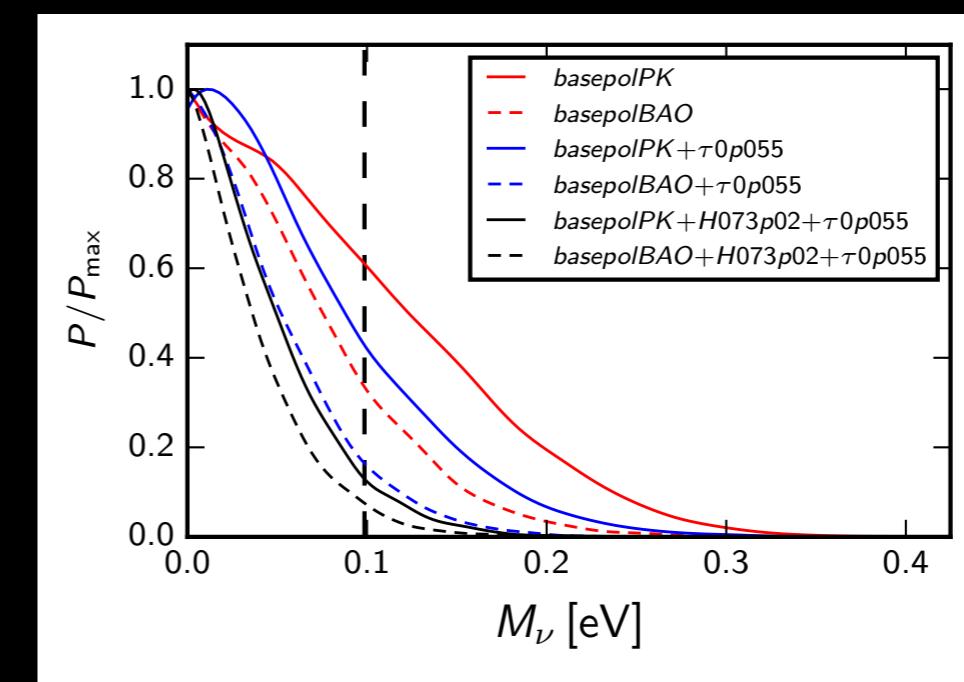
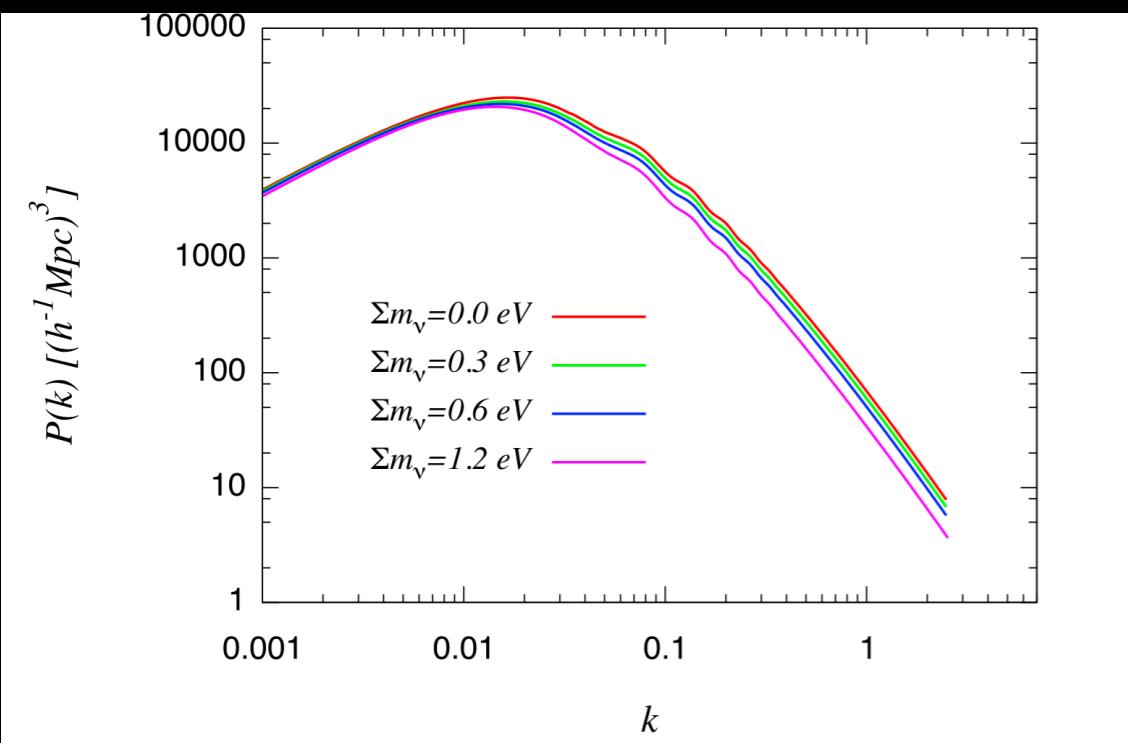


Large scale structure measurements can be interpreted either in the geometrical or shape forms

2 point correlation function

Fourier Transform
Matter power spectrum

BAO information more powerful



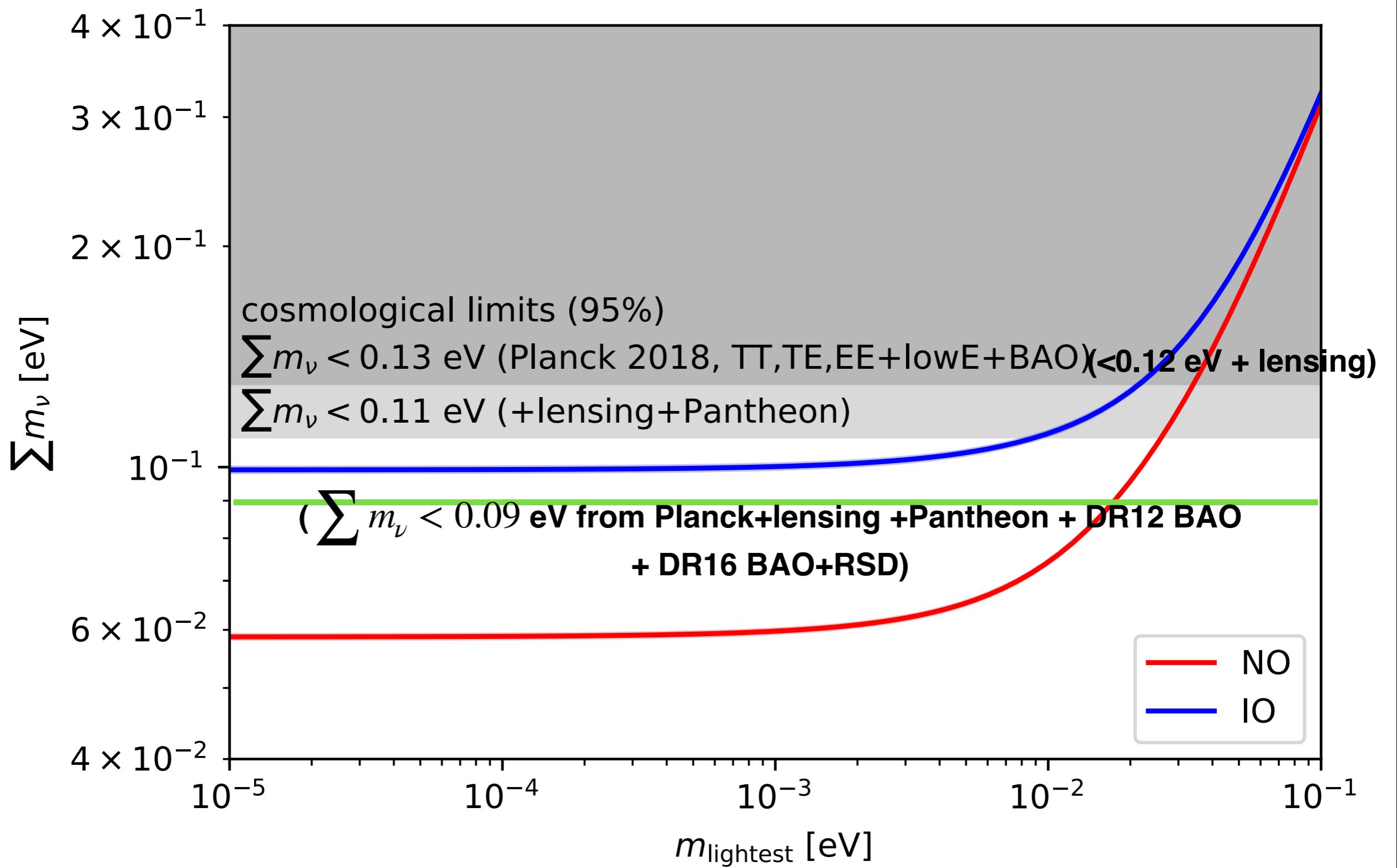
$$\sum m_\nu$$

- Planck 2018 CMB temperature polarization and lensing potential data:

$$\sum m_\nu < 0.24 \text{ eV } 95\% \text{CL}$$

- If we add large scale structure information in its BAO form

$$\sum m_\nu < 0.12 \text{ eV } 95\% \text{CL}$$



Planck TTTEEE+lowT+lowE+lensing

$$\sum m_\nu < 0.24 \text{ eV } 95\% \text{ CL}$$

+ BAO

$$\sum m_\nu < 0.12 \text{ eV } 95\% \text{ CL}$$

+ BAO + SNIa

$$\sum m_\nu < 0.11 \text{ eV } 95\% \text{ CL}$$

+ SDSS-IV (BAO + RSD) + SNIa

$$\sum m_\nu < 0.099 \text{ eV } 95 \% \text{ CL}$$

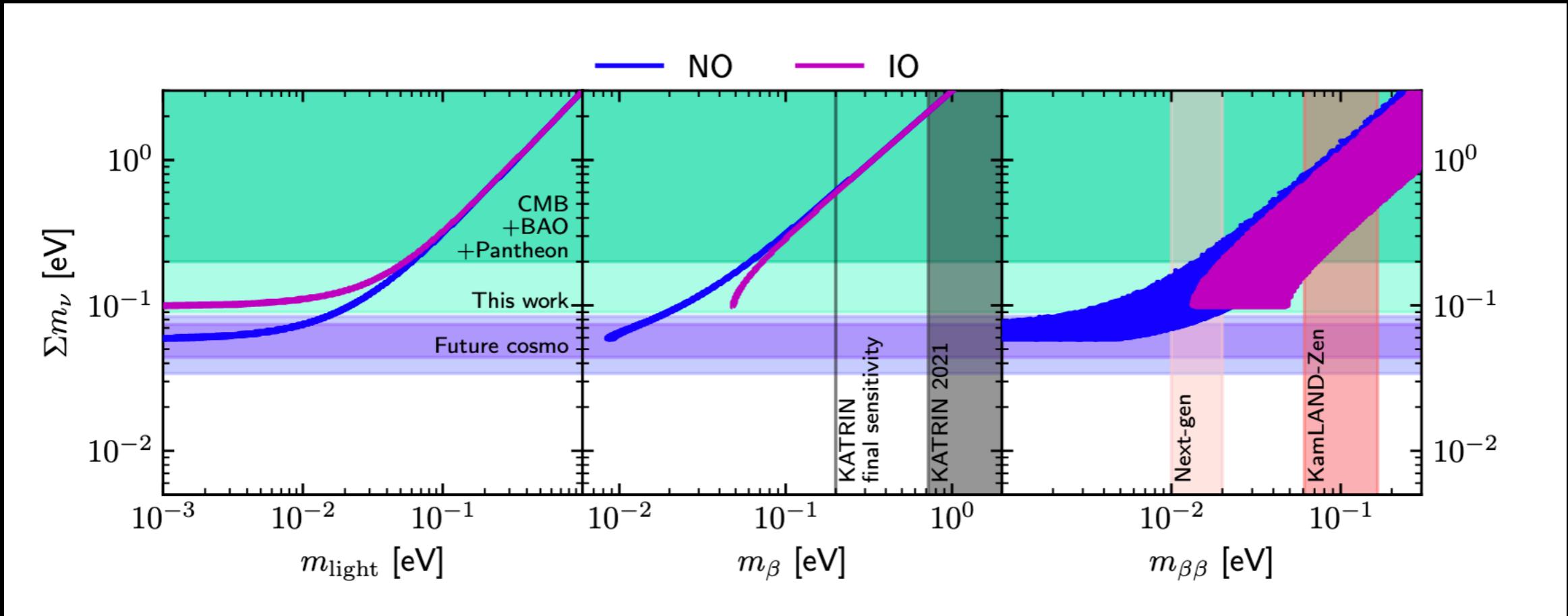
eBOSS Coll. PRD'21

+ SDSS-IV (BAO + RSD) DR16 + DR12 BAO + SNIa

$$\sum m_\nu < 0.09 \text{ eV } 95\% \text{ CL}$$

Di Valentino et al, PRD'21

Planck+lensing +Pantheon	Σm_ν [eV]	H_0 [km/s/Mpc]	Ω_m	σ_8	S_8	$\ln B_{0-NH}$	$\ln B_{NH-IH}$
+ DR12 <i>BAO only</i>	< 0.116	67.8 ± 1.0	$0.309^{+0.013}_{-0.012}$	$0.814^{+0.017}_{-0.019}$	0.826 ± 0.022	-1.3	-1.5
+ DR12 <i>BAO+RSD</i>	< 0.118	67.8 ± 1.0	$0.310^{+0.013}_{-0.012}$	$0.814^{+0.017}_{-0.019}$	$0.827^{+0.021}_{-0.022}$	-1.3	-1.7
+ DR16 <i>BAO only</i>	< 0.158	$67.5^{+1.2}_{-1.3}$	$0.314^{+0.017}_{-0.016}$	$0.811^{+0.020}_{-0.023}$	$0.830^{+0.023}_{-0.024}$	-0.7	-1.6
+DR16 <i>BAO+RSD</i>	< 0.101	$67.9^{+1.0}_{-1.1}$	$0.308^{+0.014}_{-0.013}$	$0.817^{+0.016}_{-0.017}$	0.828 ± 0.022	-1.7	-1.9
+DR12 <i>BAO only</i> + DR16 <i>BAO only</i>	< 0.121	$67.78^{+0.90}_{-0.97}$	$0.310^{+0.013}_{-0.011}$	$0.813^{+0.017}_{-0.019}$	0.826 ± 0.021	-0.9	-1.8
+DR12 <i>BAO only</i> + DR16 <i>BAO+RSD</i>	< 0.0866	$68.09^{+0.85}_{-0.88}$	0.306 ± 0.011	$0.817^{+0.015}_{-0.016}$	$0.826^{+0.020}_{-0.021}$	-1.9	-2.0
+DR12 <i>BAO+RSD</i> + DR16 <i>BAO only</i>	< 0.125	$67.71^{+0.89}_{-0.97}$	$0.311^{+0.012}_{-0.011}$	$0.813^{+0.017}_{-0.019}$	0.828 ± 0.021	-1.1	-1.4
+DR12 <i>BAO+RSD</i> + DR16 <i>BAO+RSD</i>	< 0.0934	$68.00^{+0.87}_{-0.89}$	0.307 ± 0.011	$0.817^{+0.015}_{-0.016}$	0.827 ± 0.021	-1.9	-1.8



Cosmic neuTRIVIA game steps

1. Familiarize yourself with the board's layout:

- The Λ CDM trivia: the players!
- The neutrino pie piece: decoupling in the early universe

2. Roll the dice and get:

- Number of neutrinos and Big-Bang Nucleosynthesis
- Number of neutrinos and Cosmic Microwave Background Radiation
- Neutrino masses and Cosmic Microwave Background Radiation?
- Neutrino masses and structure formation in the universe?

3. Is anyone cheating? Neutrinos and Tensions

4. Final score:

- Take home messages

PSST!

Do you know the answer to 1?

How fast the universe is expanding?

I Know the value of H₀ from Supernovae Ia....

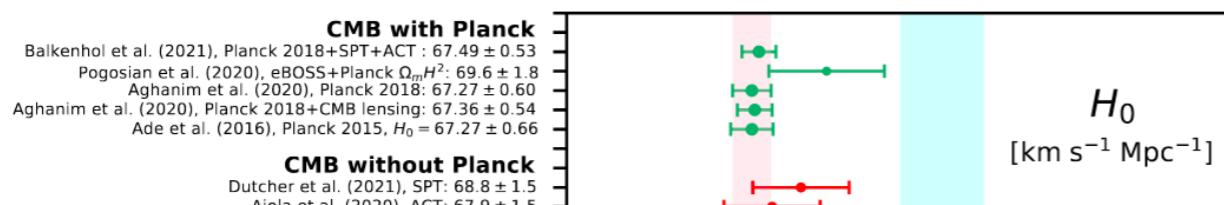
Do you know the one from the CMB?



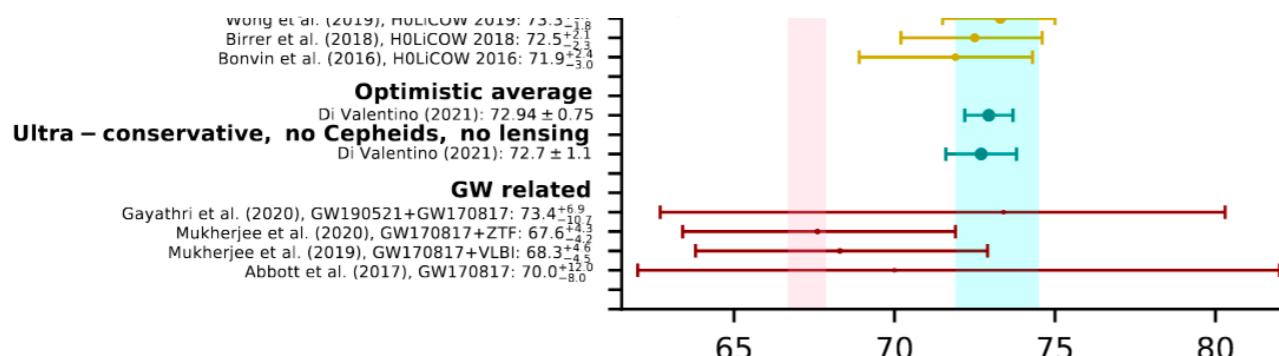
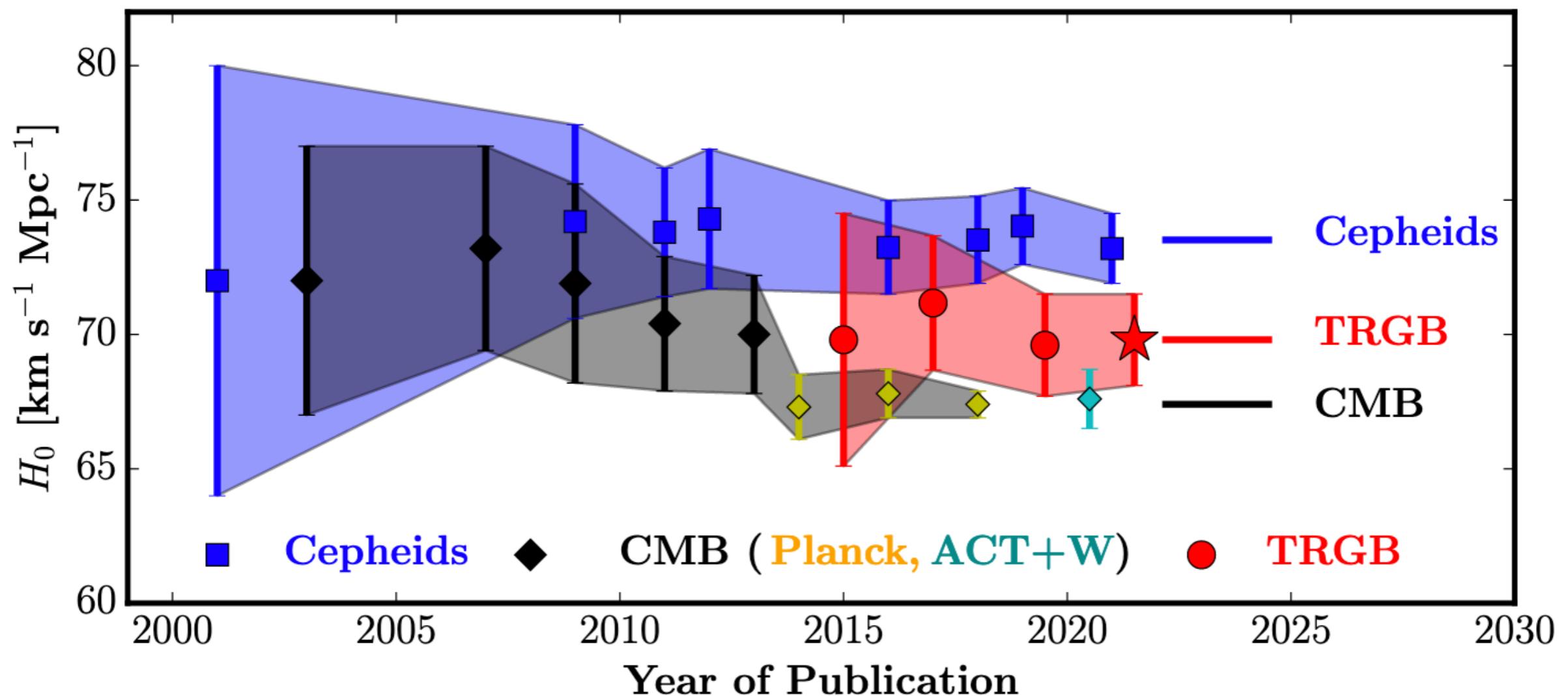
The Hubble tension

W. Freedman, APJ'21

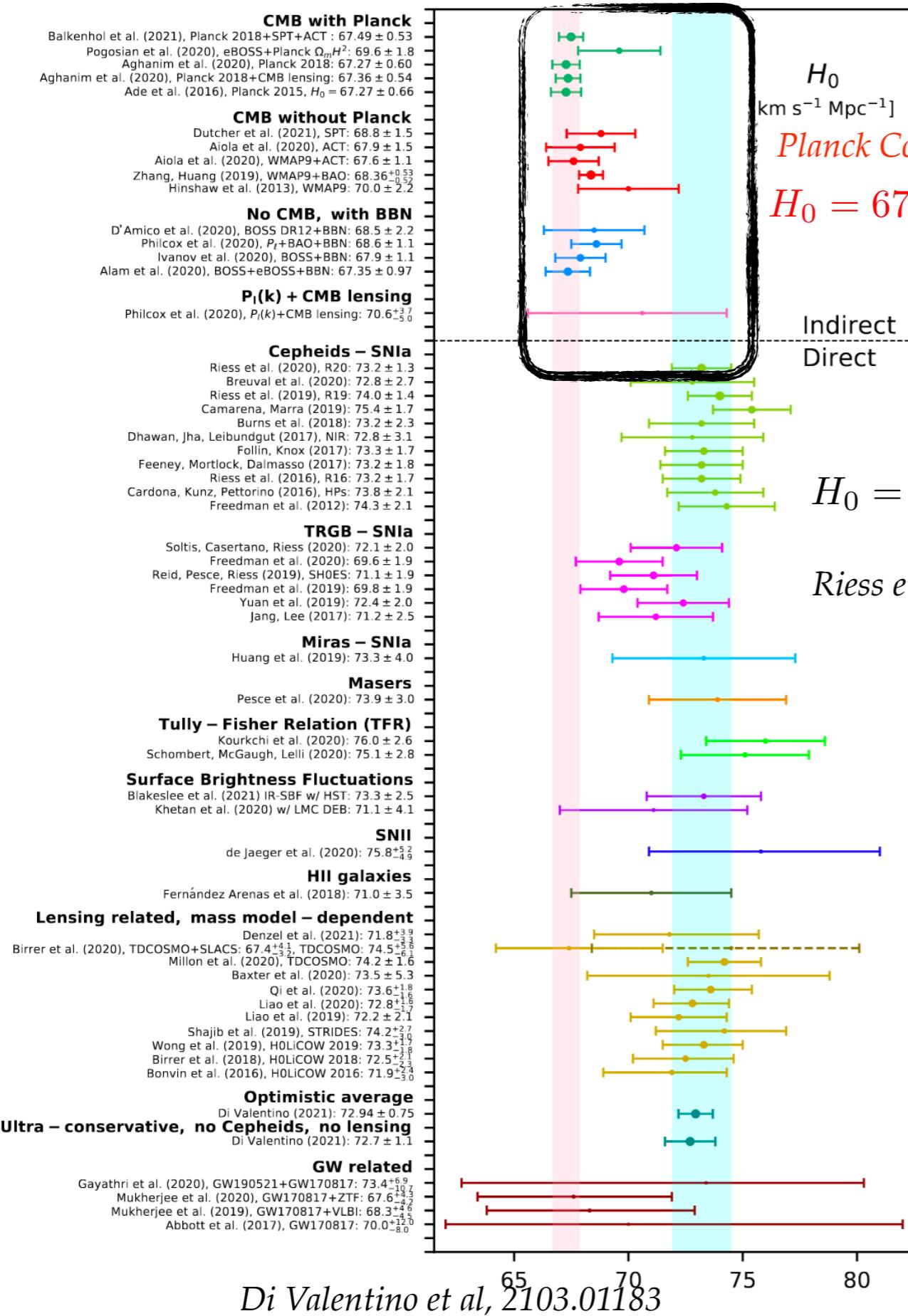
Di Valentino et al, 2103.01183



Hubble Constant Over Time



The Hubble tension



H_0
 $\text{km s}^{-1} \text{Mpc}^{-1}]$
Planck Coll. A&A'20

$$H_0 = 67.27 \pm 0.60 \text{ km/s/Mpc}$$

Indirect
Direct

$$H_0 = 73.2 \pm 1.3 \text{ km/s/Mpc}$$

Riess et al, APJ'20

Early, indirect
(they are a prediction within the Λ CDM framework) inferences tend to increase (rather than decrease) from the baseline value derived from Planck 2018 temperature anisotropy data with the inclusion of CMB polarization, BAO data, etc...

The Hubble tension

Late (I)

Best-established and unique empirical methods to measure H_0

The “distance ladder”: parallax (=geometry) is used to calibrate luminosities of stars that after can be seen at much larger distances.

Cepheids/SNIa: SHOES team

$$H_0 = 73.2 \pm 1.3 \text{ km/s/Mpc} \quad \text{Riess et al, APJ'20}$$

Reanalyses of SHOES data have fully agreed!

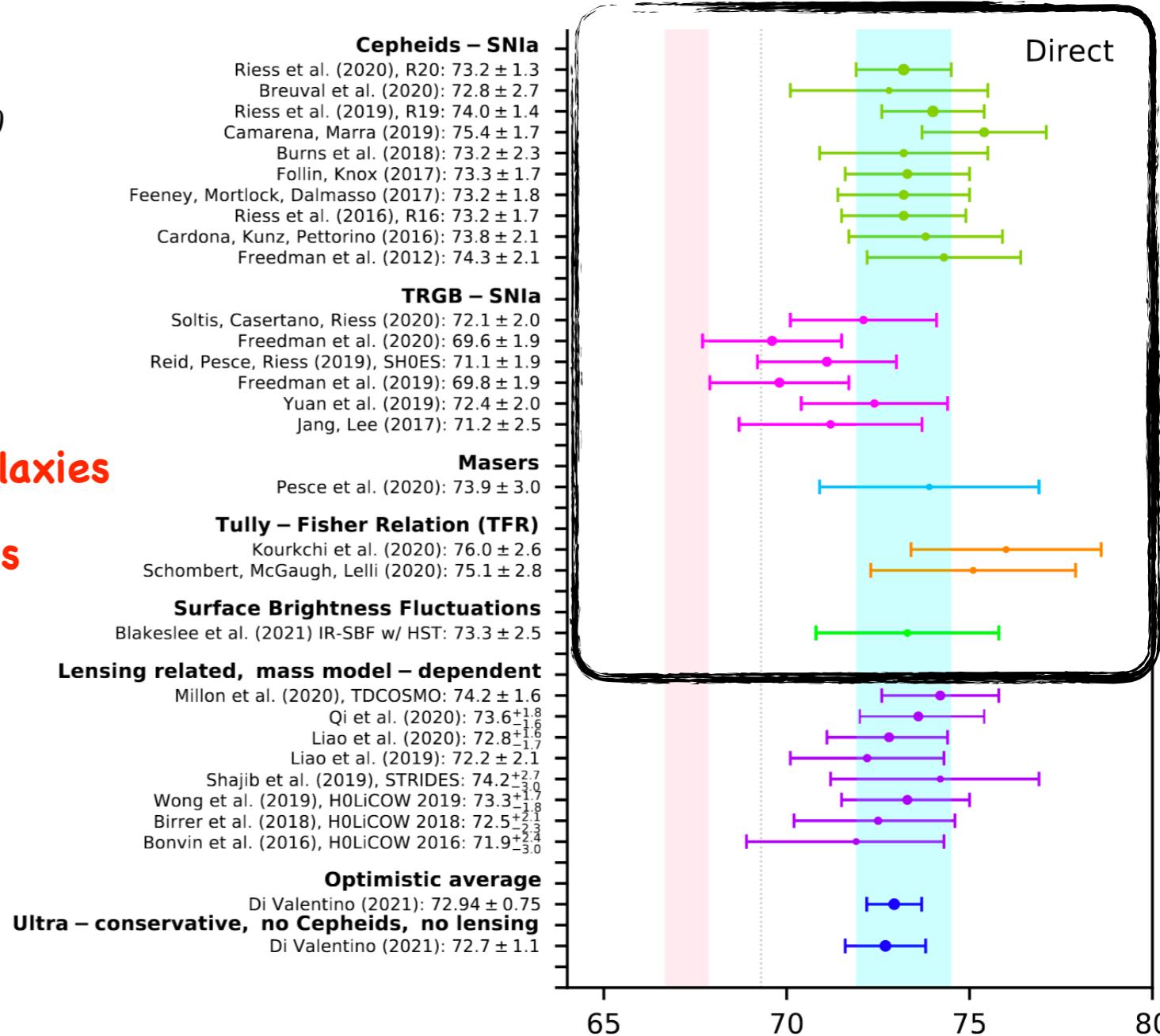
Other distance ladders:

Tip of the RedGiant Branch (TRGB)

Geometric distances to Mega-Maser hosting galaxies

Empirical relation luminosity-rotational velocities

SBF method as a distance indicator



The Hubble tension

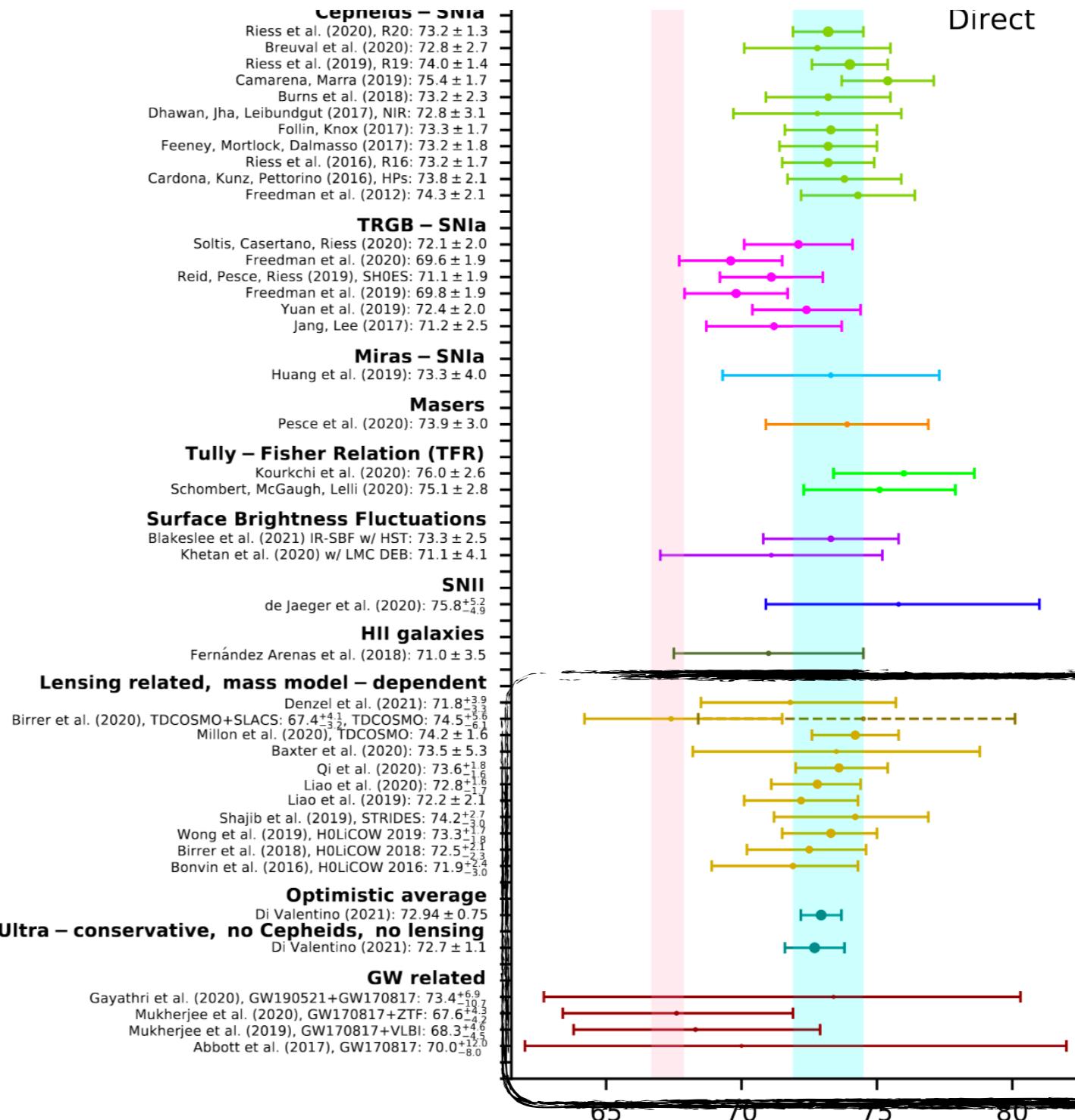
Late (II): “Gastrophysics” model-dependent



Time delays in strong lensed systems
(quasars)(lens mass distribution/foregrounds)

Late (III): Standard sirens
combination of the distance to the source
inferred purely from gravity waves,
with the recession velocity from
redshift data using electromagnetic signal

Di Valentino et al, 2103.01183



Systematics??

From V. Poulin

Several sources are required!

SNIa: Calibration issues

*Follin & Knox, MNRAS'18;
Feeney et al, MNRAS'18;
Freedman et al'19&20;
Yuan et al, APJ'19;
Soltis, Casertano & Riess, APjL'21;
Efstathiou'20*

Different populations between “local” and Hubble flow SNIa?

*Rigault et al, APJ'15
Jones et al, APJ18
Brout & Scolnic, APJ'21*

Do we live in a void?

*Wu & Huterer, MNRAS'17
Kenworthy, Scolnic & Riess, APJ'19*

Strong-lensed quasars: Lens profiles

*Blum, Castorina & Simonovic, APjL'20
Birrer et al, A&A'20*

Is the cosmological principle wrong?

*Colin et at, A&A'19
Heinesen & Bouchert, CQG'20
Secrest wt al, APjL'21*

CMB and the Hubble parameter

Within the (let's say) Λ CDM framework, the recipe is the following:

1) From measurements of the matter and baryon densities given a model:
derivation of r_s^* at the last scattering redshift z_s

2) From the position of the CMB peaks, the comoving angular diameter distance is extracted:

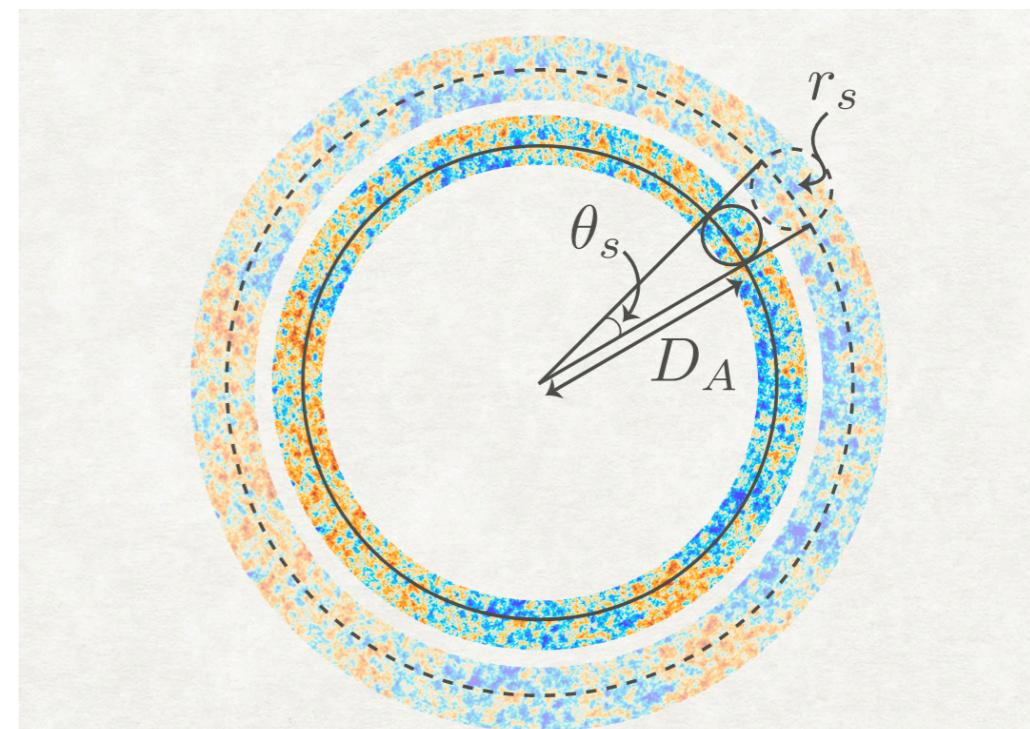
$$D_A^* \equiv r_s^* / \theta_s^*$$

3) Once we have the angular diameter distance, we can infer the value of H_0 .

$$D_A^* \propto 1/H_0$$

Could the last scattering surface be closer?

Could CMB spots be smaller?



The Hubble parameter compendium

Knox & Millea, PRD'20

$$\theta_s \equiv r_s / D_A$$

Early universe solutions

Change both r_s & D_A at recombination to have a higher H_0

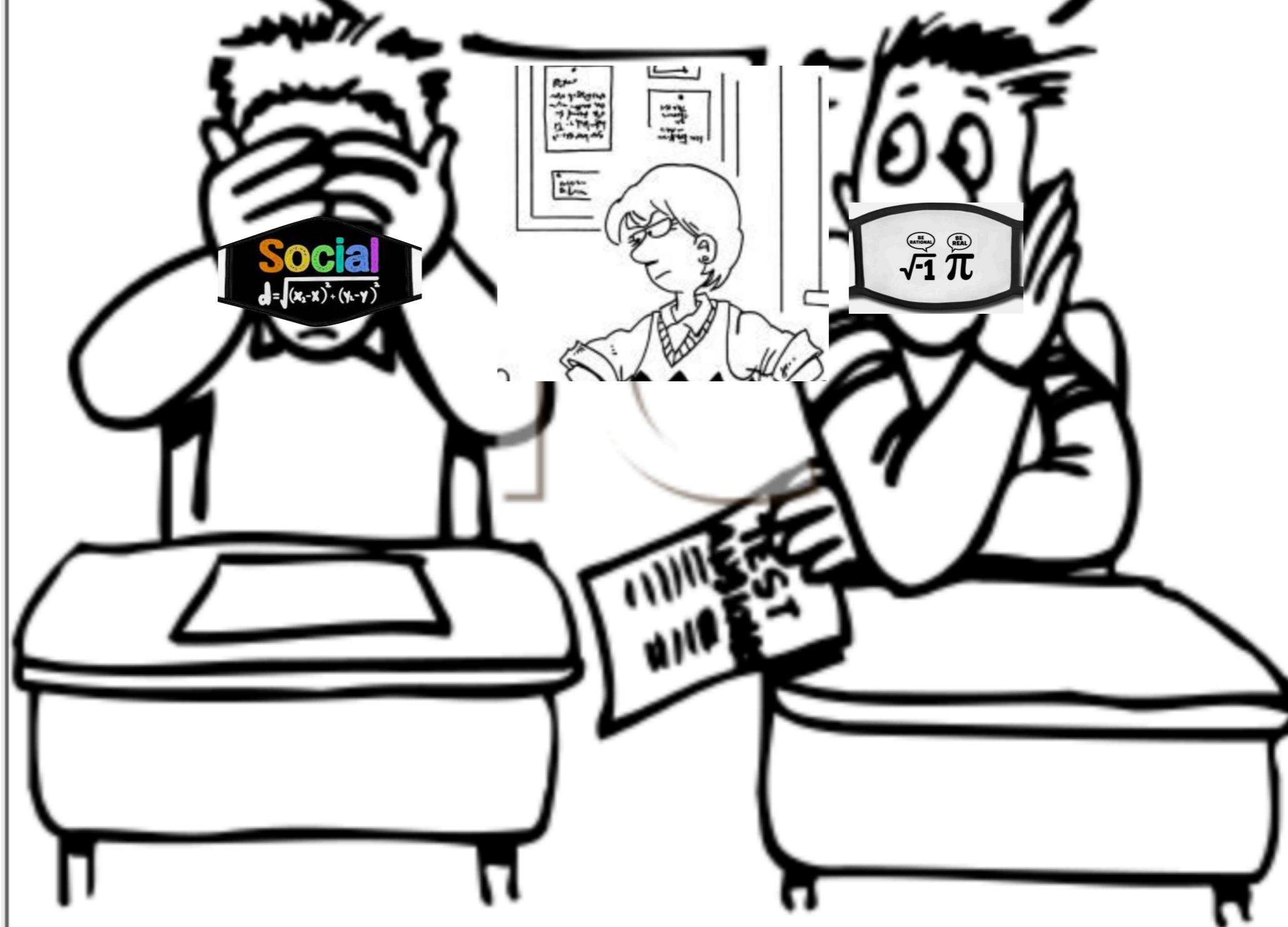
Late universe solutions

Change $D_A(z)$ at redshifts after recombination to have a lower H_0

All solutions are required to leave unchanged θ_s , θ_{EQ} & θ_D

Shhh!! masks on!!!

I'll call the neutrino major and the extra relativistic stuff.
They may know the solution for this!



H_0 versus N_{eff}

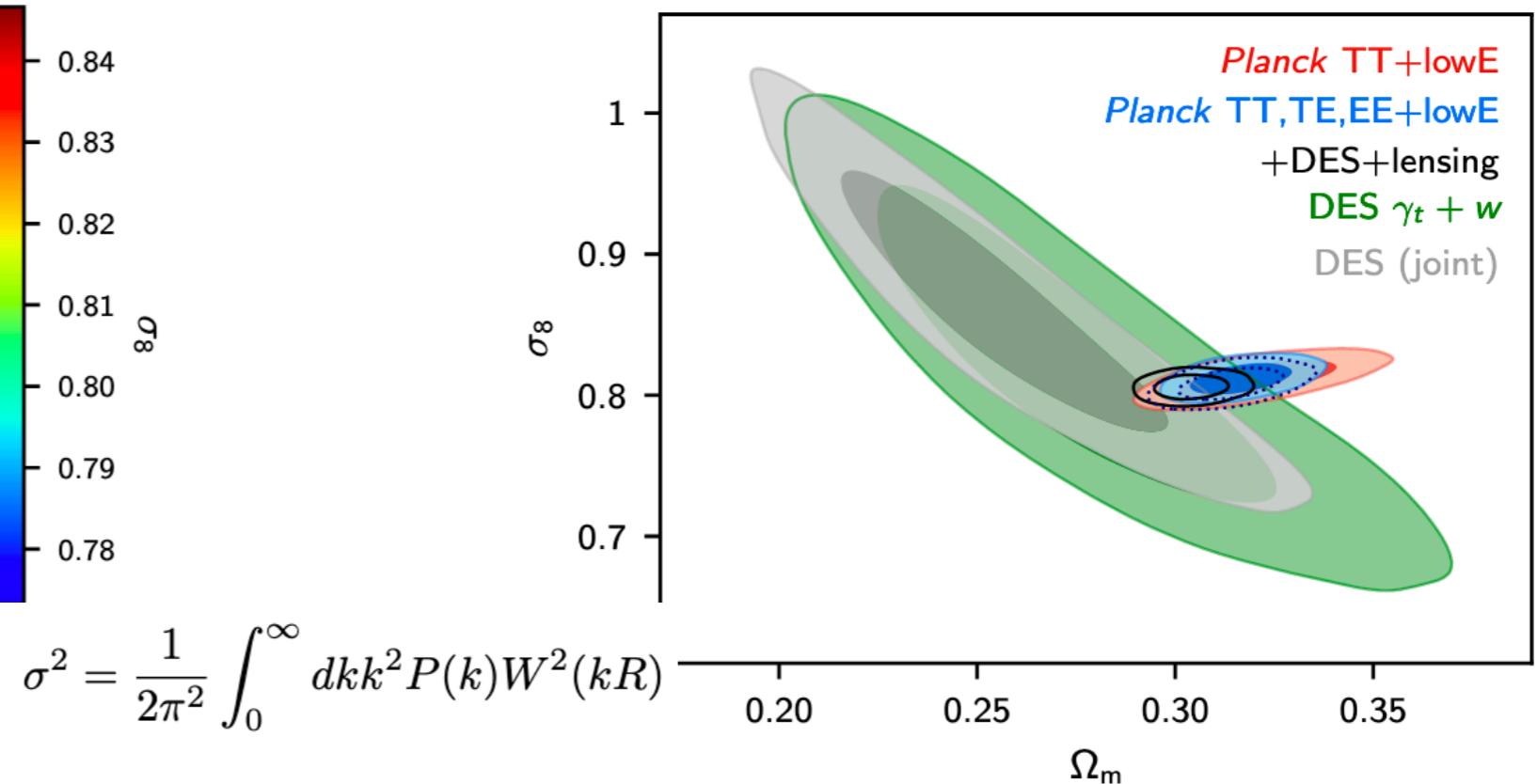
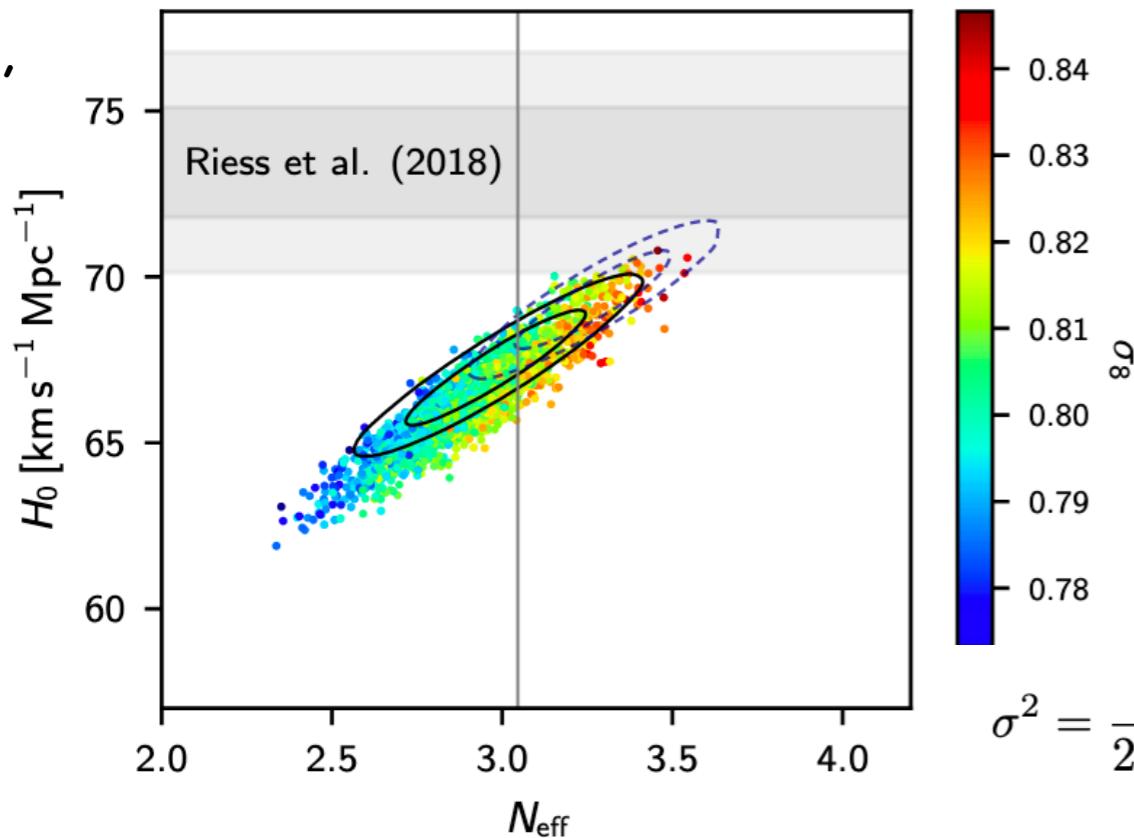
$$\Omega_r h^2 = \left(1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right) \Omega_\gamma h^2$$

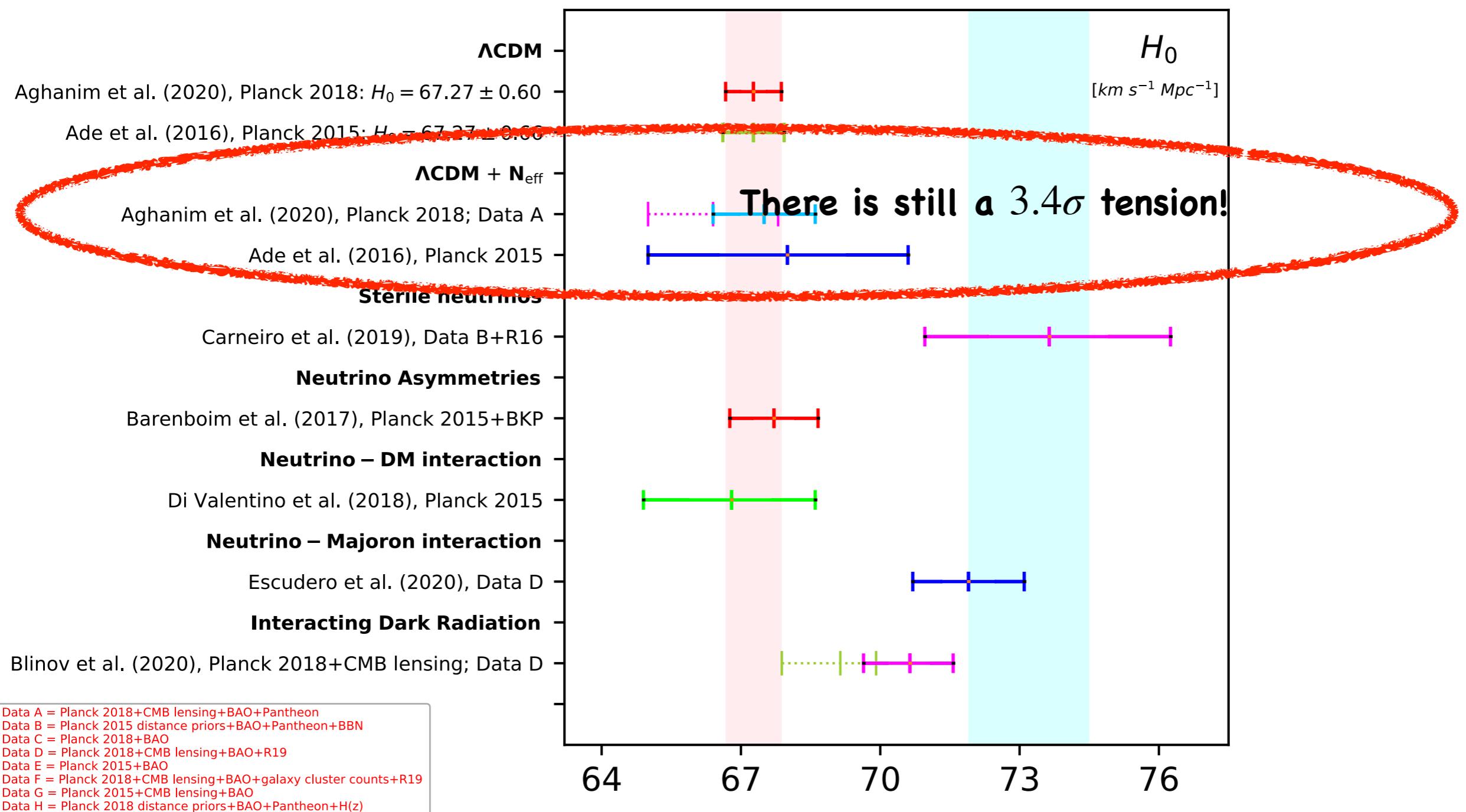
An increase of N_{eff} implies a larger expansion in the early universe: shorter r_s .
 D_A at recombination needs to be smaller: higher H_0

$$\theta_s = \frac{r_s}{D_A}$$

$$D_A = \int_0^{z_{\text{rec}}} \frac{dz}{H(z)}$$

However, this will increase the tensions with weak galaxy lensing and cluster counts data on (in order to keep $\Omega_m h^2$ constant, we need a lower Ω_m , implying a higher σ_8)

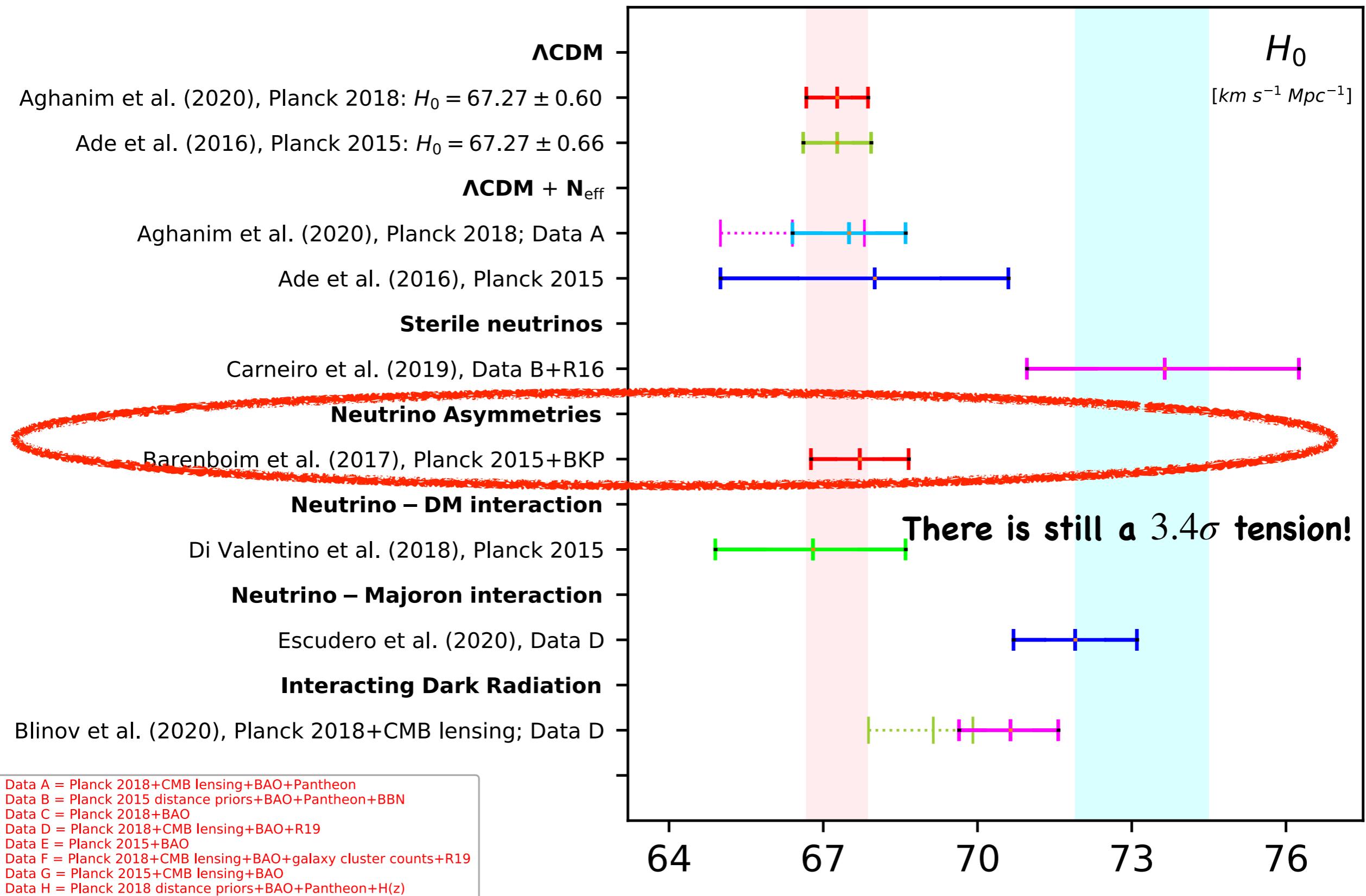


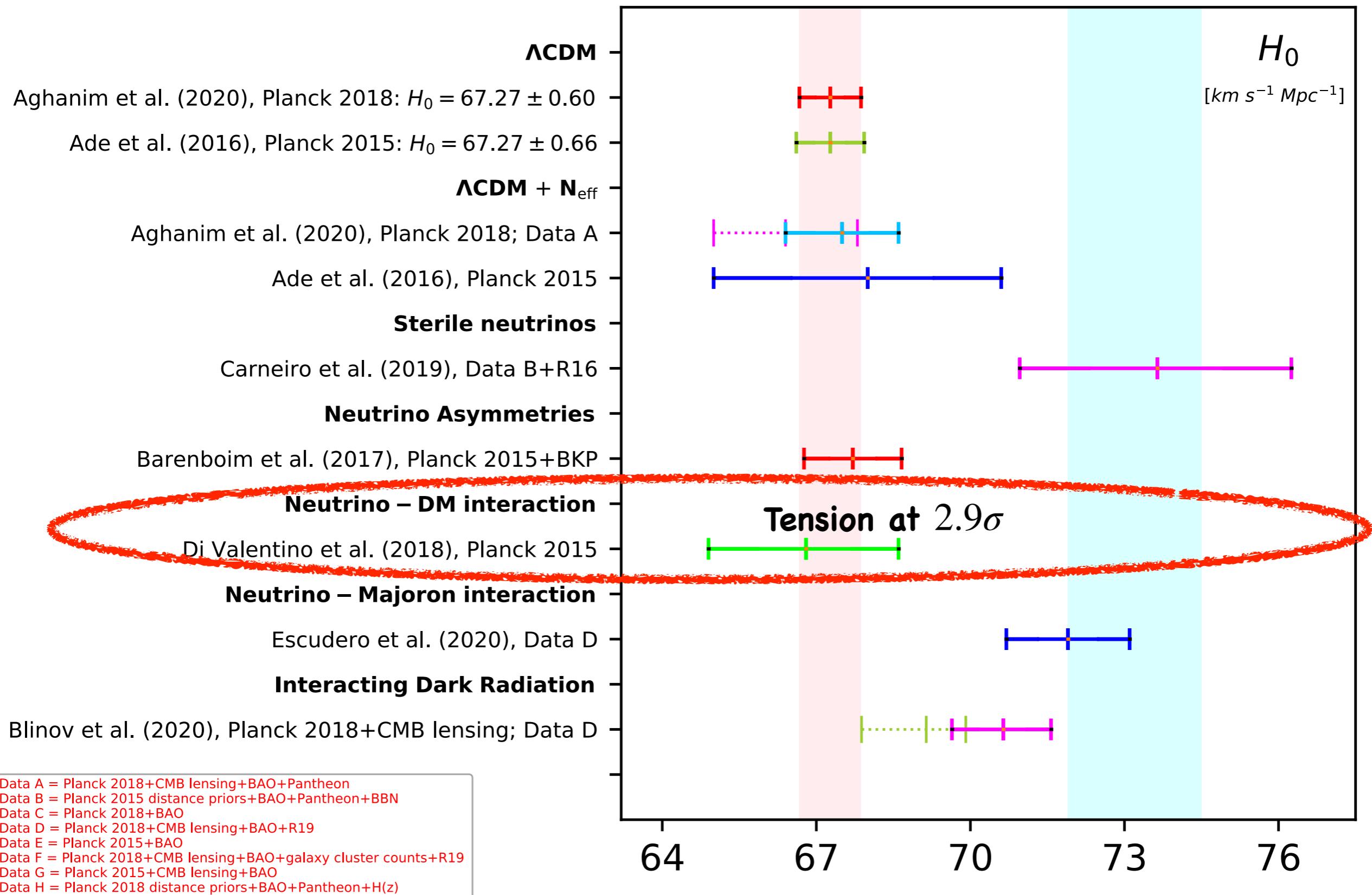


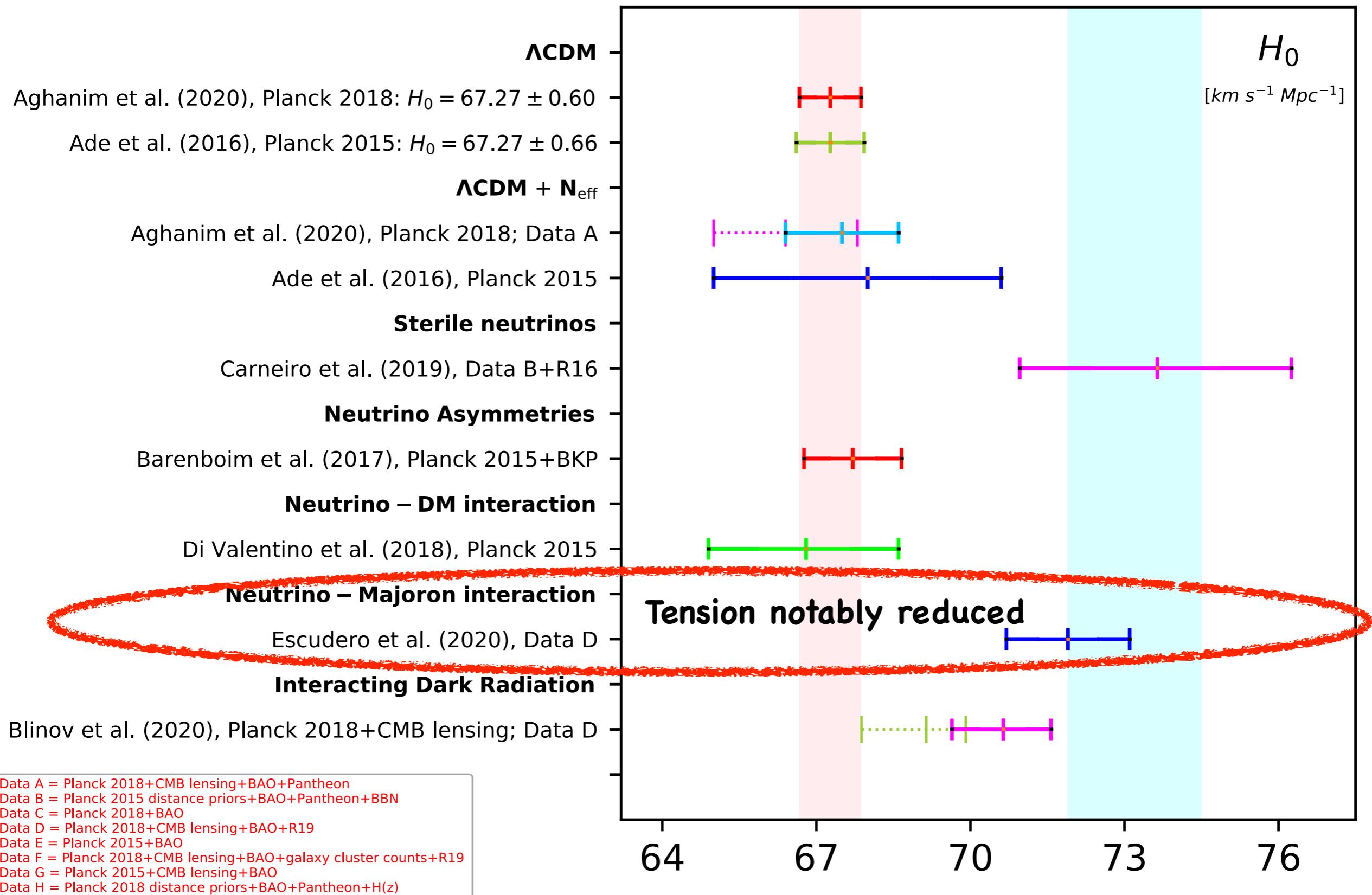
Di Valentino et al, 2103.01183

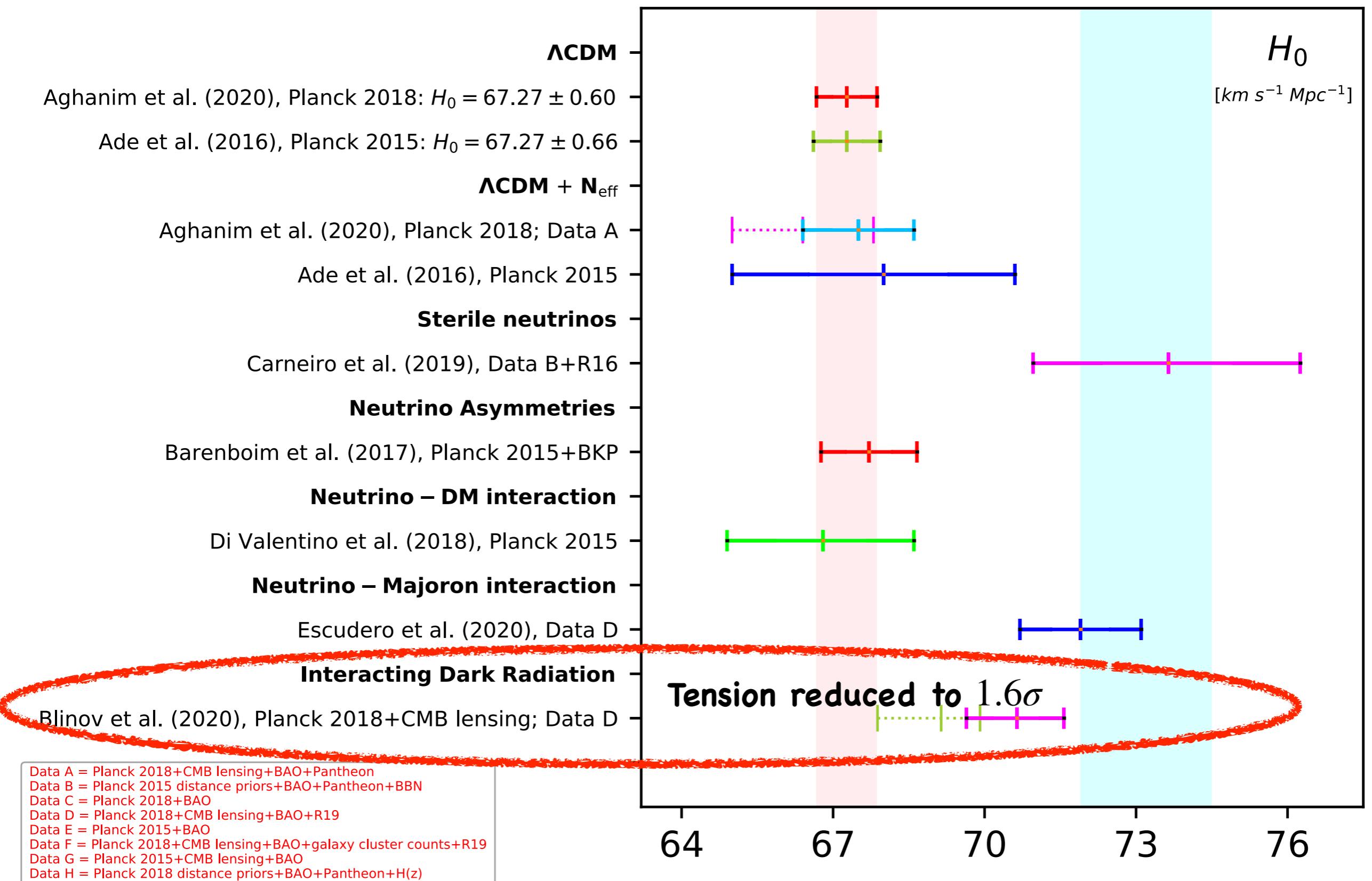
H₀ versus neutrino asymmetries

$$\Delta N_{eff} = \frac{15}{7} \left(\frac{\xi_i}{\pi} \right)^2 \left(2 + \left(\frac{\xi_i}{\pi} \right)^2 \right)$$









Interacting neutrinos

Free-streaming neutrinos travel supersonically through the photon-baryon plasma at early times, inducing a net phase shift in the CMB power spectra towards larger scales (smaller multipoles), as well as a slight suppression of its amplitude.

$$\delta\phi \simeq 0.1912\pi \frac{\rho_\nu}{\rho_r}$$

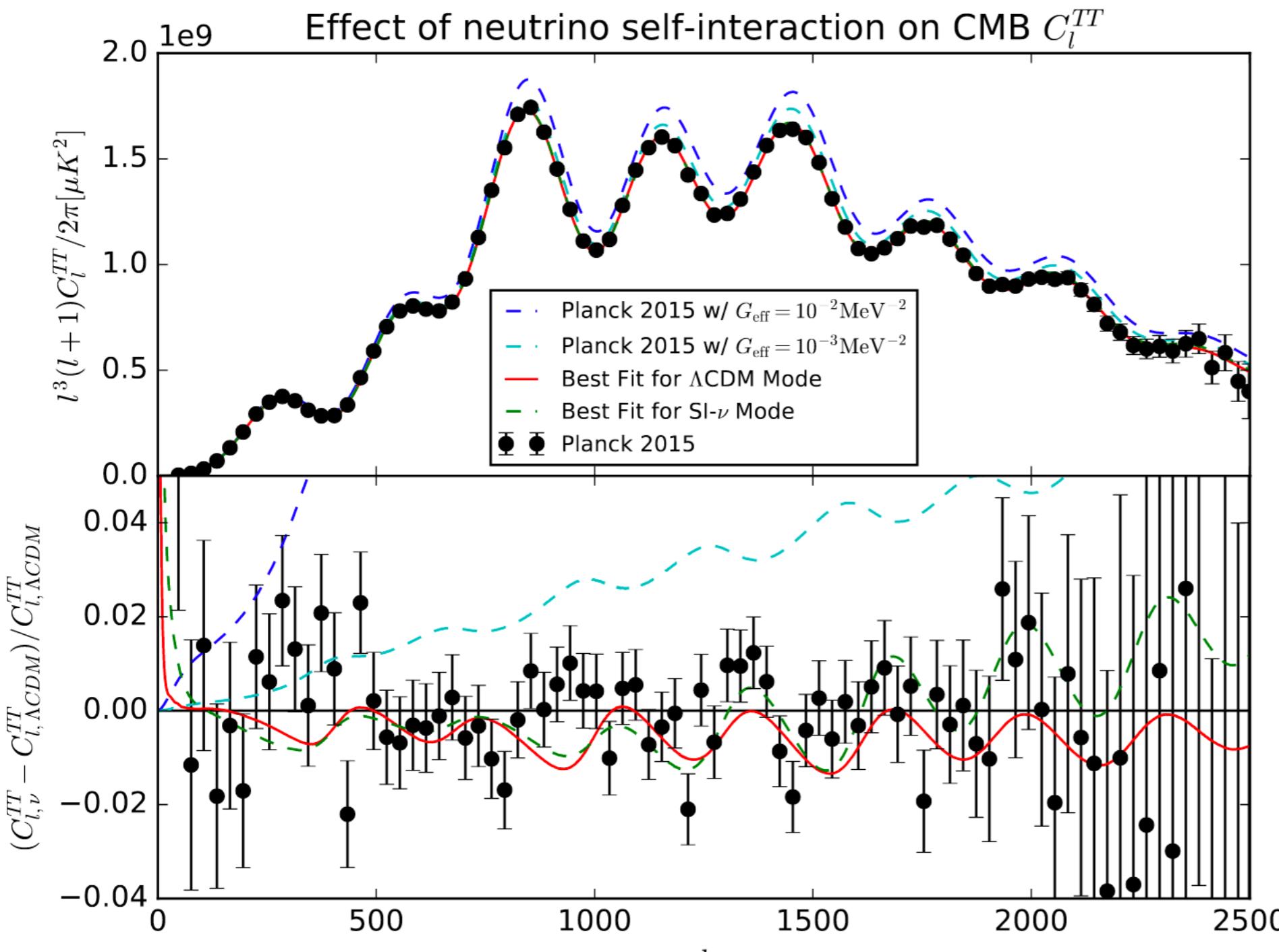
Bashinsky & Seljak, PRD'04

Bashinsky & Seljak, PRD'04; Follin et al, PRL'15;
Baumann et al, JCAP'16; Choi, Chiang & LoVerde, JCAP'18;
Baumann, Green & Zaldarriaga, JCAP'17;
Lancaster et al, JCAP'17; Oldengott, Tram & Wong, JCAP'17;
Kreisch, Cyr-Racine & Doré, PRD'20

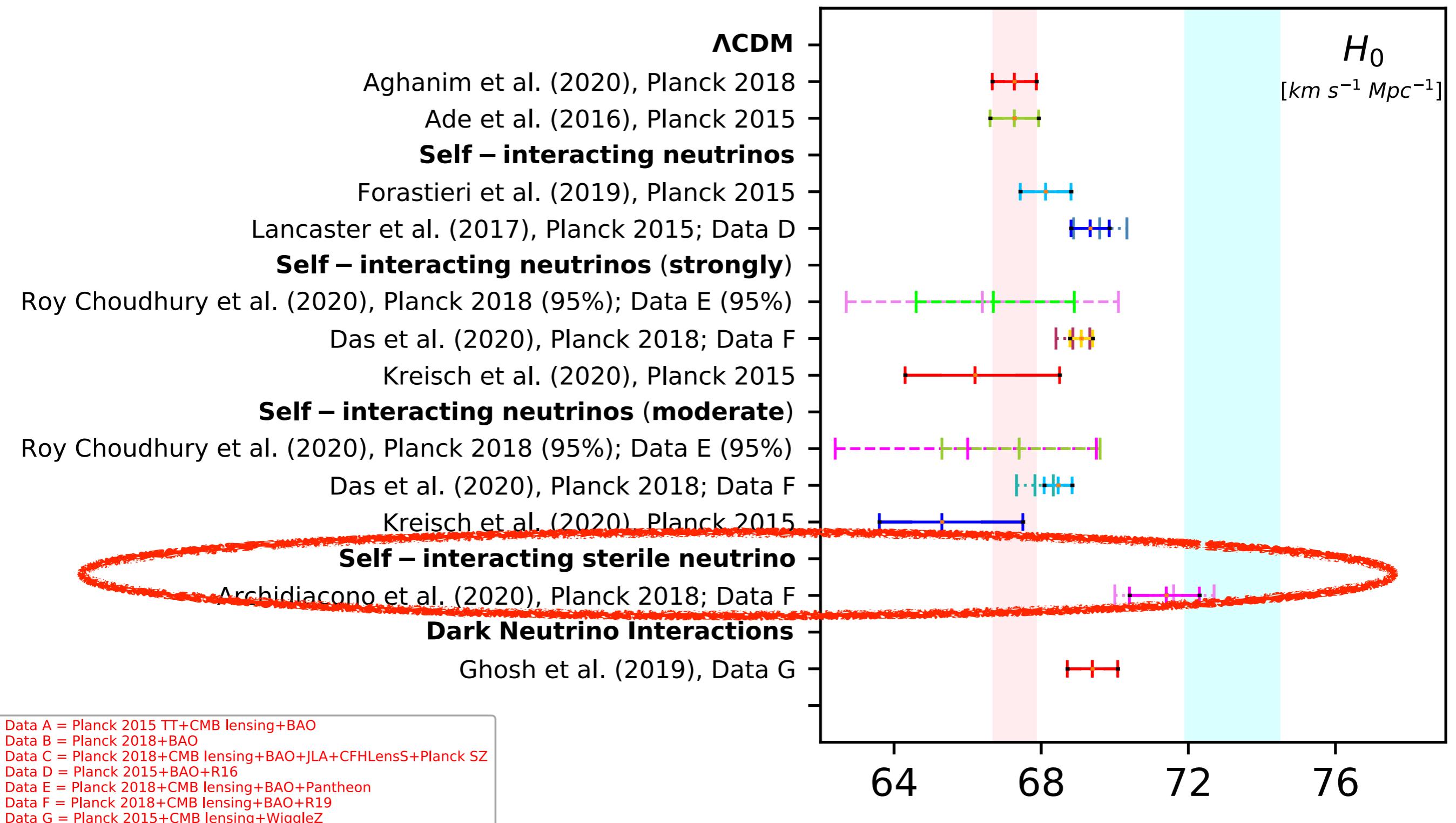
Free-streaming neutrinos lead to a physical size of the photon sound horizon at last scattering that is slightly larger.

Interacting neutrinos shift the power spectrum towards smaller scales and reduce the physical size of photon sound horizon at last scattering: a smaller value of D_A = higher value of H_0 is required!

Interacting neutrinos



Lancaster et al, JCAP'17



tension $\leq 1\sigma$ “Excellent models”	tension $\leq 2\sigma$ “Good models”	tension $\leq 3\sigma$ “Promising models”
Dark energy in extended parameter spaces [292]	Early Dark Energy [238]	Early Dark Energy [232]
Dynamical Dark Energy [312]	Phantom Dark Energy [11]	Decaying Warm DM [477]
Metastable Dark Energy [317]	Dynamical Dark Energy [11, 284, 312]	Neutrino-DM Interaction [509]
PEDE [395, 397]	GEDE [400]	Interacting dark radiation [520]
Elaborated Vacuum Metamorphosis [403–405]	Vacuum Metamorphosis [405]	Self-Interacting Neutrinos [703, 704]
IDE [317, 639, 640, 642, 655, 660, 664–666]	IDE [317, 656, 659, 664, 666, 673]	IDE [659]
Self-interacting sterile neutrinos [714]	Critically Emergent Dark Energy [997]	Unified Cosmologies [750]
Generalized Chaplygin gas model [747]	$f(\mathcal{T})$ gravity [817]	Scalar-tensor gravity [859]
Galileon gravity [879, 885]	Über-gravity [59]	Modified recombination [986]
Power Law Inflation [966]	Reconstructed PPS [978]	Super Λ CDM [1007]
$f(\mathcal{T})$ [821]		Coupled Dark Energy [653]

Table B1. Models solving the H_0 tension with R20 within the 1σ , 2σ and 3σ confidence levels considering the *Planck* dataset only.



Cosmic neuTRIVIA game steps

1. Familiarize yourself with the board's layout:

- The Λ CDM trivia: the players!
- The neutrino pie piece: decoupling in the early universe

2. Roll the dice and get:

- Number of neutrinos and Big-Bang Nucleosynthesis
- Number of neutrinos and Cosmic Microwave Background Radiation
- Neutrino masses and Cosmic Microwave Background Radiation?
- Neutrino masses and structure formation in the universe?

3. Is anyone cheating? Neutrinos and Tensions

4. Final score:

- Take home messages

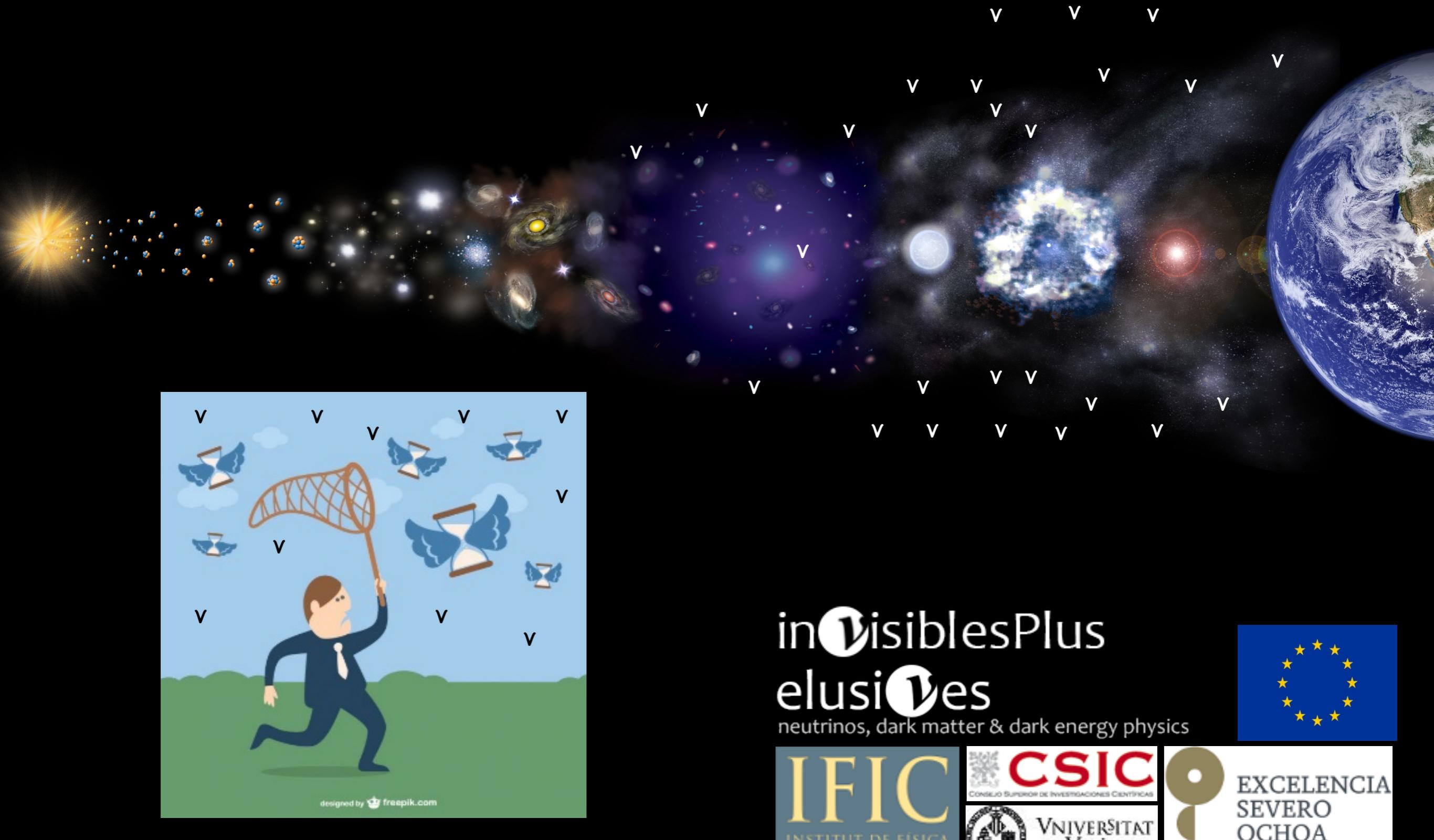
The “Take Home” messages

- ν masses & abundances leave key signatures in cosmological observables.
- NO hints so far for neutrino masses or extra dark radiation species!
- N_{eff} @BBN: Light element abundances (${}^4\text{He}$) abundances.
- N_{eff} @CMB: damping tail
- $N_{\text{eff}} = 2.99^{+0.34}_{-0.33}$, (95% CL) from 2018 Planck TTTEEE+lensing, perfectly consistent with BBN.
- Cosmology provides currently the tightest bounds to neutrino masses.
- ν masses@CMB: Early ISW, gravitational lensing
- ν masses@LSS: Free streaming
- $\sum m_\nu < 0.099 \text{ eV (95%CL)}$ from 2018 Planck TTTEEE+lensing plus RSD+BAO +SNIa data





BACKUP SLIDES



The Hubble parameter compendium

Knox & Millea, PRD'20

Pre-recombination solutions

Low sound horizon solutions

Confusions in determinations of
 w_m (neutrino interactions,
modify gravity, different $P(k)$,...)

Sound speed reduction

Higher recombination temperature

Photon cooling previous to recombination

Increasing $H(z)$ with additional components
(e.g. via N_{eff} or early dark energy)

$$\theta_s \equiv r_s / D_A$$

Post-recombination solutions

High sound horizon solutions

Wiggles in $H(z)$

Violation in the distance-duality
relation (axion dimming)

Cepheid mis-calibration

Low sound horizon solutions

Confusions in determinations of
 w_b w_m
Post recombination evolution of
 r_s

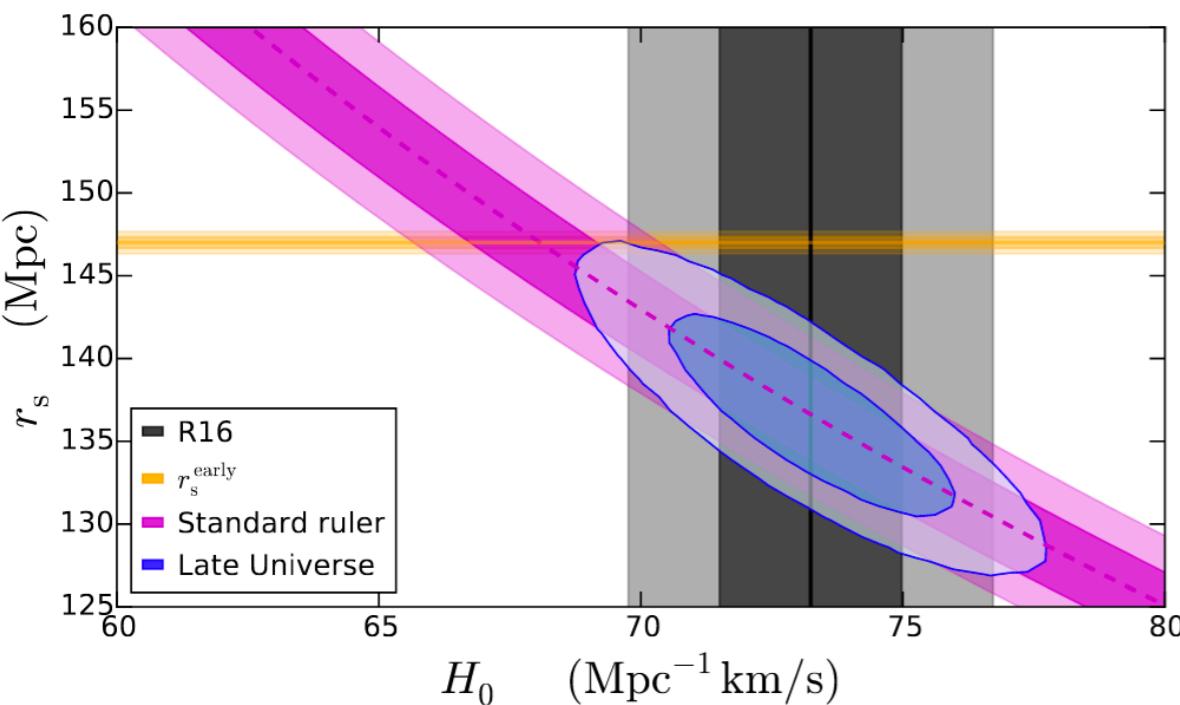
All solutions are required to leave unchanged θ_s , θ_{EQ} & θ_D

The ~~H~~ r_s tension?

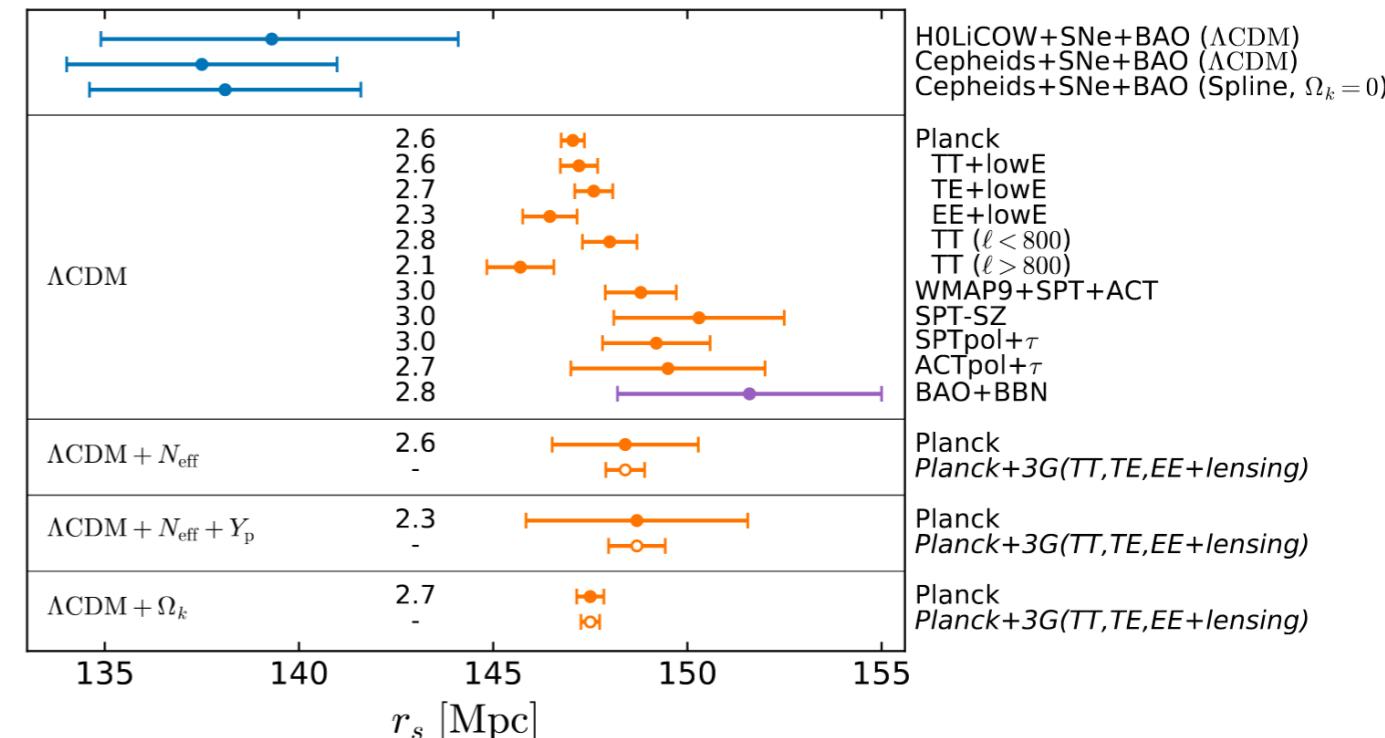
BAO observations determine D_V/r_s (or Hr_s & $D_A r_s$). One can use BAO, SNIa & Cepheids data to infer r_s :

$$r_s = 137.7 \pm 3.6 \text{ Mpc}$$

Casting the tension between Cosmic distance leaders and Λ CDM + CMB datasets in terms of r_s , **weakens the statistical significance**, but helps to clarify the physics that could reconcile these datasets.



Bernal, Verde & Riess, JCAP'16



Aylor et al, APJ'19

Baryon Acoustic Oscillations

2020= SDSS IV

Zhao et al, SDSS IV Coll. MNRAS'21

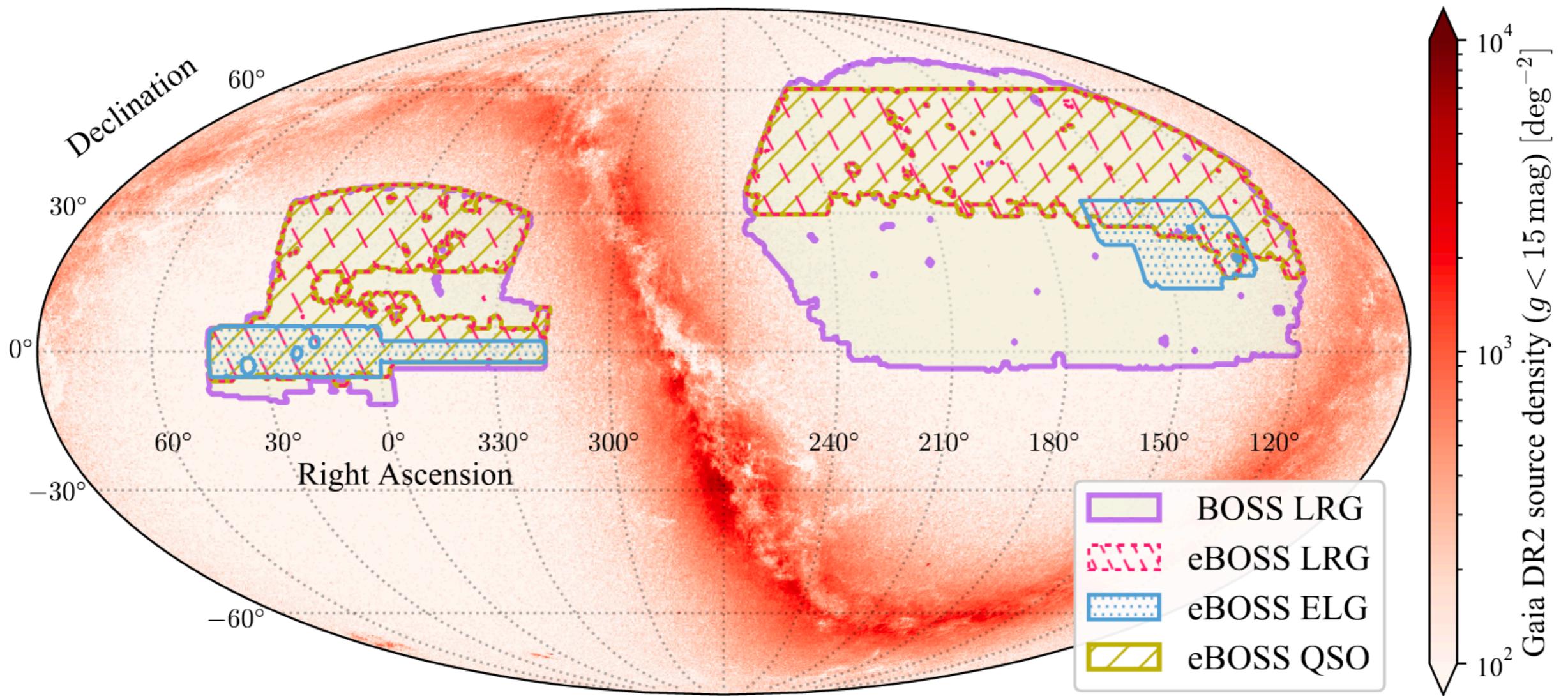


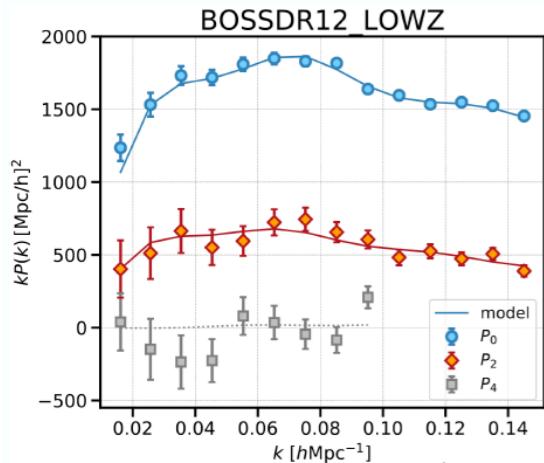
Figure 1. The sky coverage of eBOSS DR16 tracers and BOSS DR12 LRGs, as well as the density map of Gaia DR2 sources with $g < 15$ mag.

Baryon Acoustic Oscillations

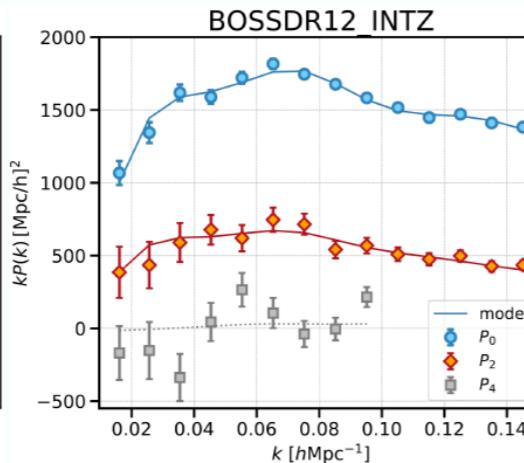
Hector Gil-Marín (ICCUB)

Cosmology from galaxy redshift surveys: current results and future prospects

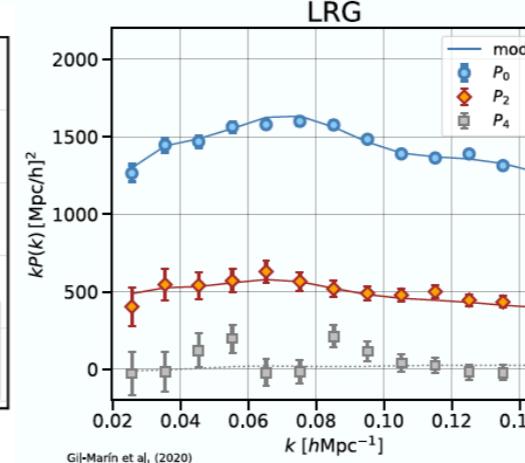
$0.2 < z < 0.5$



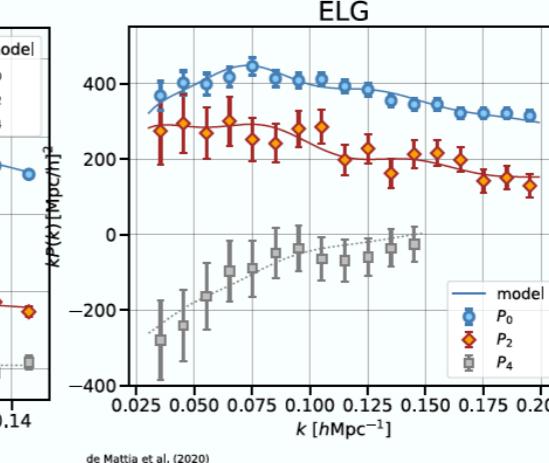
$0.4 < z < 0.6$



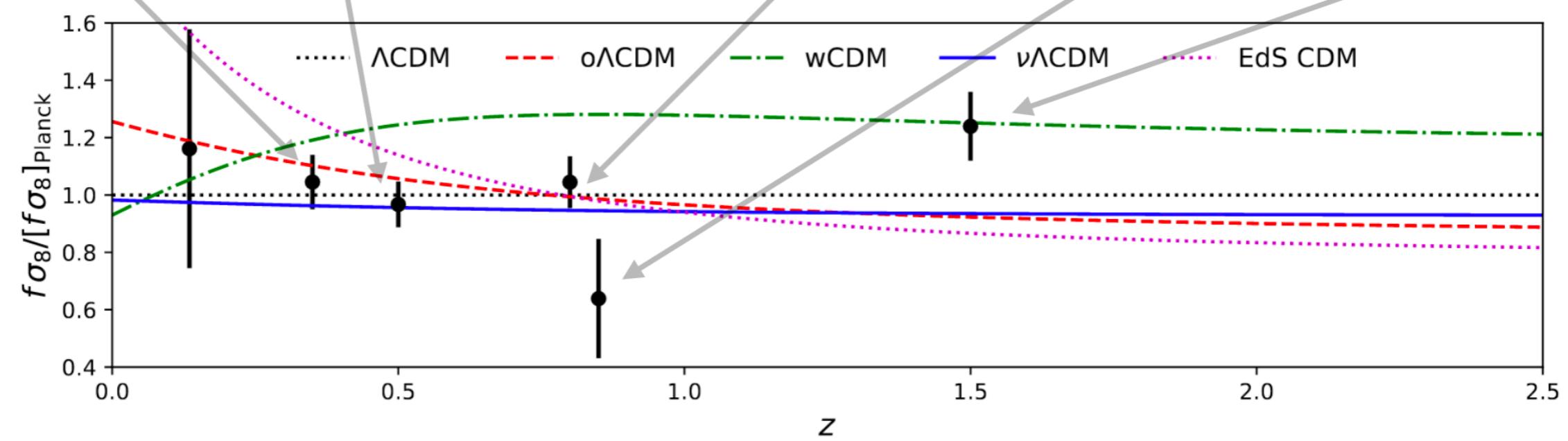
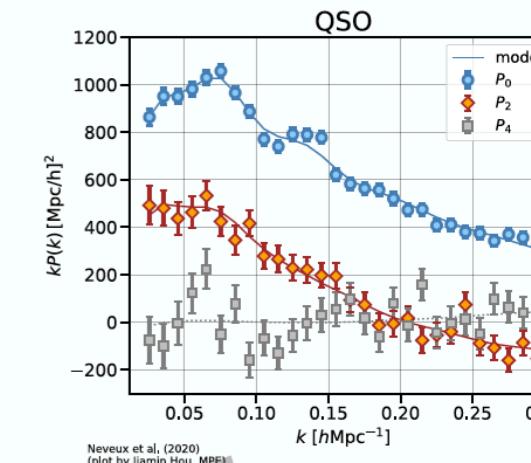
$0.6 < z < 1.0$



$0.6 < z < 1.1$



$0.8 < z < 2.2$



A gravitational-wave standard siren measurement of the Hubble constant

The LIGO Scientific Collaboration and The Virgo Collaboration, The 1M2H Collaboration, The Dark Energy Camera GW-EM Collaboration and the DES Collaboration, The DLT40 Collaboration, The Las Cumbres Observatory Collaboration, The VINROUGE Collaboration & The MASTER Collaboration

[Affiliations](#) | [Contributions](#) | [Corresponding authors](#)

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Citation



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Article metrics

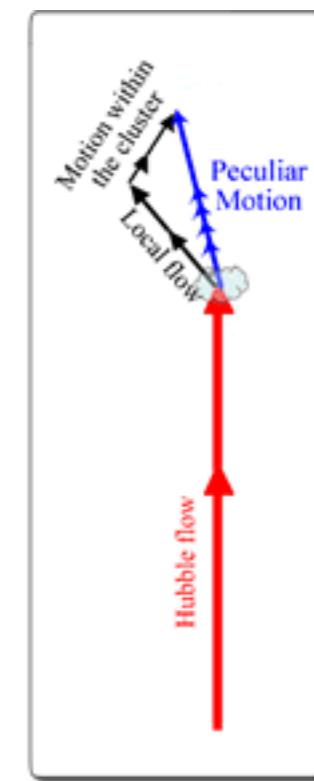
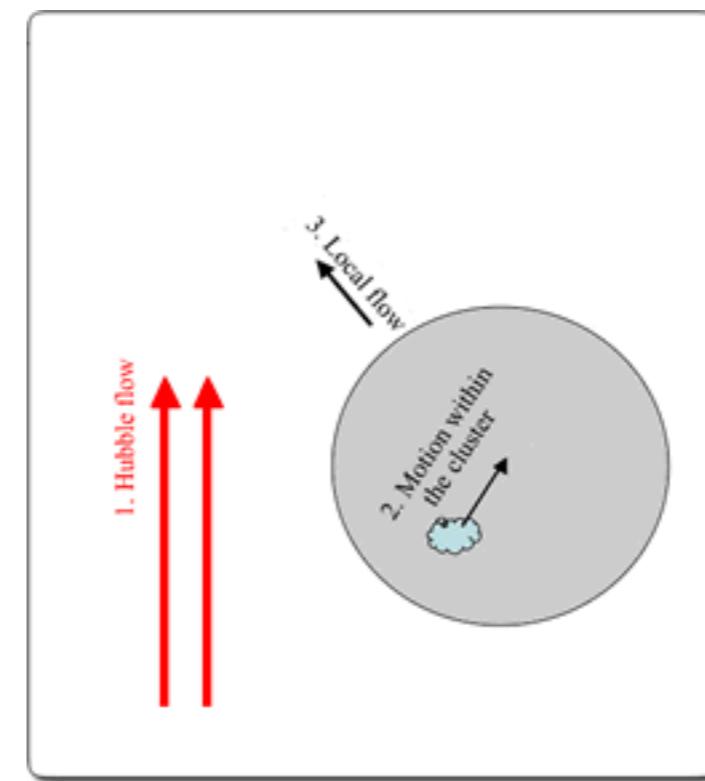
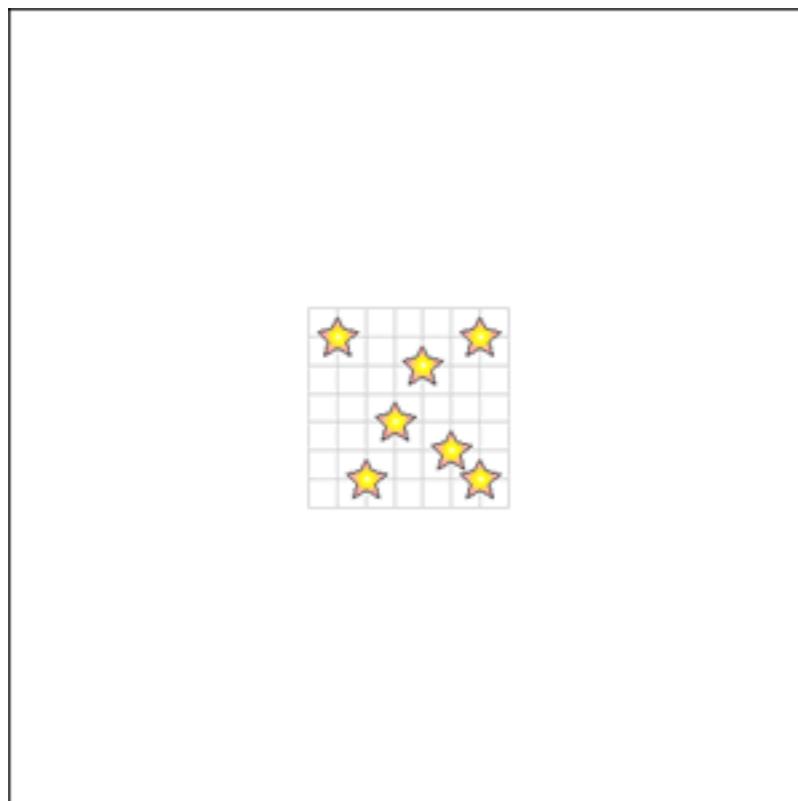
On 17 August 2017, the Advanced LIGO¹ and Virgo² detectors observed the gravitational-wave event GW170817—a strong signal from the merger of a binary neutron-star system³. Less than two seconds after the merger, a γ-ray burst (GRB 170817A) was detected within a region of the sky consistent with the LIGO–Virgo-derived location of the gravitational-wave source^{4, 5, 6}. This sky region was subsequently observed by optical astronomy facilities⁷, resulting in the identification^{8, 9, 10, 11, 12, 13} of an optical transient signal within about ten arcseconds of the galaxy NGC 4993. This detection of GW170817 in both gravitational waves and electromagnetic waves represents the first ‘multi-messenger’ astronomical observation. Such observations enable GW170817 to be used as a ‘standard siren’^{14, 15, 16, 17, 18} (meaning that the absolute distance to the source can be determined directly from the gravitational-wave measurements) to measure the Hubble constant. This quantity represents the local expansion rate of the Universe, sets the overall scale of the Universe and is of fundamental importance to cosmology. Here we report a measurement of the Hubble constant that combines the distance to the source inferred purely from the gravitational-wave signal with the recession velocity inferred from measurements of the redshift using the electromagnetic data. In contrast to previous measurements, ours does not require the use of a cosmic ‘distance ladder’¹⁹: the gravitational-wave analysis can be used to estimate the luminosity distance out to cosmological scales directly, without the use of intermediate astronomical distance measurements. We determine the Hubble constant to be about 70 kilometres per second per megaparsec. This value is consistent with existing measurements^{20, 21}, while being completely independent of them. Additional standard siren measurements from future gravitational-wave sources will enable the Hubble constant to be constrained to high precision.

The method combines the **distance to the source** inferred purely from the **gravitational-wave** signal with the **recession velocity** inferred from measurements of the redshift using **electromagnetic data**.

$$v_H = H_0 d$$

$$d = 43.8^{+2.9}_{-6.9} \text{ Mpc}$$

Using the optical identification of the host galaxy NGC 4993, they derive the Hubble flow velocity. **PROBLEM:** the random relative motion of galaxies (peculiar velocity) needs to be taken into account! In practice, the motions of galaxies are influenced by more than just the Hubble flow: the local flow, and the motion of the galaxy within its cluster and/or group environment. These deviations from the pure Hubble flow are **referred to as peculiar motions**. The peculiar velocity is about 10% of the measured recessional velocity.



The **Hubble flow** causes all **galaxies** to recede from each other.

The **local flow** and the motion of the galaxy within its **cluster environment** also contribute.

NGC 4993 is part of a collection of galaxies, ESO 508, which has a center-of-mass recession velocity relative to the frame of the cosmic CMB of 3327 ± 72 km/s. The authors correct the group velocity by 310 km/s, due to the local gravitational fields.

The standard error on their estimate of the peculiar velocity is 69 km/s, but recognizing that this value may be sensitive to details of the bulk flow motion, in their analysis adopt a more conservative estimate of 150 km/s for the uncertainty on the peculiar velocity at the location of NGC 4993 and fold this in their estimate of the uncertainty on v_H . From this, they obtain a Hubble velocity $v_H = 3017 \pm 166$ km/s.

Using this recessional velocity, one can find $H_0 = 68.9$ km/s.

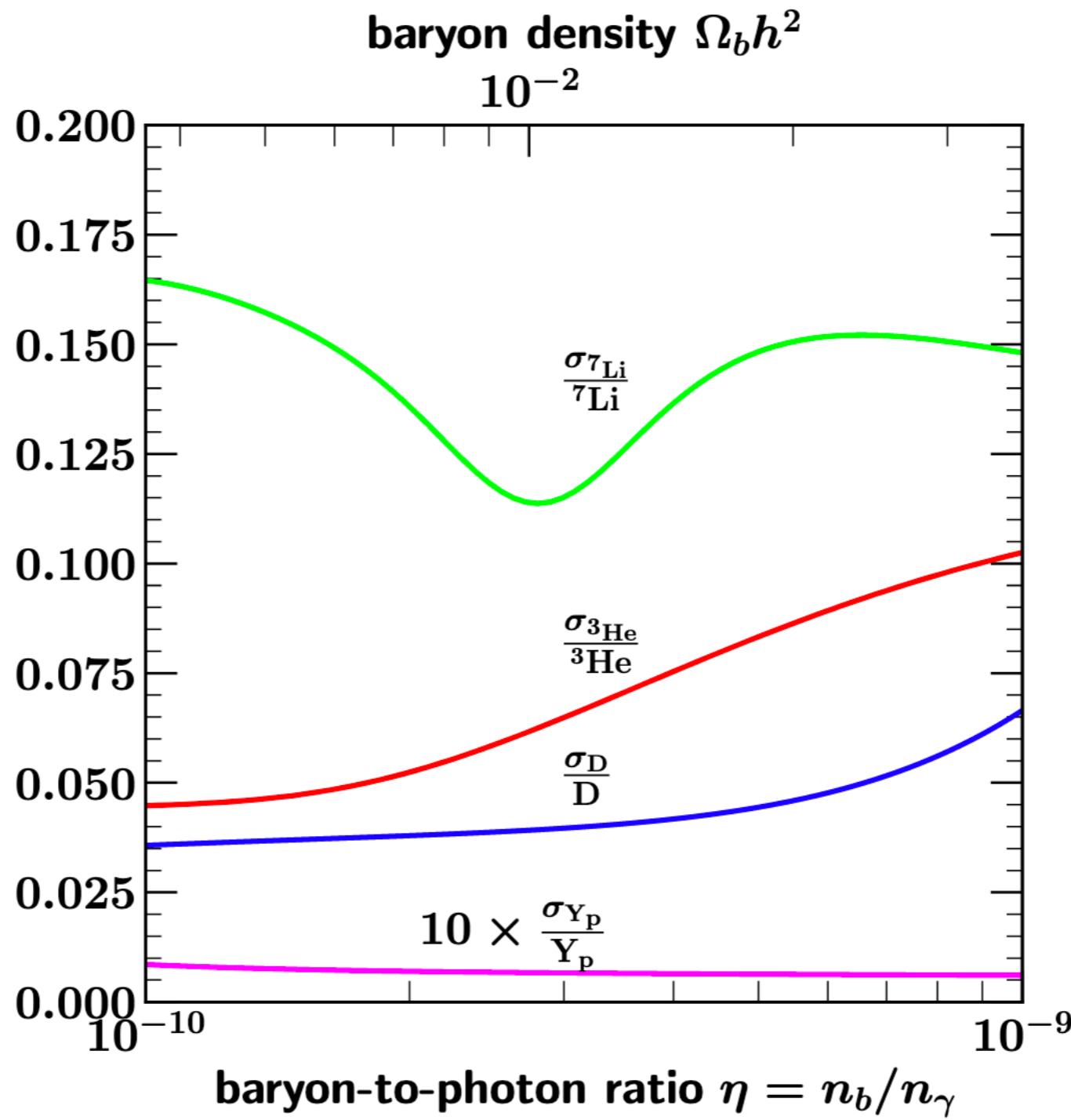


FIG. 3. Fractional uncertainties in the light element abundance predictions shown in Fig. 2. For each species i , we plot ratio of the standard deviation σ_i to the mean μ_i , as a function of baryon-to-photon ratio. The relative uncertainty of the ${}^4\text{He}$ abundance has been multiplied by a factor of 10.

Standard-model corrections to $N_{\text{eff}}^{\text{SM}}$	Leading-digit contribution
m_e/T_d correction	+0.04
$\mathcal{O}(e^2)$ FTQED correction to the QED EoS	+0.01
Non-instantaneous decoupling+spectral distortion	-0.005
$\mathcal{O}(e^3)$ FTQED correction to the QED EoS	-0.001
Flavour oscillations	+0.0005
Type (a) FTQED corrections to the weak rates	$\lesssim 10^{-4}$

Bennett et al, 2012.02726

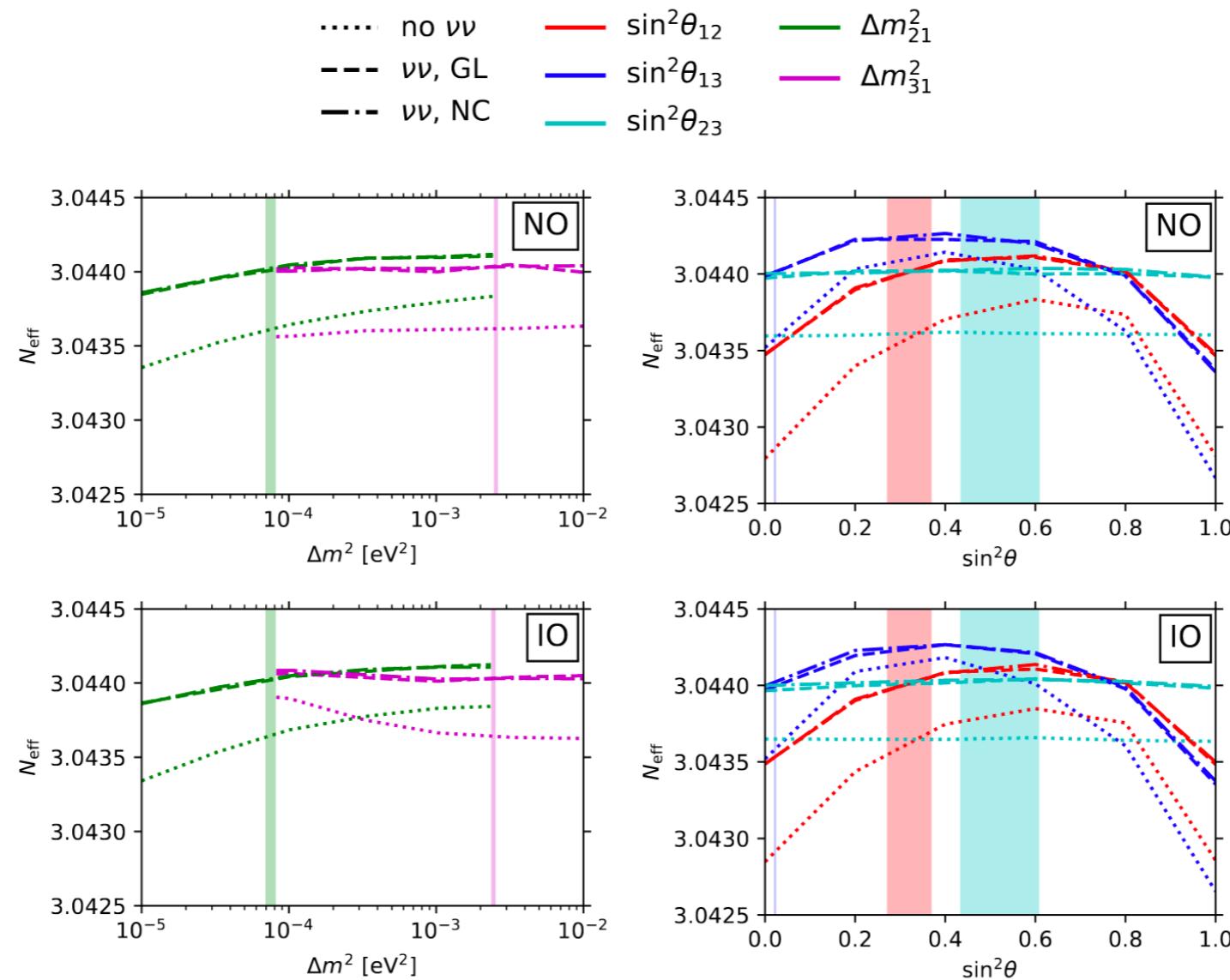
The ultra relativistic approximation:

$$T_d/m_e \rightarrow \infty$$

is not well satisfied in reality!

10^{-4} Uncertainty due to measurement errors on the solar mixing angle

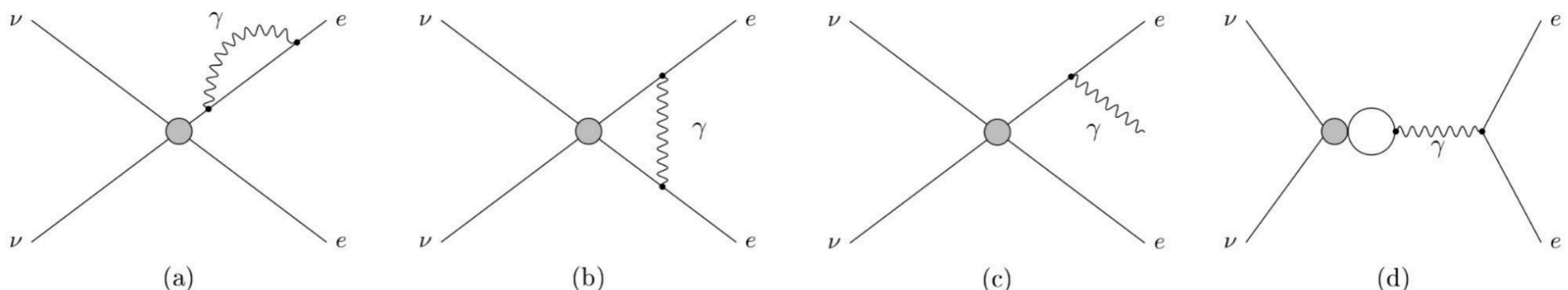
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$\mathcal{O}(e^5)$ FTQED correction to the QED EoS	-0.001
Flavour oscillations	+0.0005
Type (a) FTQED corrections to the weak rates	$\lesssim 10^{-4}$

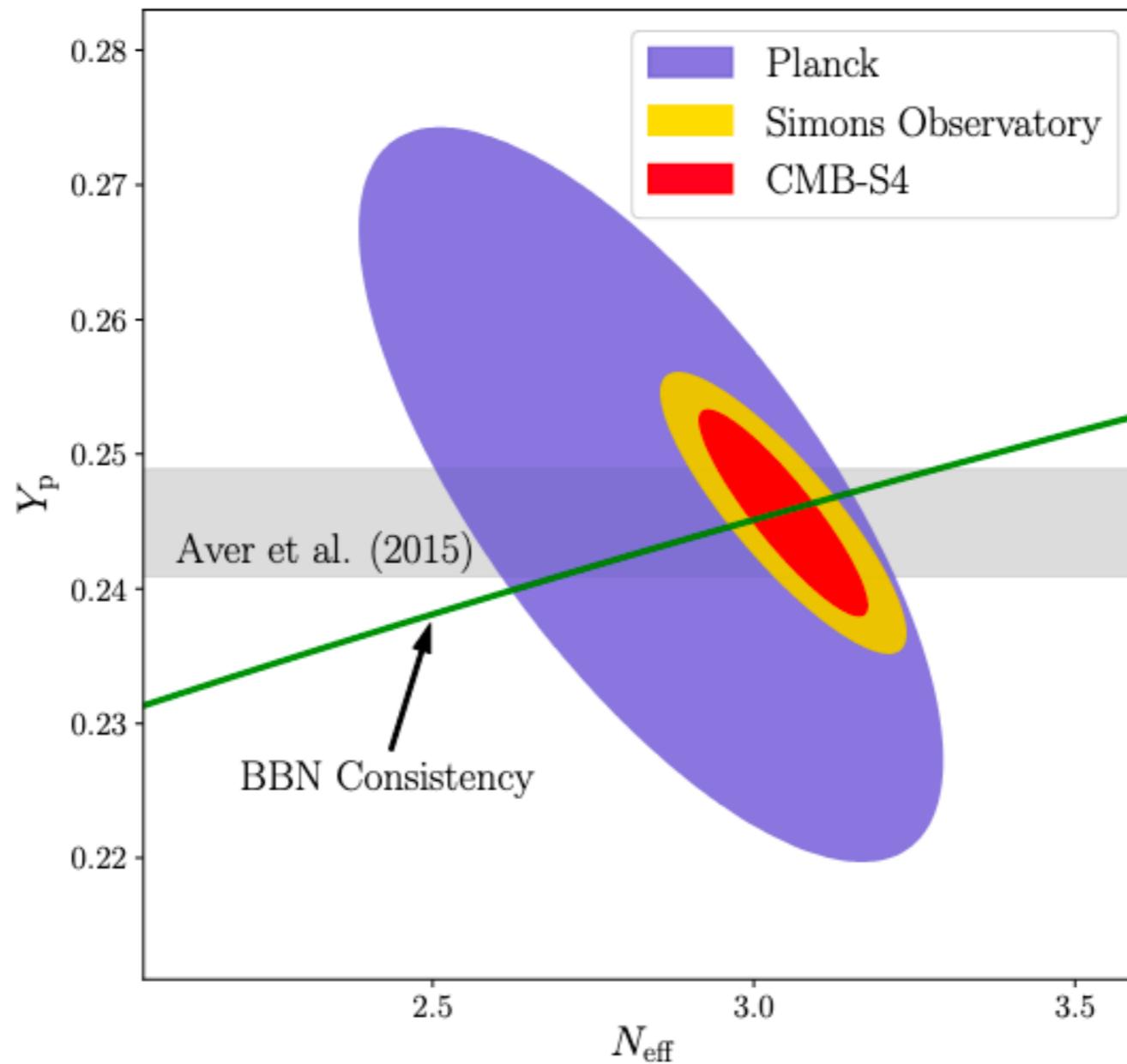
Bennett et al, 1911.04504

$$\ln Z^{(2)} + \ln Z^{(3)} = -\frac{1}{2} \cdot \text{(loop diagram)} + \frac{1}{2} \left[\frac{1}{2} \cdot \text{(diagram with 2 loops)} - \frac{1}{3} \cdot \text{(diagram with 3 loops)} + \frac{1}{4} \cdot \text{(diagram with 4 loops)} + \dots \right]$$



CMB Stage IV: N_{eff}

$$\Delta N_{\text{eff}} < 0.06 \text{ 95%CL}$$



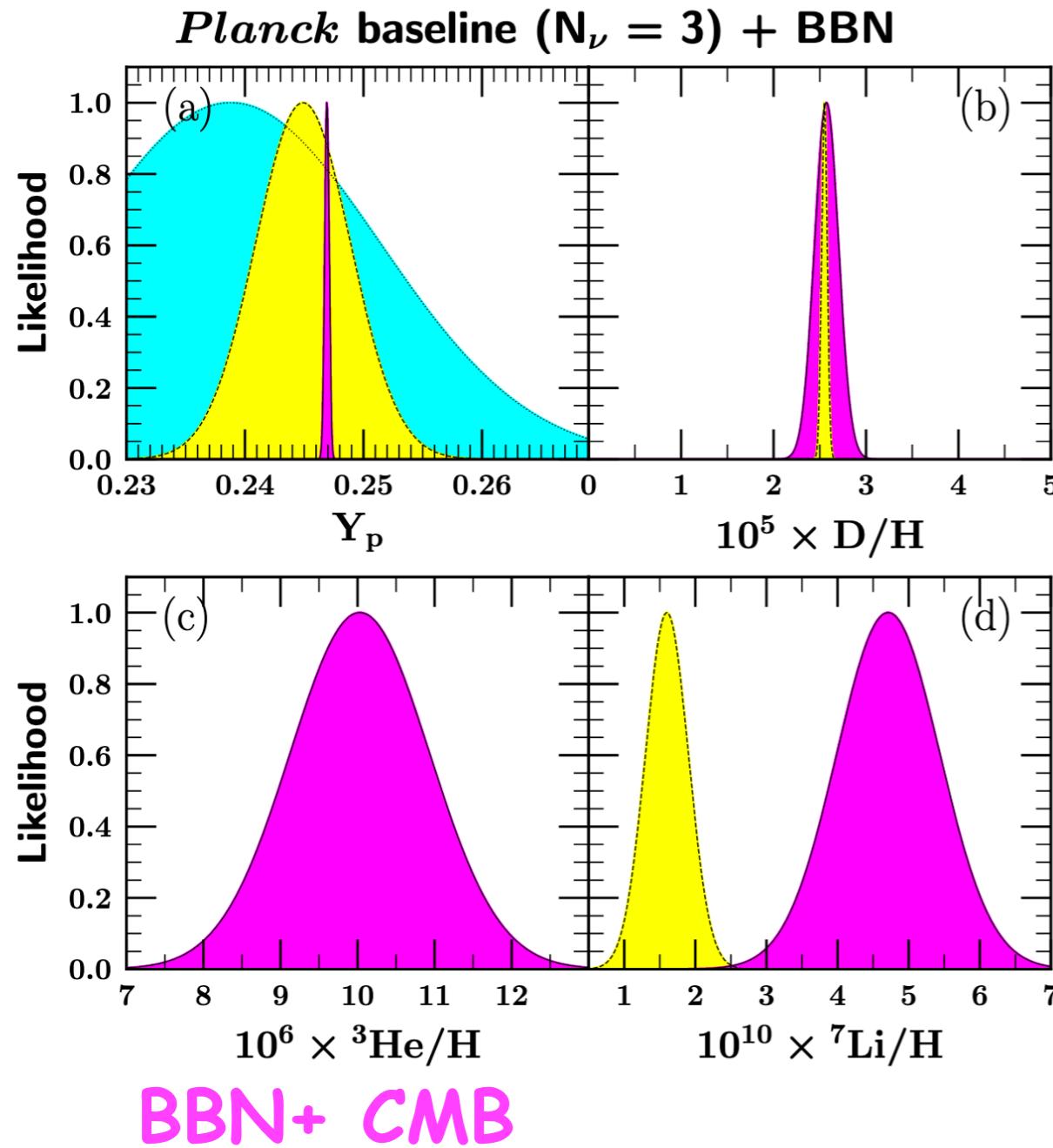
CMB-S4 Science Case, Reference Design, and Project Plan, 1907.04473

CMB: N_{eff}

$$Y_p = 0.24691 \pm 0.00018 \quad (0.24691)$$

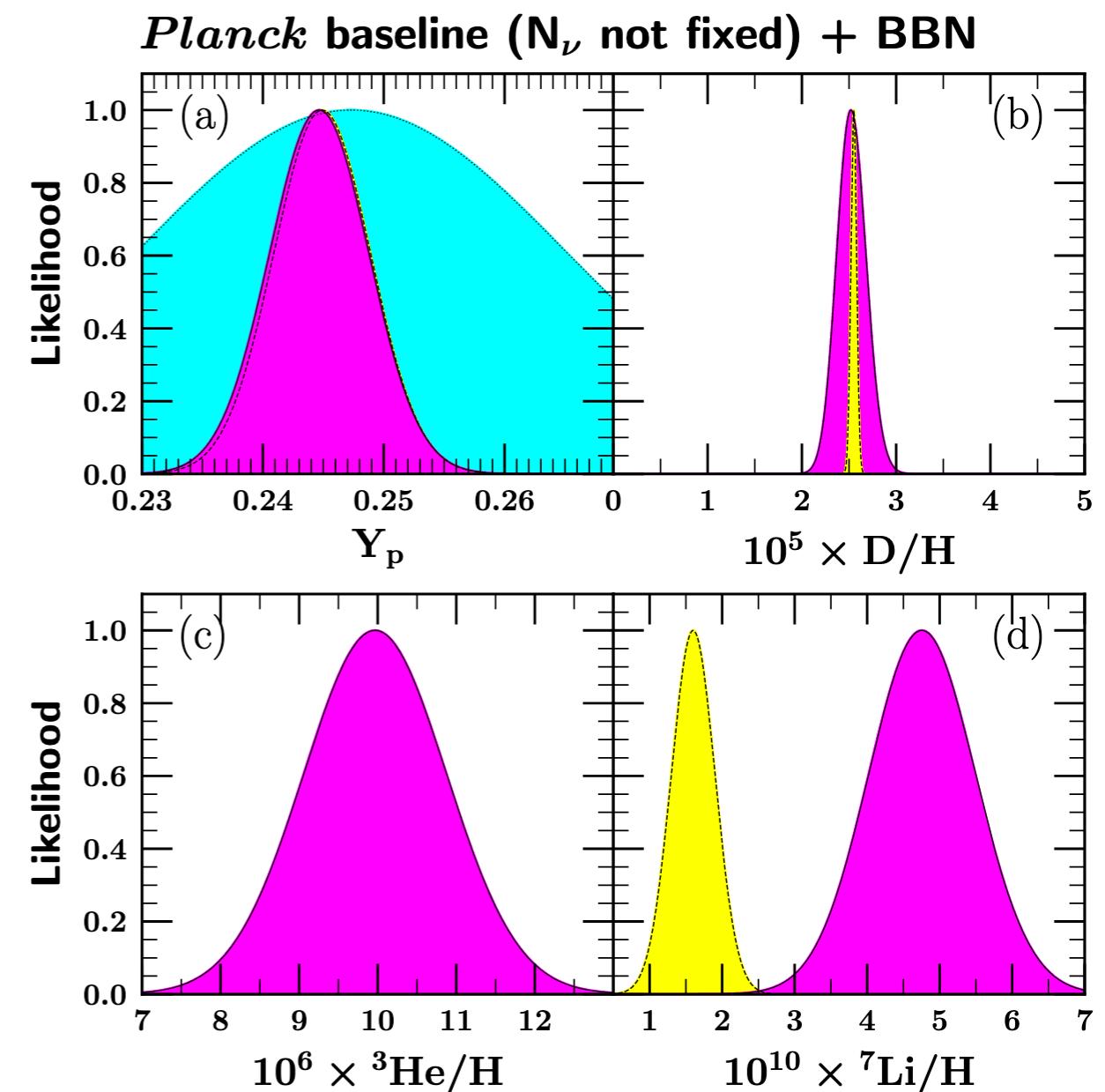
$$Y_p = 0.24465 \pm 0.00410 \quad (0.24498)$$

Fields, Olive, Yeh & Young JCAP '20

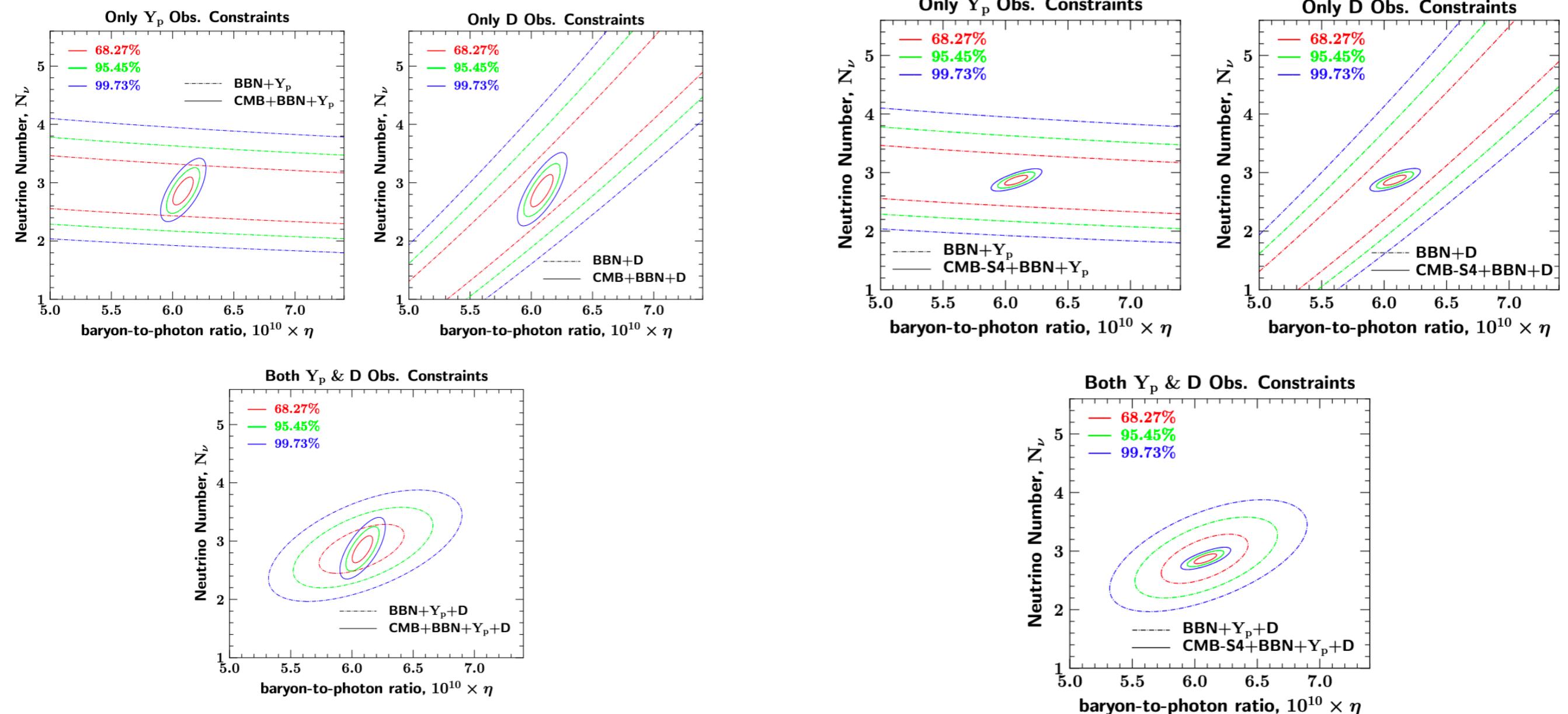


BBN+ CMB

Astronomical measurements
CMB He determinations

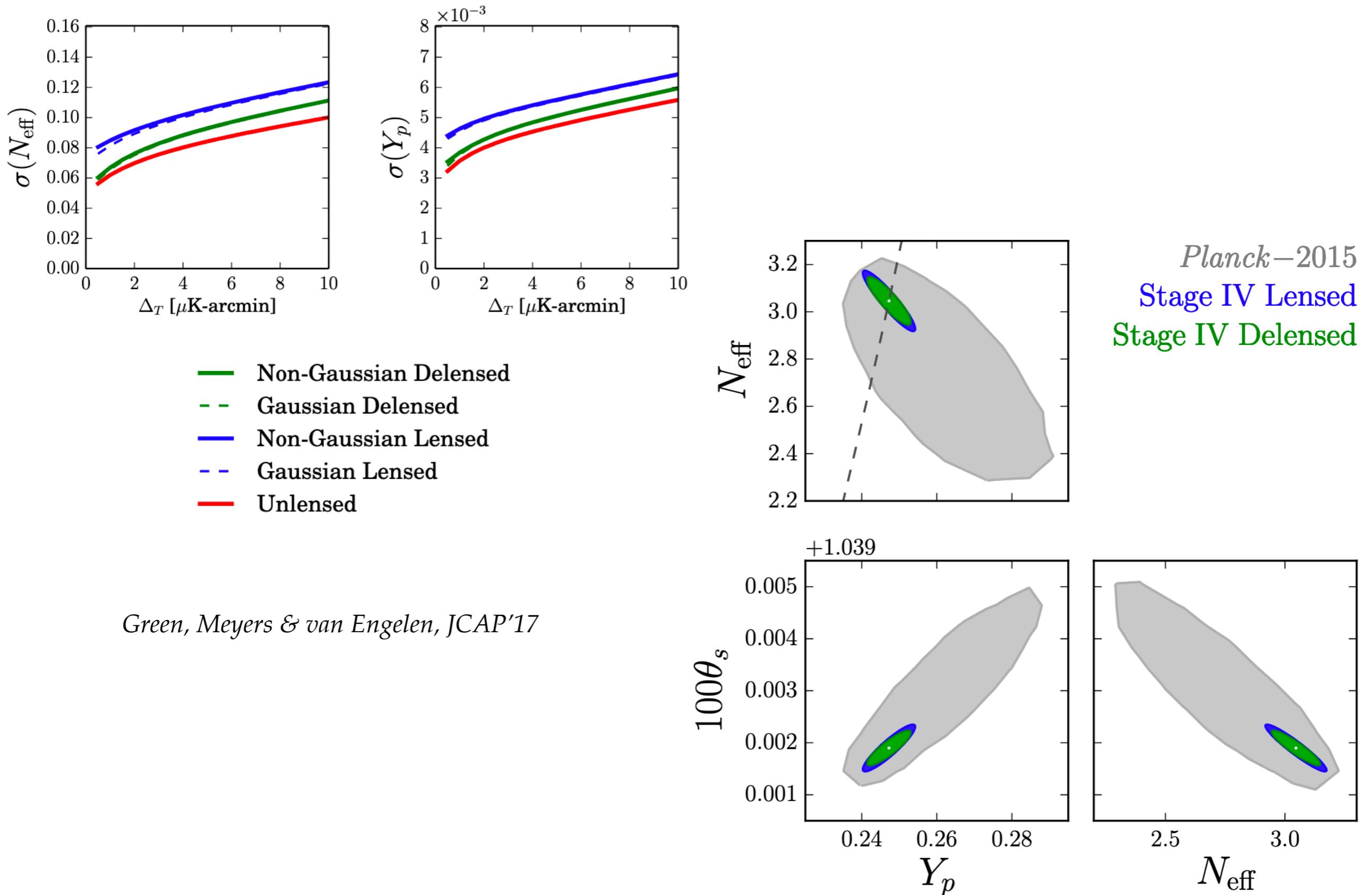


CMB Stage IV: N_{eff}



Fields, Olive, Yeh & Young JCAP '20

CMB Stage IV: N_{eff}



B. Strongly interacting neutrino mode

The existence of the $\text{SI}\nu$ mode was first pointed out in Ref. [55], and further studied in Refs. [65, 66]. As discussed there, the $\text{SI}\nu$ cosmology arises due to a multi-parameter degeneracy that opens up in CMB data when the onset of neutrino free-streaming is delayed until redshift $z \sim 8000$. This approximately coincides with the epoch when Fourier modes corresponding to multipole $\ell \approx 400$ enters the causal horizon [65], which lies somewhere between the first and second peak of the CMB temperature spectrum. We review below the properties of this alternate cosmology, emphasizing its differences with the standard ΛCDM model.

neutrinos means that the CMB spectra do not receive the standard phase shift, and thus appear slightly displaced toward larger ℓ as compared to the corresponding ΛCDM spectra. In order to fit the data, we must compensate for this shift by *increasing* the value of θ_* . Thus, the difference between the values of θ_* in the $\text{SI}\nu$ and ΛCDM models directly reflects the absence of the free-streaming neutrino phase shift in the former.

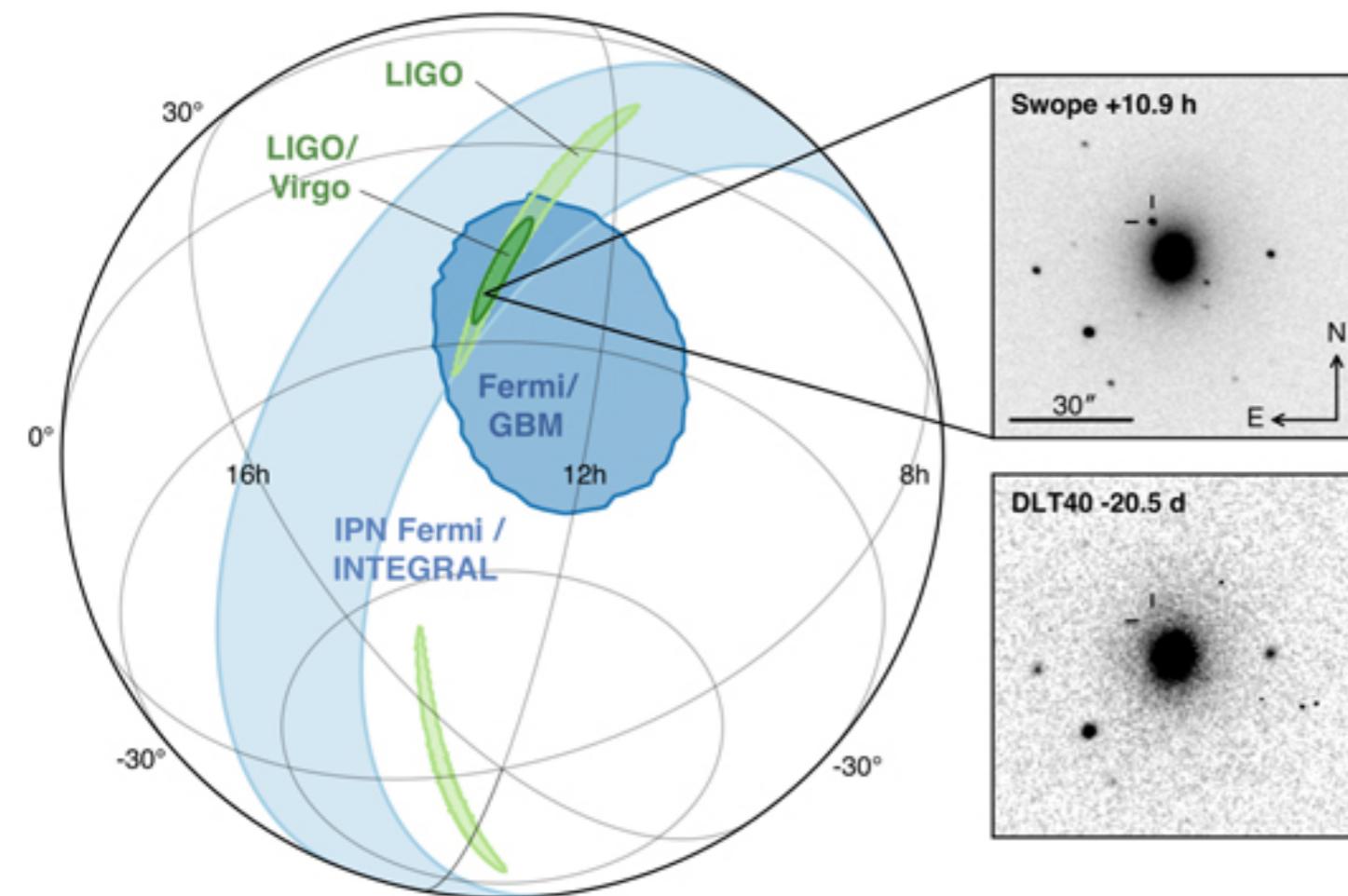
We note that it was a priori far from obvious that such a dramatic change in the angular size of the sound horizon was possible without introducing other artifacts that would significantly worsen the fit to CMB and BAO data. Our analysis shows that the larger value of θ_* is achieved by increasing H_0 and $\Omega_{\text{c}}h^2$ above their ΛCDM values.

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We note that it was a priori far from obvious that such a dramatic change in the angular size of the sound horizon was possible without introducing other artifacts that would significantly worsen the fit to CMB and BAO data. Our analysis shows that the larger value of θ_* is achieved by increasing H_0 and $\Omega_{\text{c}}h^2$ above their ΛCDM values.

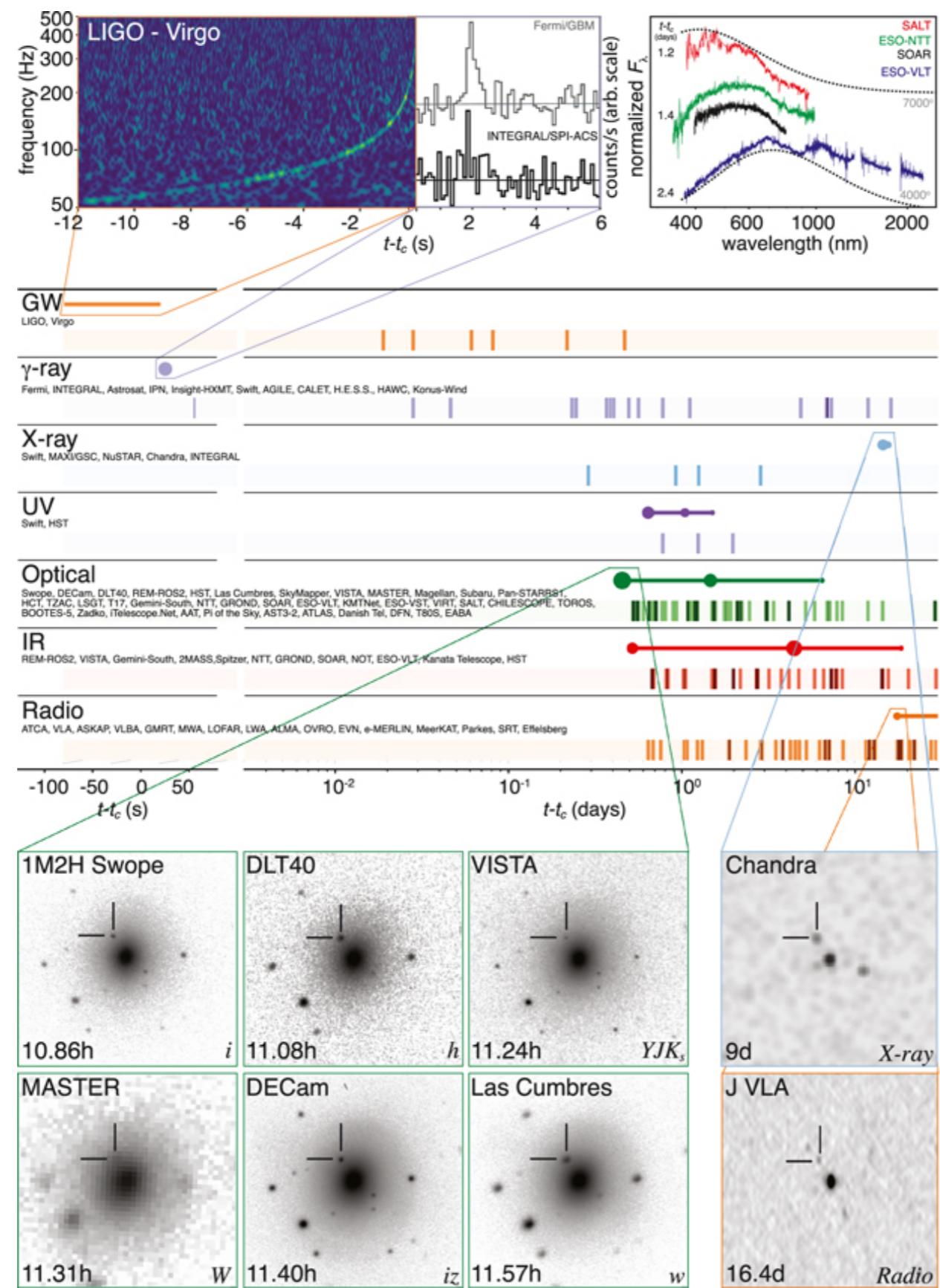
GW170817

The detection of GW170817 in both gravitational waves and electromagnetic waves represents the first ‘multi-messenger’ astronomical observation



Localization of the gravitational-wave, gamma-ray, and optical signals.
The left panel shows an orthographic projection of the 90% credible regions from LIGO (light green), the initial LIGO-Virgo localization (dark green), IPN triangulation from the time delay between Fermi and INTEGRAL (light blue), and Fermi-GBM (dark blue). The inset shows the location of the apparent host galaxy NGC 4993 in the Swope optical discovery image at 10.9 hr after the merger (top right) and the DLT40 pre-discovery image from 20.5 days prior to merger (bottom right). The reticle marks the position of the transient in both images.

GW170817



A gravitational-wave standard siren measurement of the Hubble constant

The LIGO Scientific Collaboration and The Virgo Collaboration, The 1M2H Collaboration, The Dark Energy Camera GW-EM Collaboration and the DES Collaboration, The DLT40 Collaboration, The Las Cumbres Observatory Collaboration, The VINROUGE Collaboration & The MASTER Collaboration

[Affiliations](#) | [Contributions](#) | [Corresponding authors](#)

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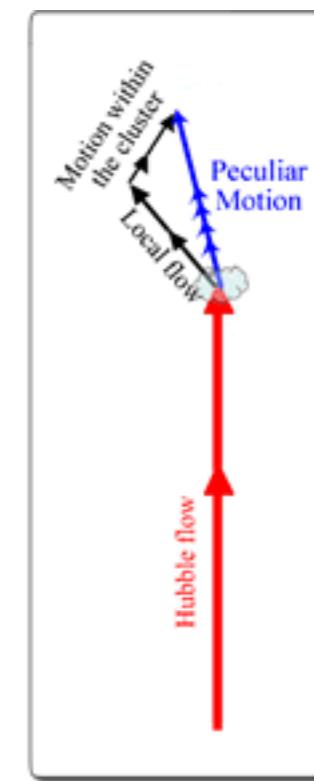
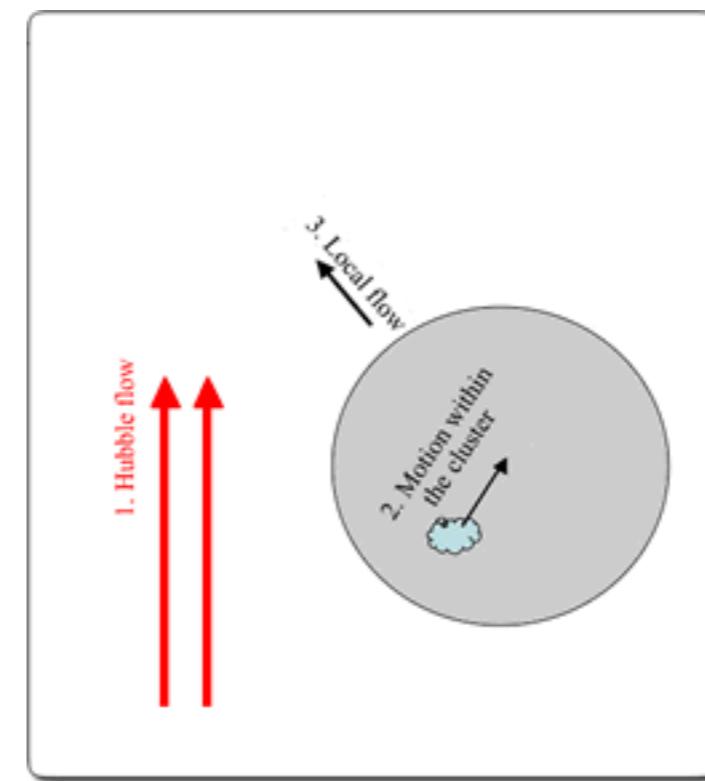
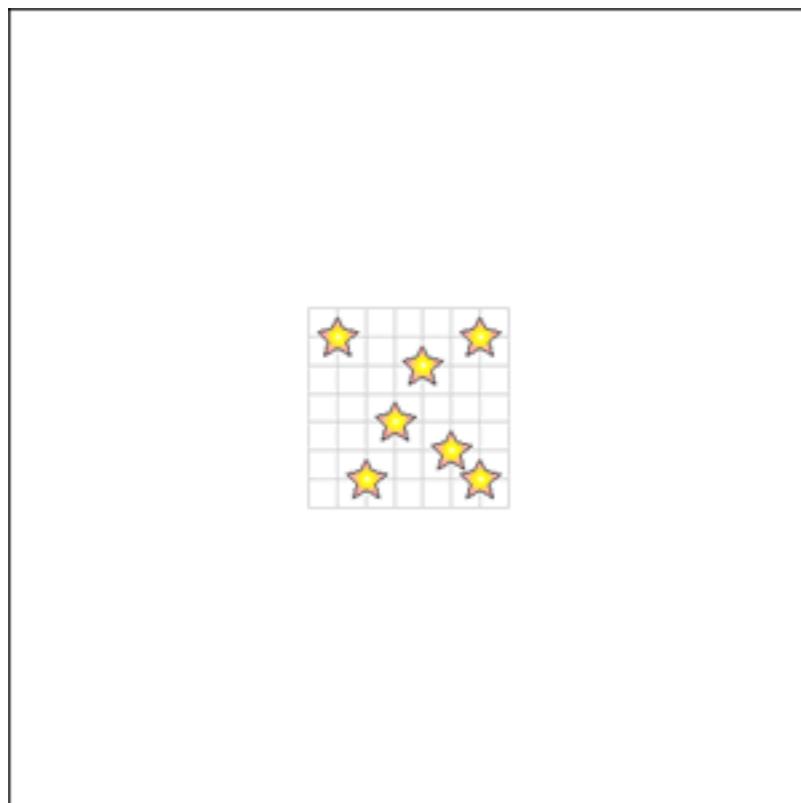
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The method combines the **distance to the source** inferred purely from the **gravitational-wave** signal with the **recession velocity** inferred from measurements of the redshift using **electromagnetic data**.

$$v_H = H_0 d$$

$$d = 43.8^{+2.9}_{-6.9} \text{ Mpc}$$

Using the optical identification of the host galaxy NGC 4993, they derive the Hubble flow velocity. **PROBLEM:** the random relative motion of galaxies (peculiar velocity) needs to be taken into account! In practice, the motions of galaxies are influenced by more than just the Hubble flow: the local flow, and the motion of the galaxy within its cluster and/or group environment. These deviations from the pure Hubble flow are **referred to as peculiar motions**. The peculiar velocity is about 10% of the measured recessional velocity.



The Hubble flow causes all **galaxies** to recede from each other.

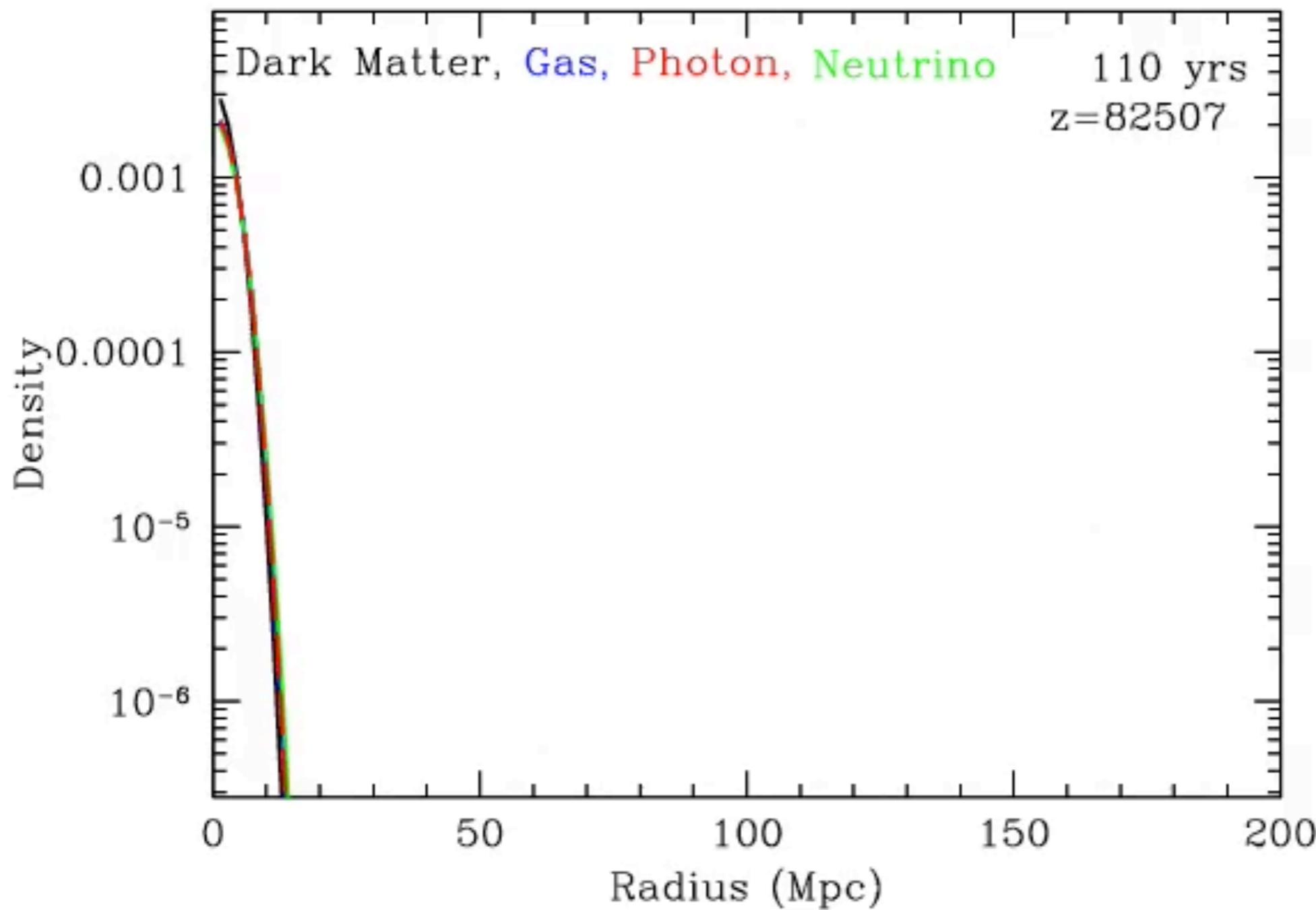
The **local flow** and the motion of the galaxy within its **cluster environment** also contribute.

NGC 4993 is part of a collection of galaxies, ESO 508, which has a center-of-mass recession velocity relative to the frame of the cosmic CMB of 3327 ± 72 km/s. The authors correct the group velocity by 310 km/s, due to the local gravitational fields.

The standard error on their estimate of the peculiar velocity is 69 km/s, but recognizing that this value may be sensitive to details of the bulk flow motion, in their analysis adopt a more conservative estimate of 150 km/s for the uncertainty on the peculiar velocity at the location of NGC 4993 and fold this in their estimate of the uncertainty on v_H . From this, they obtain a Hubble velocity $v_H = 3017 \pm 166$ km/s.

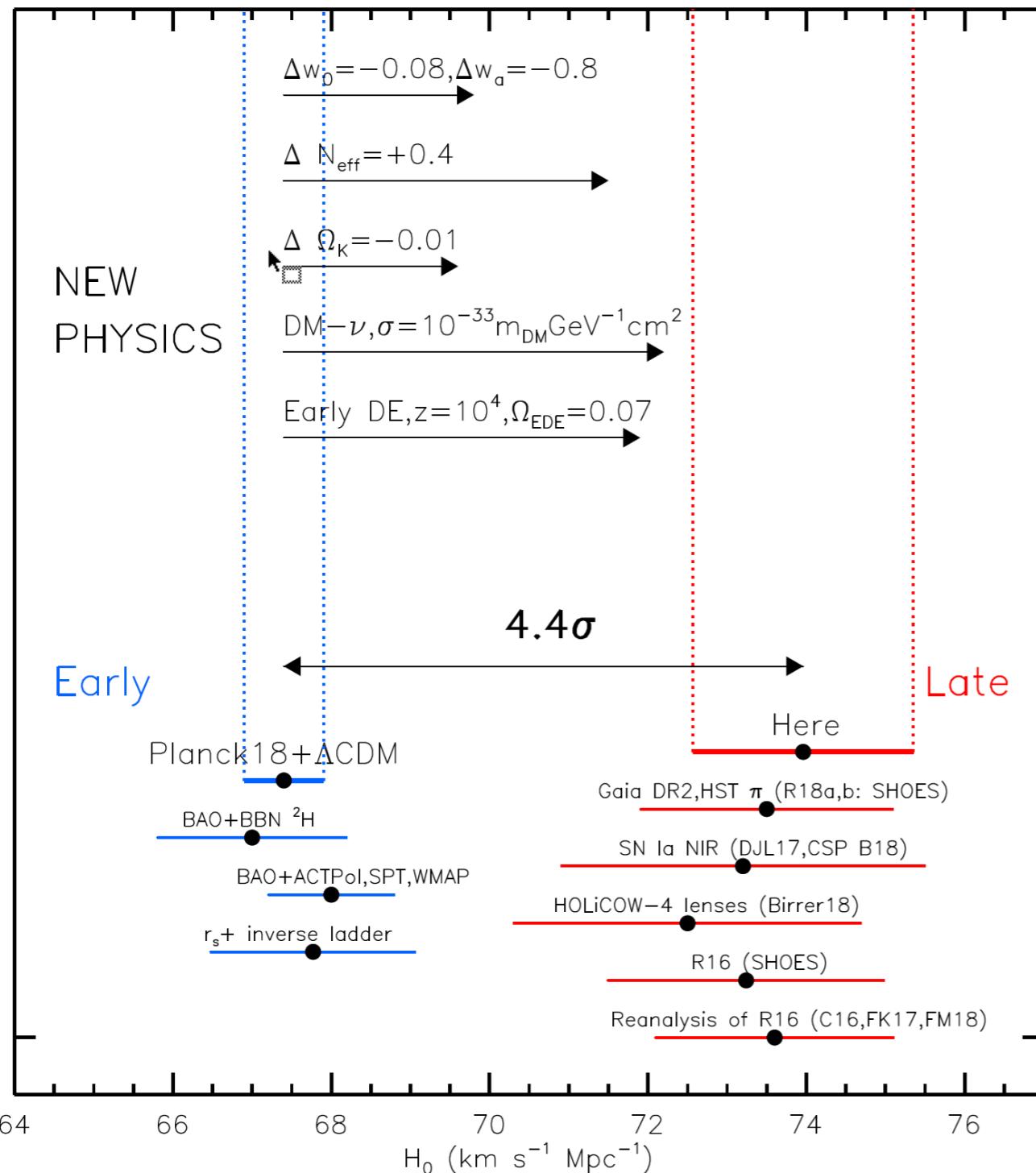
Using this recessional velocity, one can find $H_0 = 68.9$ km/s.

Baryon Acoustic Oscillations



From D. Eisenstein and M. White

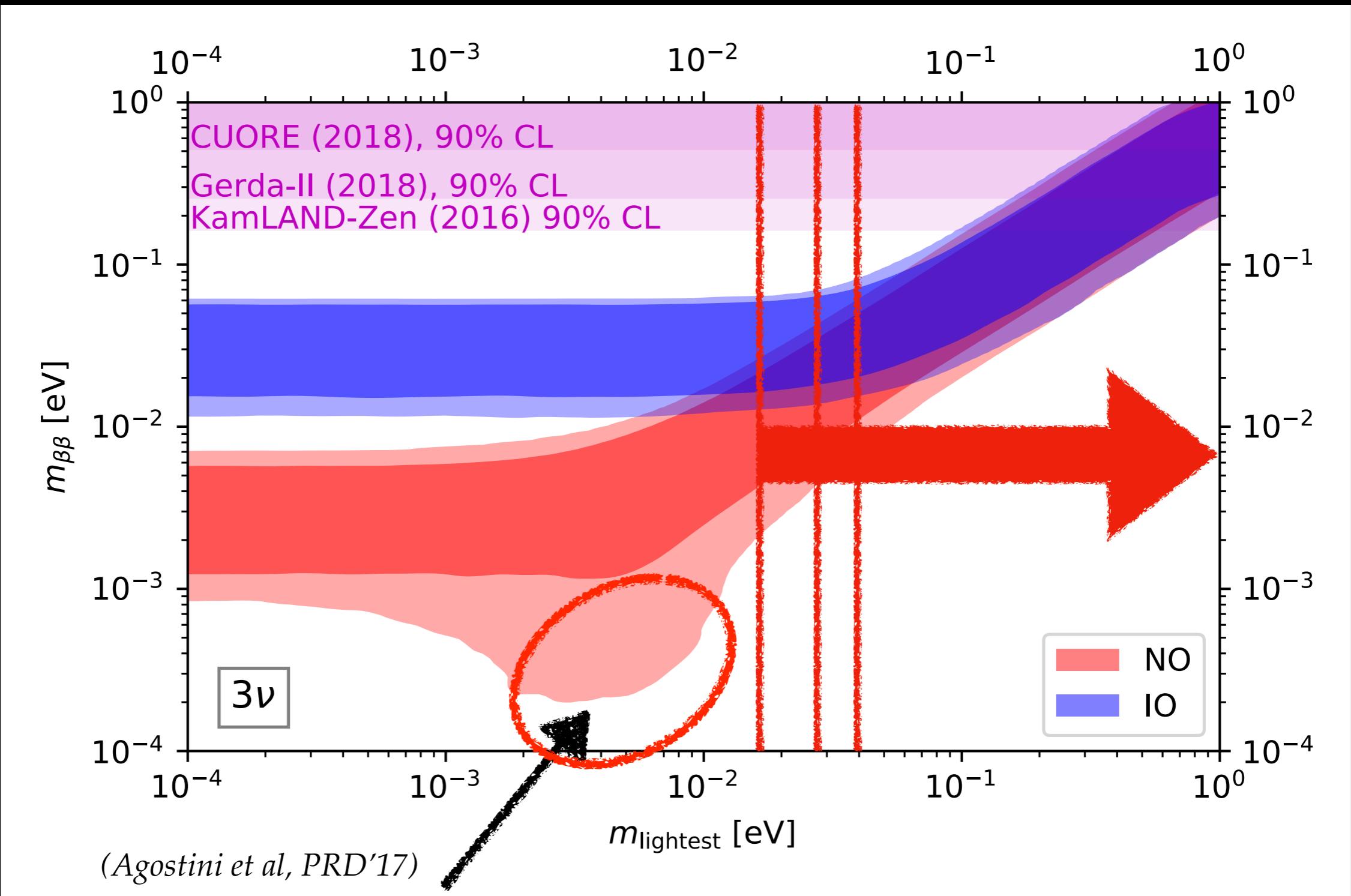
*A. G. Riess, S. Casertano, W. Yuan, L. M. Macri and
D. Scolnic, c.*



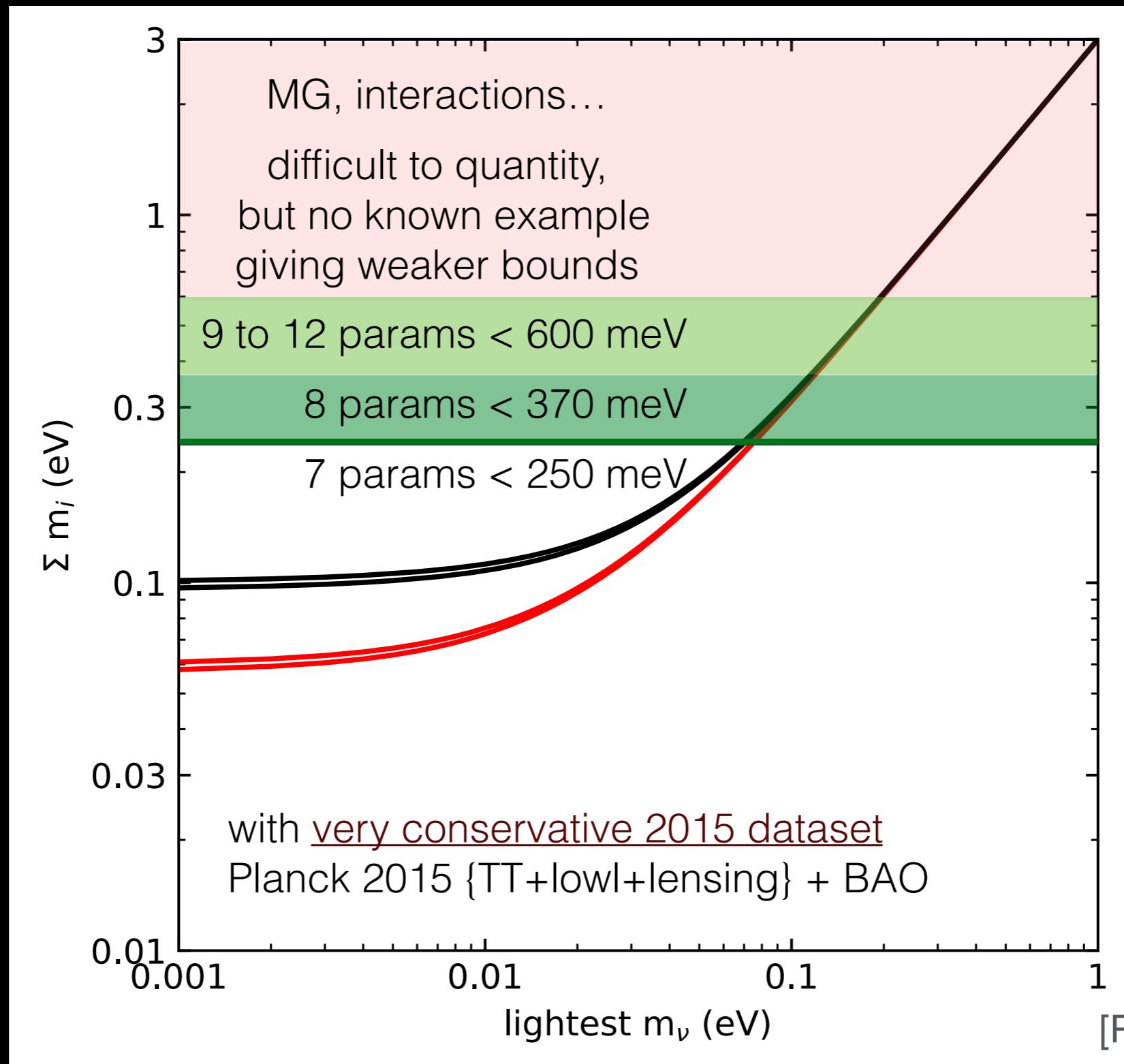
Decaying dark matter
Interacting dark energy
Curvature
Early dark energy
Modifications of gravity

$$H_0 = 67.37 \pm 0.54 \text{ km/s/Mpc} \quad H_0 = 74.03 \pm 1.42 \text{ km/s/Mpc}$$

What we know



$m_{\beta\beta} < 2 \cdot 10^{-4}$ would require
some fine tuning in the Majorana phases



Methods to detect non-non-relativistic neutrinos: PTOLEMY

Today neutrinos have a mean temperature:

$$T_{\nu 0} = \left(\frac{4}{11}\right)^{1/3} T_0 = 1.945 \text{ K} \sim 1.697 \times 10^{-4} \text{ eV}$$

And two neutrinos have a mass above: $\sqrt{\Delta m_{12}^2} \simeq 0.008 \text{ eV}$

at least one of these neutrinos has a mass larger than $\sqrt{|\Delta m_{13}^2|} \simeq 0.05 \text{ eV}$

Therefore there are at least two non-relativistic neutrino states.

A process without energy threshold is mandatory!

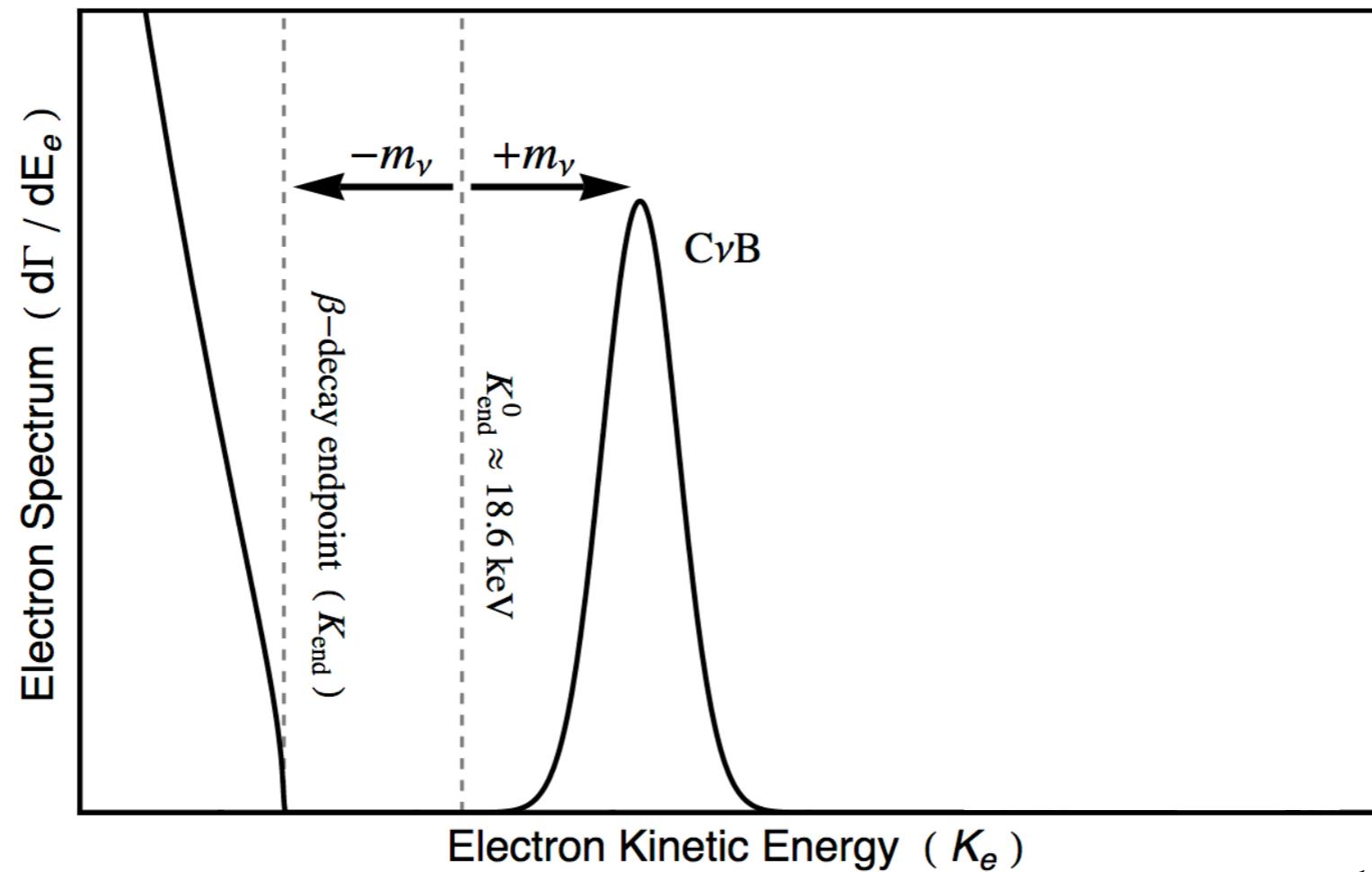
(Anti)neutrino capture on β -decaying nuclei



$$M(N) - M(N') = Q_\beta > 0$$

Methods to detect non-relativistic neutrinos:

(Anti)neutrino capture on β -decaying nuclei



Long, Lunardini & Sabancilar, JCAP'14

For finite m_ν , the electron kinetic energy is $Q_\beta + E_\nu \geq Q_\beta + m_\nu$, while electrons emerging from the analogous beta decay have at most an energy $Q_\beta - m_\nu$, neglecting nucleus recoil energy. A minimum gap of $2m_\nu$ is thus present and this at least in principle allows to distinguish between beta decay and NCB interaction: **GOOD ENERGY RESOLUTION!**

PTOLEMY (PonTecorvo Observatory for Light, Early-universe Massive-neutrino Yield) @ LNGS

PTOLEMY: A Proposal for Thermal Relic Detection of Massive Neutrinos
and Directional Detection of MeV Dark Matter

E. Baracchini³, M.G. Betti¹¹, M. Biasotti⁵, A. Boscá¹⁶, F. Calle¹⁶, J. Carabe-Lopez¹⁴, G. Cavoto^{10,11},
C. Chang^{22,23}, A.G. Cocco⁷, A.P. Colijn¹³, J. Conrad¹⁸, N. D'Ambrosio², P.F. de Salas¹⁷,
M. Faverzani⁶, A. Ferella¹⁸, E. Ferri⁶, P. Garcia-Abia¹⁴, G. Garcia Gomez-Tejedor¹⁵, S. Gariazzo¹⁷,
F. Gatti⁵, C. Gentile²⁵, A. Giachero⁶, J. Gudmundsson¹⁸, Y. Hochberg¹, Y. Kahn²⁶, M. Lisanti²⁶,
C. Mancini-Terracciano¹⁰, G. Mangano⁷, L.E. Marcucci⁹, C. Mariani¹¹, J. Martínez¹⁶, G. Mazzitelli⁴,
M. Messina²⁰, A. Molinero-Vela¹⁴, E. Monticone¹², A. Nucciotti⁶, F. Pandolfi¹⁰, S. Pastor¹⁷,
J. Pedrós¹⁶, C. Pérez de los Heros¹⁹, O. Pisanti^{7,8}, A. Polosa^{10,11}, A. Puiu⁶, M. Rajteri¹²,
R. Santorelli¹⁴, K. Schaeffner³, C.G. Tully²⁶, Y. Raitses²⁵, N. Rossi¹⁰, F. Zhao²⁶, K.M. Zurek^{21,22}

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⁷INFN Sezione di Napoli, Napoli, Italy

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PTOLEMY (PonTecorvo Observatory for Light, Early-universe, Massive-neutrino Yield) @ LNGS

The expected rate is:

$$\Gamma_{C\nu B} = [n_0(\nu_{h_R}) + n_0(\nu_{h_L})] N_T \bar{\sigma} \sum_{i=1}^3 |U_{ei}|^2 f_c(m_i),$$

For unclustered neutrinos (i.e. $f_c = 1$) and 100 g of tritium, the expected number of events per year:

$$\Gamma_{C\nu B}^D \simeq 4 \text{ yr}^{-1}, \quad \Gamma_{C\nu B}^M = 2\Gamma_{C\nu B}^D \simeq 8 \text{ yr}^{-1}$$

If neutrinos are Majorana particles, the expected number of events is doubled with respect to the Dirac case. The reason is related to the fact that, during the transition from ultra-relativistic to non-relativistic particles, helicity is conserved, but not chirality. The population of relic neutrinos is then composed by left- and right-helical neutrinos in the Majorana case, and only left-helical neutrinos in the Dirac case. Since the neutrino capture can only occur for left-chiral electron neutrinos, the fact that in the Majorana case the right-handed neutrinos can have a left-chiral component leads to a doubled number of possible interactions.

masses (meV)	matter halo	overdensity f_c {best fit best fit + baryons}	$\Gamma_{C\nu B}^D$ (yr $^{-1}$) best fit best fit + baryons	$\Gamma_{C\nu B}^M$ (yr $^{-1}$) optimistic
any	any	no clustering	4.06	8.12
degenerate $m_{\nu_{1,2,3}} = 150$	NFW	2.18 2.44 2.88	8.8 9.9 11.7	17.7 19.8 23.4
	Einasto	1.68 1.87 2.43	6.8 7.6 9.9	13.6 15.1 19.7
minimal (IO) $m_{\nu_3} = 60$	NFW	1.15 1.18 1.21	4.07 4.08 4.08	8.15 8.15 8.16
	Einasto	1.09 1.12 1.18	4.07 4.07 4.08	8.14 8.14 8.15
minimal (NO) $m_{\nu_{1,2}} = 60$	NFW	1.15 1.18 1.21	4.66 4.78 4.89	9.31 9.55 9.77
	Einasto	1.09 1.12 1.18	4.42 4.54 4.78	8.84 9.07 9.55

$$r_s(z_{drag}) = \int_0^{\eta(z)} d\eta c_s(1+z) , \quad (16)$$

where $c_s = 1/\sqrt{3(1+R)}$ is the sound speed, and $R \equiv 3\rho_b/4\rho_\gamma$. The drag epoch corresponds to the redshift at which the *drag* optical depth τ_d is equal to one:

$$\tau_d(z_{drag}) = \int_0^{z_{drag}} dz \frac{d\eta}{da} \frac{x_e(z)\sigma_T}{R} \equiv 1 . \quad (17)$$

COLD dark matter

HOT dark matter

HDM

CDM

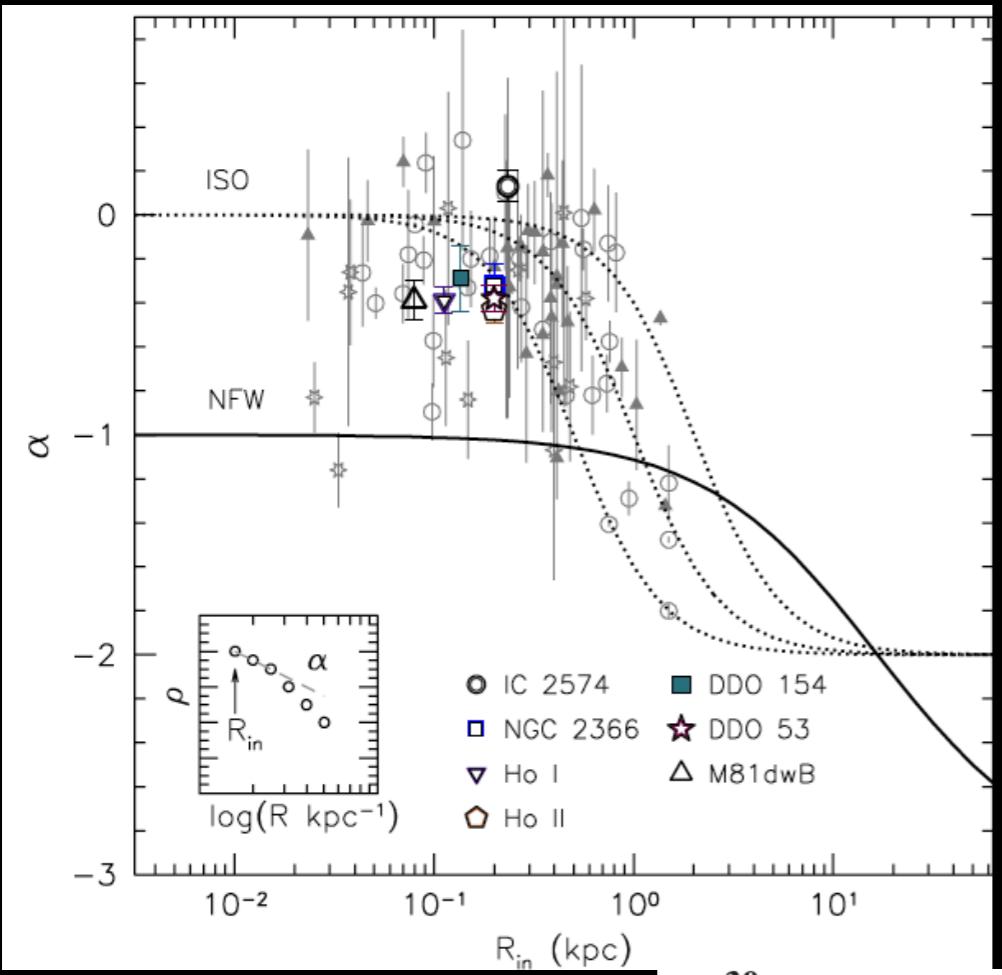
WDM

WARM dark matter

CREATED BY

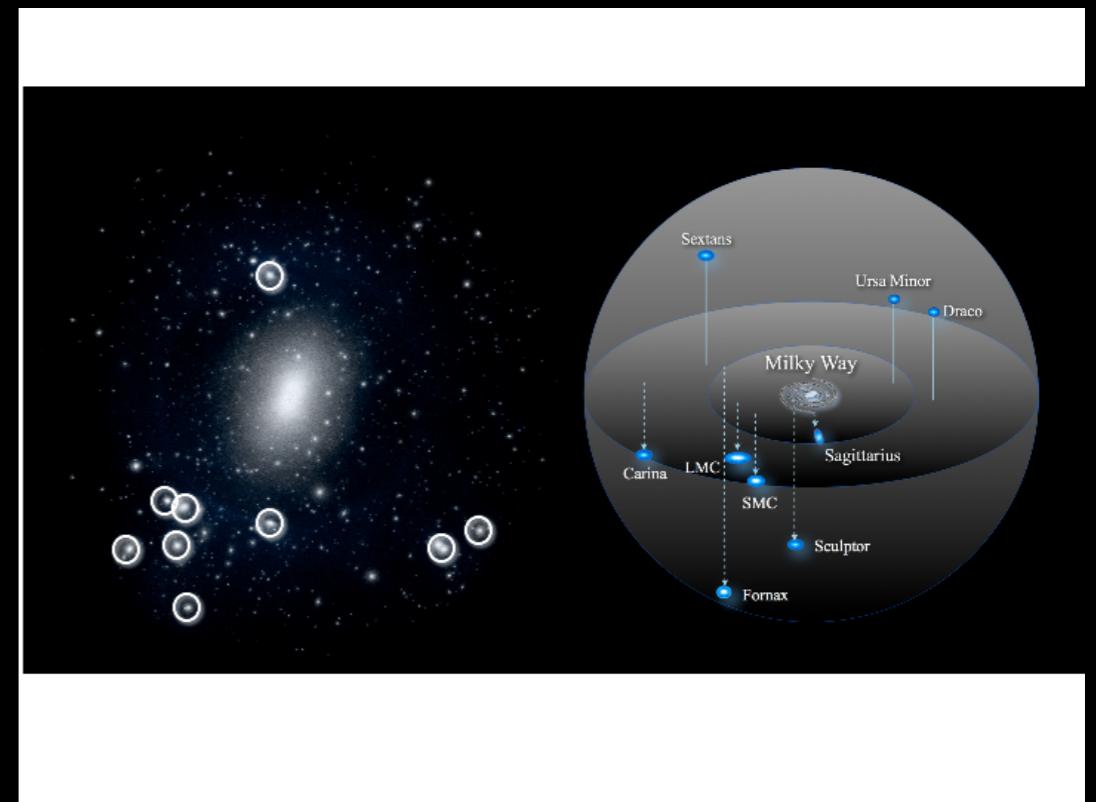
JOSEPH MELLOZZI

Small scale crisis of Λ CDM@galactic and sub-galactic scales



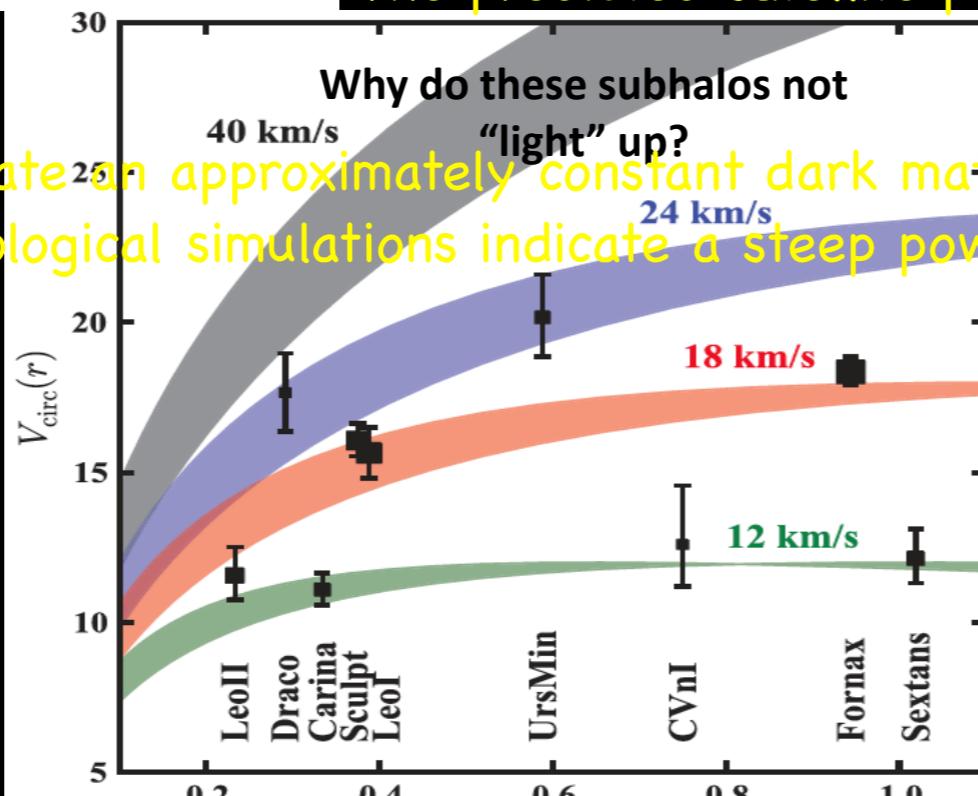
Core/Cusp problem

Observations seem to indicate an approximately constant dark matter density in the inner parts of galaxies (core), while cosmological simulations indicate a steep power-law-like behaviour (cusp) problem



Missing satellite problem

The predicted satellite population far exceeds the observed one



Massive dark subhalos are too dense to match data.
Expected 10 subhalos in the Milky Way with $v > 30$ km/s, only 3 known!

(Boylan et al, MNRAS'11)

Sterile keV ($0.01 m_e$) neutrino as a warm dark matter candidate?

A controversial unidentified line has been detected at with a significance $> 3\sigma$ in two independent samples of X-ray clusters with XMM-Newton.



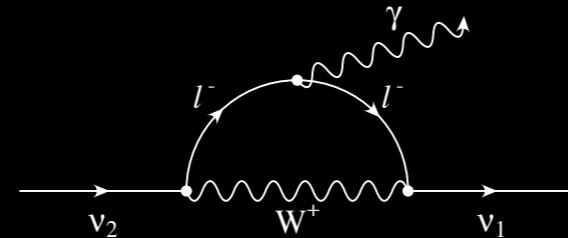
It is independently seen by the same group in the Perseus Cluster with Chandra data.



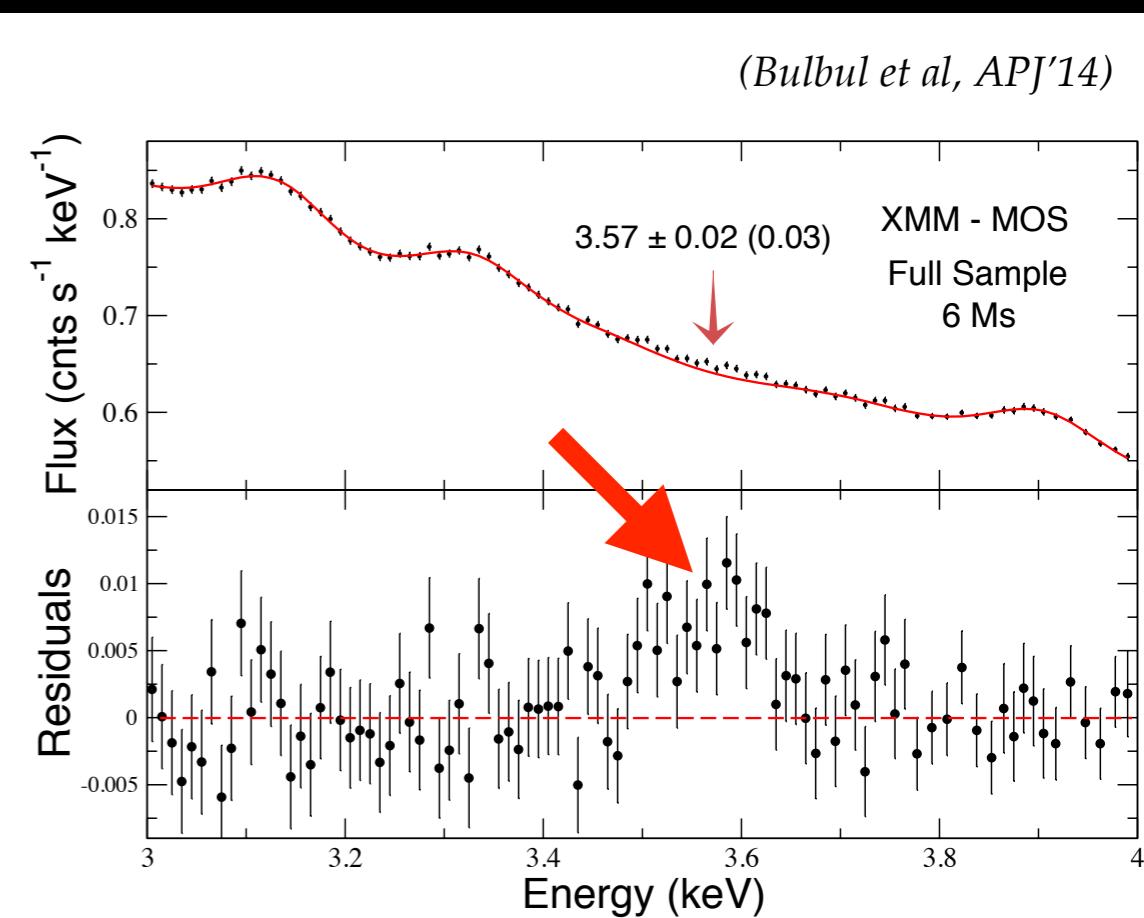
(Bulbul et al, APJ'14)

An independent group finds a line at the same energy toward Andromeda and Perseus with XMM-Newton, with a combined statistical evidence of 4.4σ . (Boyarsky et al, PRL'14)

$$\nu_s \rightarrow \nu_\alpha + \gamma$$

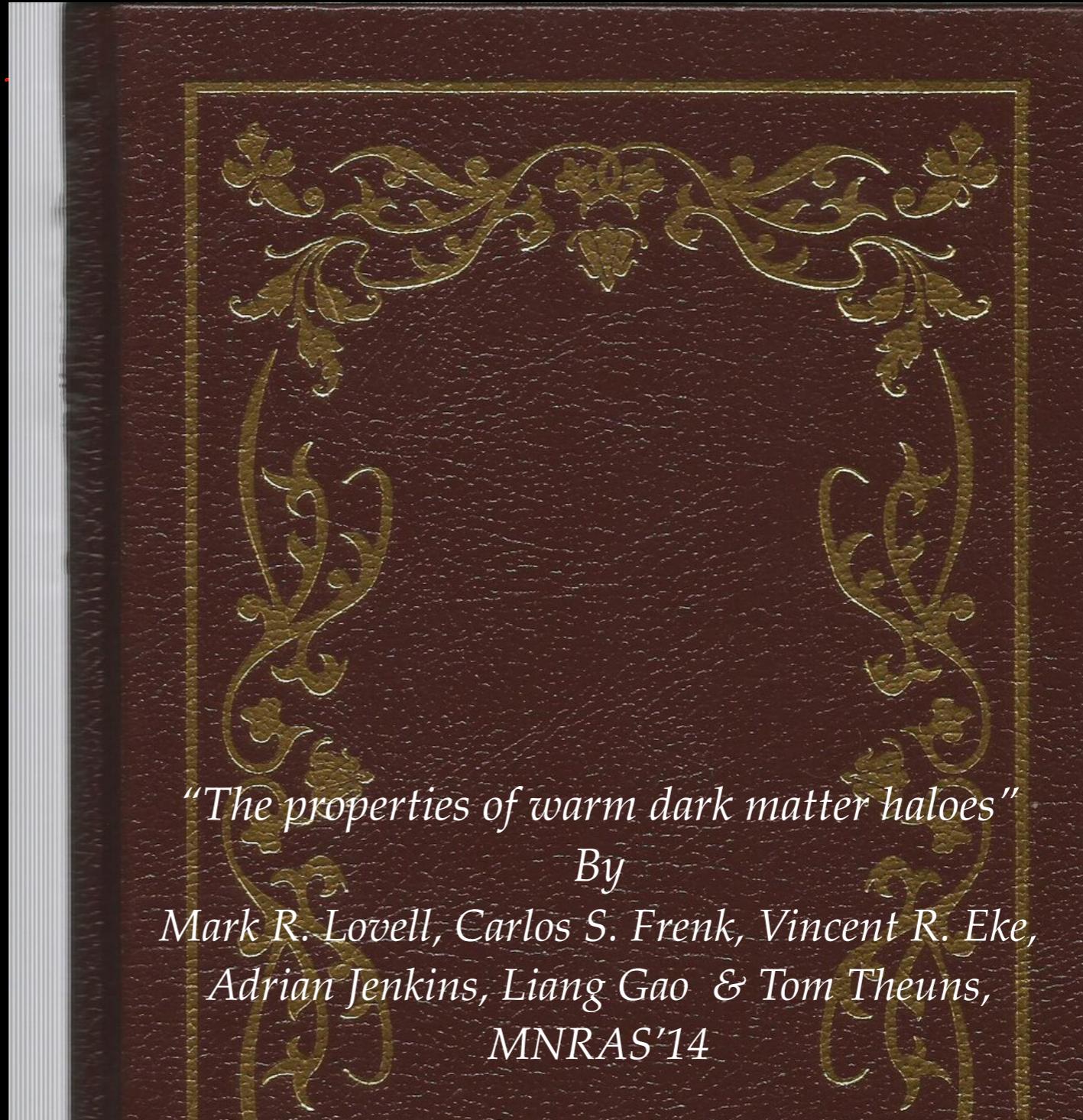


$$m_s = 2E = 7.1 \text{ keV}$$



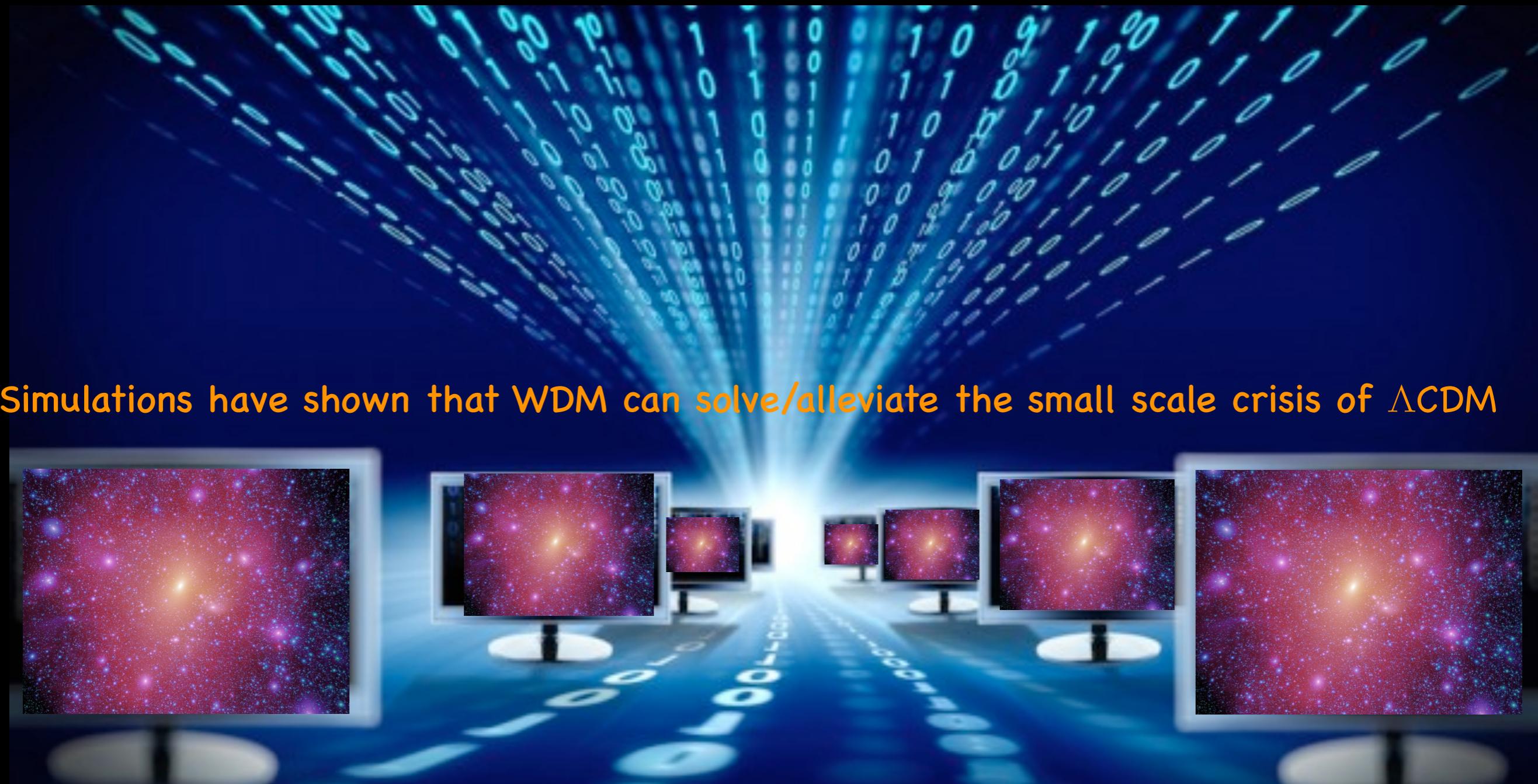
Sterile keV ($0.01 m_e$) neutrino as a **warm** dark matter candidate?

WDM leads to an identical large scale structure pattern than CDM, but very different subhaloes abundance, structure and dynamics: the free streaming of a keV sterile neutrino will reduce power at the small scales, delaying structure formation and lowering the haloes concentration.



Sterile keV ($0.01 m_e$) neutrino as a warm dark matter candidate?

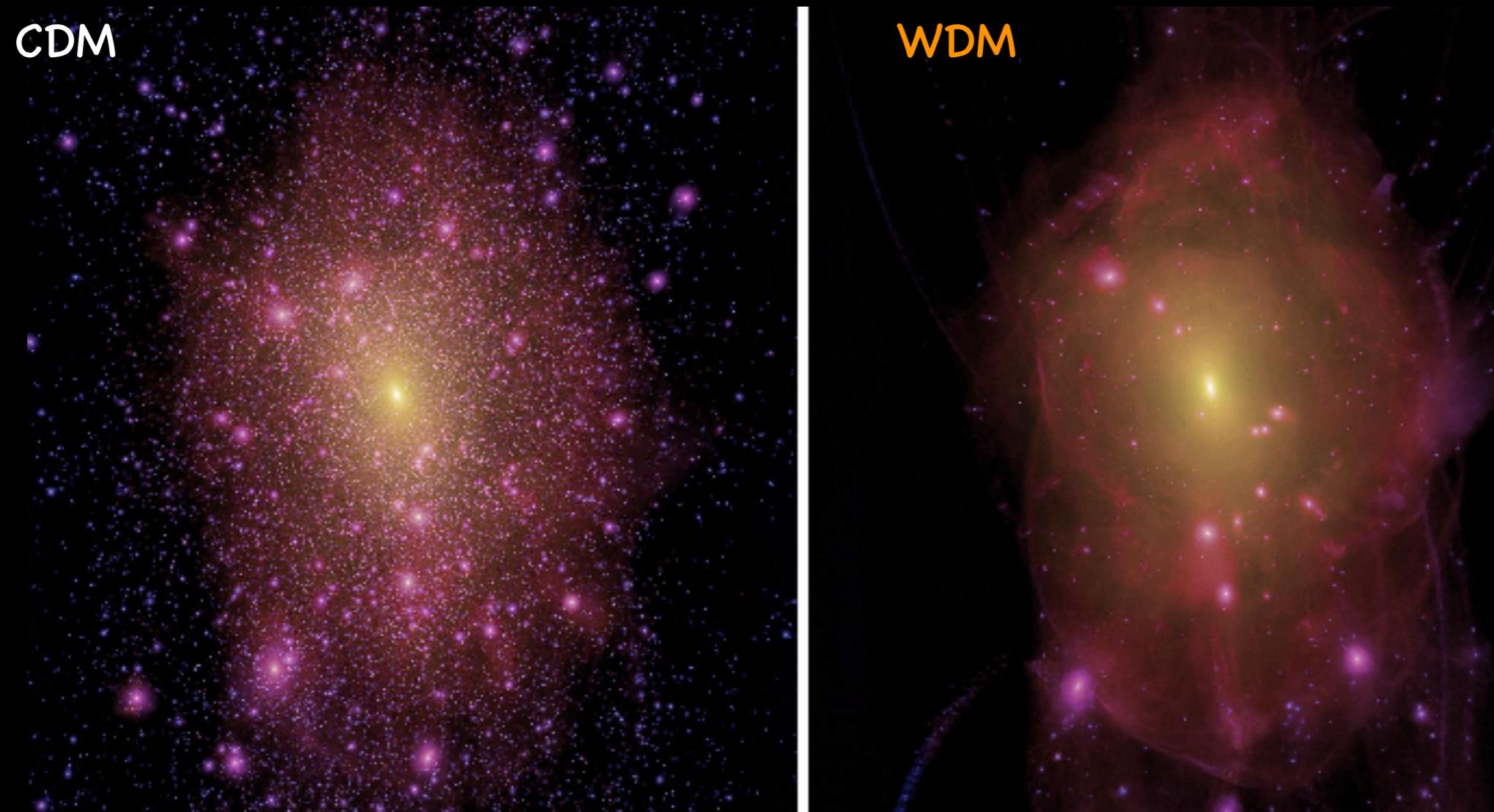
WDM leads to an identical large scale structure pattern than CDM, but very different subhaloes abundance, structure and dynamics: the free streaming of a keV sterile neutrino will reduce power at the small scales, delaying structure formation and lowering the haloes concentration.



Simulations have shown that WDM can solve/alleviate the small scale crisis of Λ CDM

Sterile keV ($0.01 m_e$) neutrino as a **warm** dark matter candidate?

WDM could reconcile theory with observations!



"The Haloes of Bright Satellite Galaxies in a Warm Dark Matter Universe", Mark R. Lovell, Vincent R. Eke, Carlos S. Frenk, Liang Gao, Adrian Jenkins, Jie Wang, D.M. White, Alexey Boyarsky & Oleg Ruchayskiy MNRAS'12

"The properties of warm dark matter haloes", Mark R. Lovell, Carlos S. Frenk, Vincent R. Eke, Adrian Jenkins, Liang Gao & Tom Theuns, MNRAS'14

Big Bang Nucleosynthesis: N_{eff}

N_{eff} changes the freeze out temperature of weak interactions:

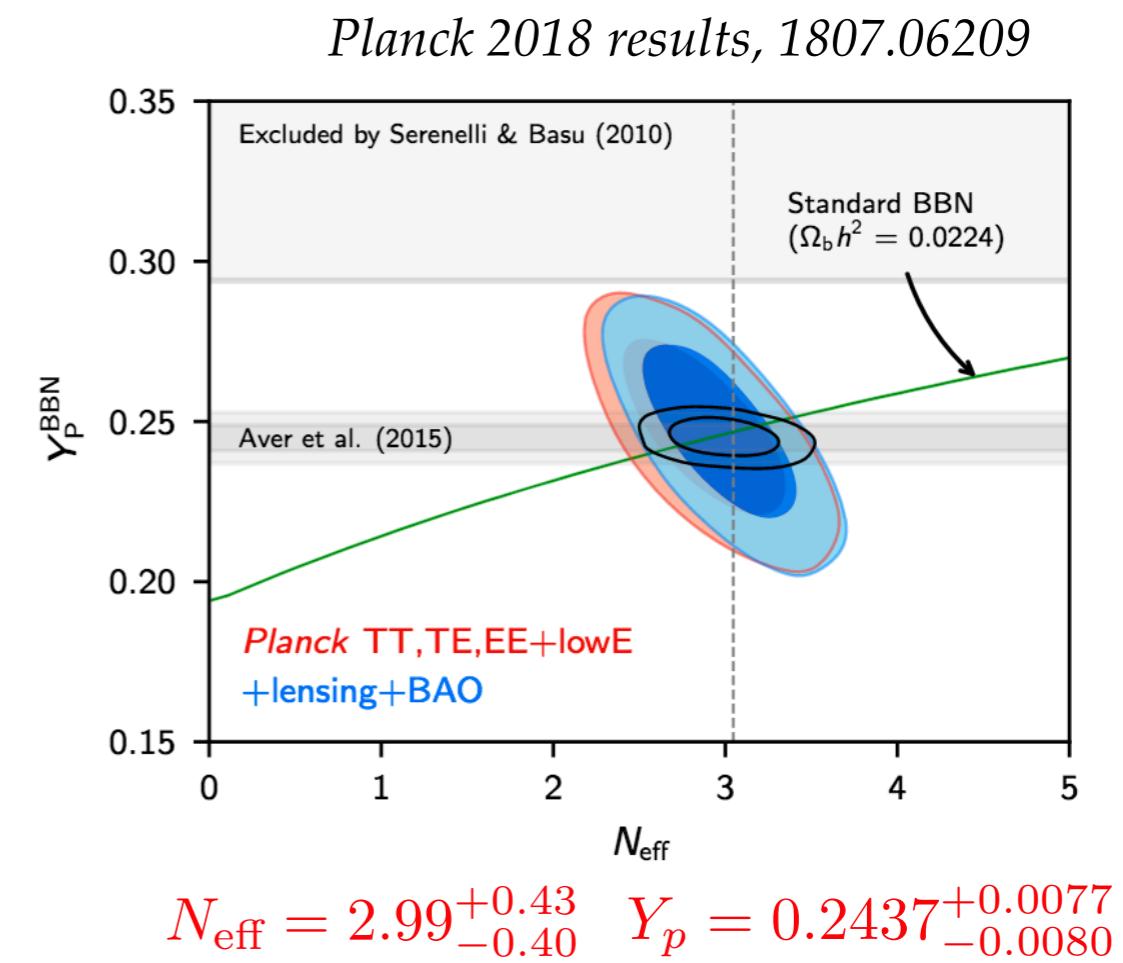
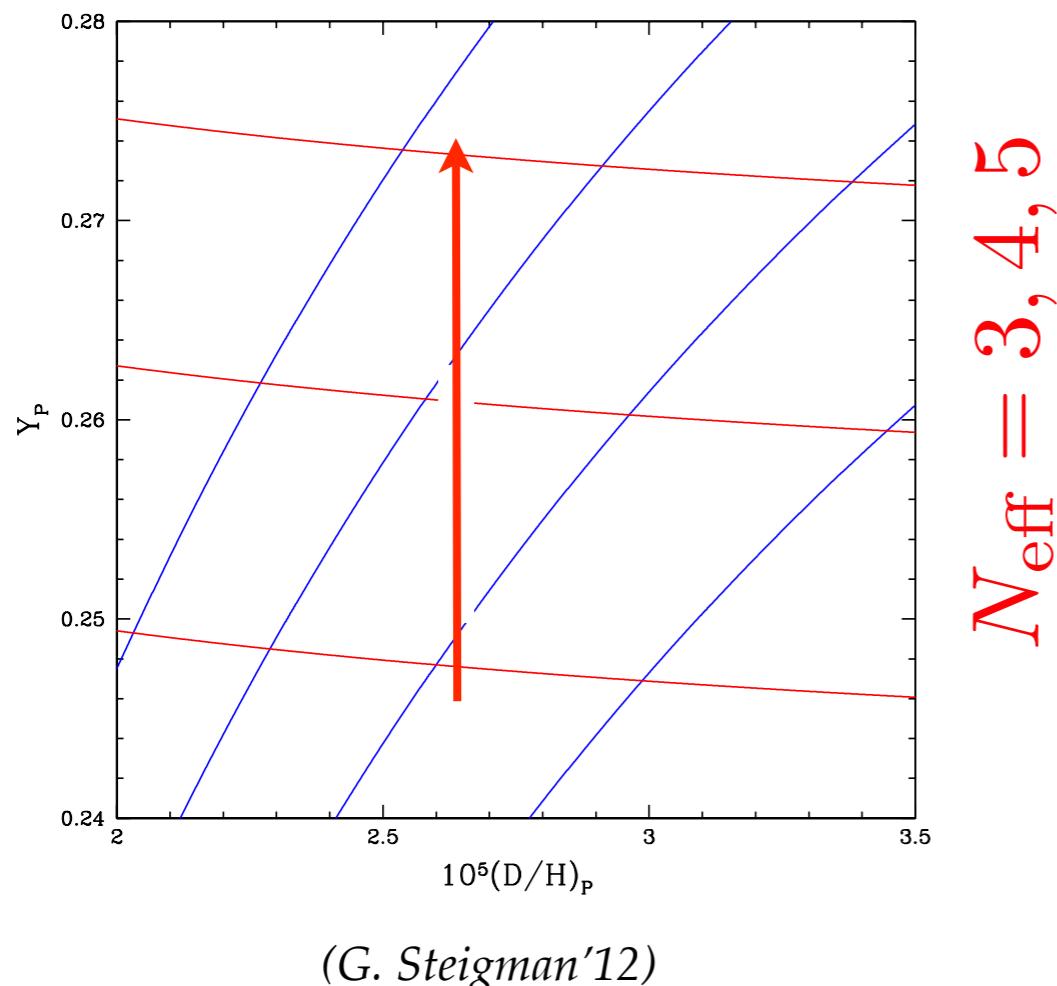
$$\Gamma_{n \leftrightarrow p} \sim H$$

MORE NEUTRINOS:

Higher N_{eff} : larger expansion rate & freeze out temperature, MORE HELIUM 4

$$n/p \simeq e^{-\frac{m_n - m_p}{T_{\text{freeze}}}}$$

$$Y_p = \frac{2(n/p)}{1 + n/p}$$



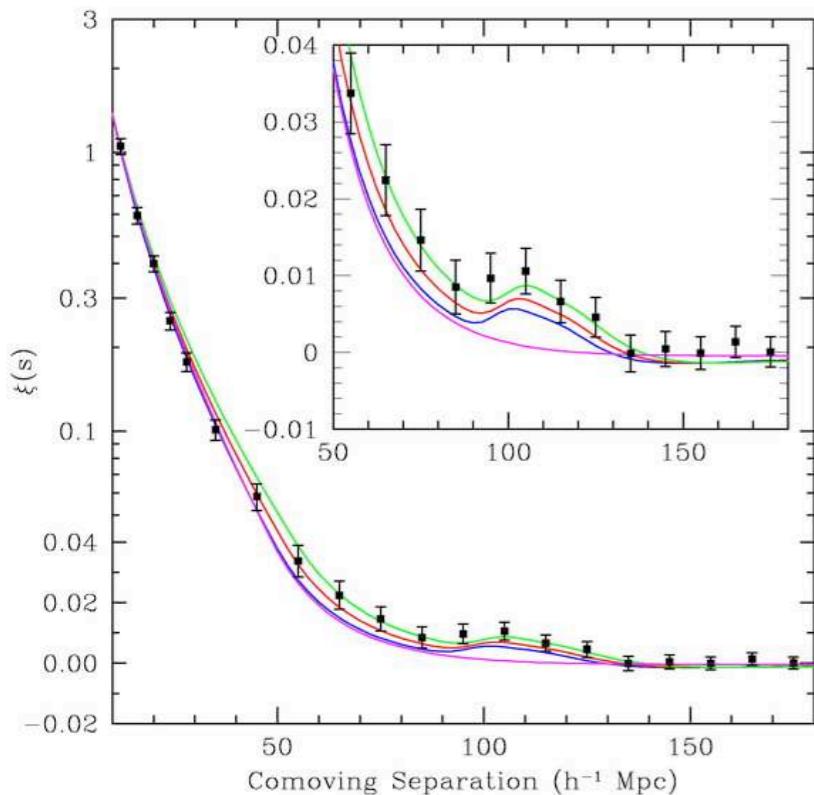
Los catálogos de galaxias miden la función de correlación:

$$\xi(\vec{r}) \equiv \langle \delta(\vec{x})\delta(\vec{x} + \vec{r}) \rangle_{\text{Volume}}$$

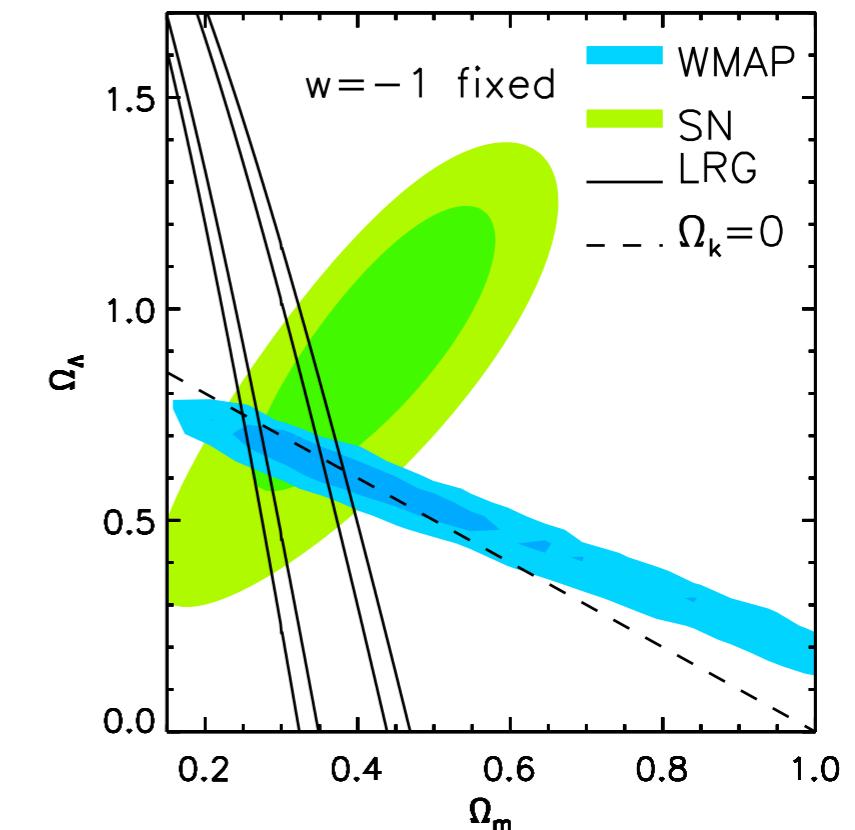
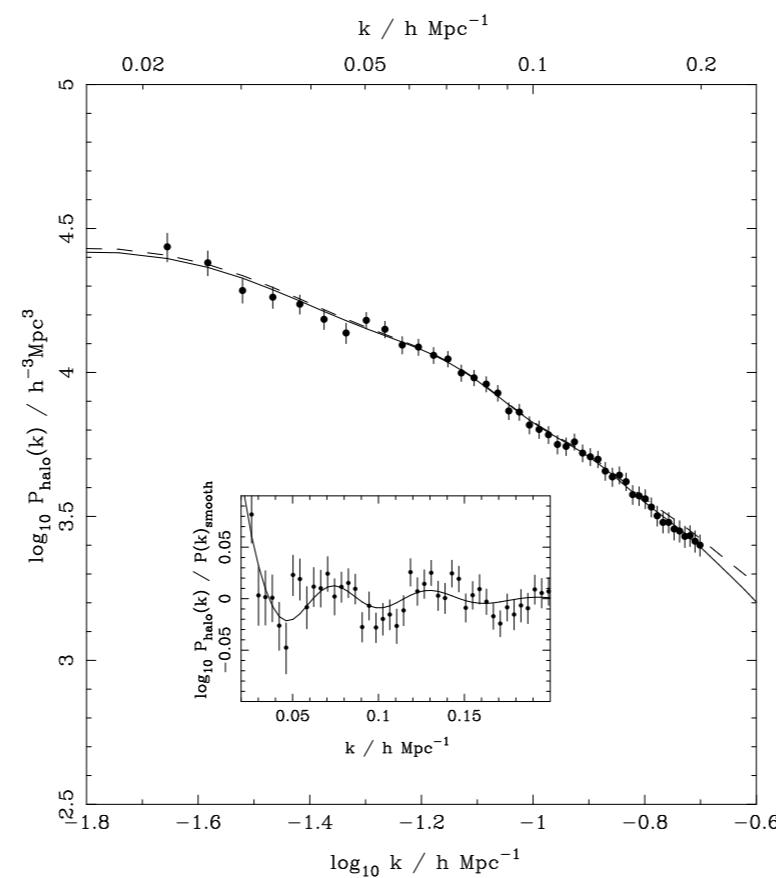
$$\langle \tilde{\delta}(\vec{k})\tilde{\delta}(\vec{k}') \rangle_{\text{Volume}} = (2\pi)^3 P(k)\delta^3(\vec{k} - \vec{k}')$$

$$\delta(\vec{x}) \equiv \frac{\rho(\vec{x}) - \bar{\rho}(\vec{x})}{\bar{\rho}(\vec{x})}$$

$$\tilde{\delta}(\vec{k}) \equiv \int d^3\vec{r} e^{i\vec{k}\vec{r}} \delta(\vec{r})$$



Eisenstein et al'05



Reid et al'09

SSDS 2005: Primera detección de la señal BAO (3.4s) (47000 LRGs, 4000 deg² , z=0.35)

SDSS II 2009: 110 000 LRGs, 8000 deg² , z=0.35.